#### **Interpreting Higgs results**

Adam Falkowski LPT Orsay

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Carmi,AA,Kuflik,Volansky [1202.3144] AA,Rychkov,Urbano [1202.1532] + oven fresh unpublished



#### 2 Higgs - observations







#### Who broke electroweak symmetry?

- Since LEP we know for a fact fundamental interactions of matter obey  $SU(2) \times U(1)$  local symmetry that is however spontaneously broken (non-linearly realized), as W, Z and fermions have masses
- The Question for the LHC is the precise nature of electroweak symmetry breaking
- More rigorously, the question is what sort of physics stops the growth of the scattering amplitudes of W and Z bosons:
  - In the SM (without Higgs) the tree-level amplitude for longitudinally polarized W's and Z's grows with energy,  ${\cal M}\sim s/v^2$
  - Unitarity requires  ${\rm Re}\,\bar{\cal M}^J < 1/2$  for all partial waves. Perturbative unitarity is lost at TeV
  - Something else must enter before that scale!

#### **Options for Electroweak Symmetry Breaking**

3 basic possibilities. Unitarity saved by

- Non-Perturbative effects in the SM (no concrete framework so far)
- Strongly Coupled: composite vectors and/or scalars to WW and WZ
- Weakly Coupled: fundamental scalar coupled to WW and ZZ, otherwise known as the Higgs
- $\bullet$  ...or a combination of the above, for example Composite Higgs weakly coupled up to  $\sim$  3 TeV, then strongly coupled



picture stolen from C. Grojean

- Current experimental data strongly suggest that the weakly coupled option is approximately true, at least for  $E\lesssim 1$  TeV, and likely up to much higher scales
  - Electroweak precision tests
  - No new vector or tensor states observed at the Tevatron and LHC
  - Higgs-like excess near 125 GeV
- Furthermore, they point to the simplest realization with a single Higgs boson resposible for unitarizing WW scattering
  - Approximate global symmetries of SM, such as flavor and CP seem to be very well preserved

#### (almost) Unshakable Arguments

- $\blacksquare$  Observed neutrino masses imply new physics (at least, right-handed neutrinos) somewhere between 1 keV and  $10^{15}~{\rm GeV}$
- 0 Existence of dark matter requires new physics somewhere between sub-eV and  $10^{19}~\text{GeV}$
- 0 Domination of matter over anti-matter requires new physics between 100 GeV and  $10^{16}~\text{GeV}$

unfortunately, none of above guarantees new physics showing up in LHC

#### Some Esthetic Arguments

- Fermion masses and mixings suggests another sector generating the observed structures, at any scale above TeV and Planck
- $\bullet$  Approximate unification of gauge couplings suggests new states at any scale between 100 and  $10^{14}~{\rm GeV}$
- $\bullet$  Higgs potential metastability suggests new physics between 100 GeV and  $10^{10}~{\rm GeV}$
- Instability of Higgs mass against radiative corrections suggests new states at 100 GeV

only one, somewhat shaky argument clearly points to new physics in LHC

$$\delta m_H^2 = \cdots \delta m + \cdots \delta m$$

- Hierarchy problem dominated model building for last 30 years
- Two important classes of solutions
  - Supersymmetry: fermion-boson cancellation, may be weakly coupled up to Planck scale
  - Composite/Little Higgs: boson-boson or fermion-fermion cancellation, weakly coupled up to 3-10 TeV, then strongly coupled
- All existing models introduce a multitude of new particles at weak scale, and require serious conspiracy why they preserve approximate accidental symmetries of the SM, to avoid showing up indirectly in numerous precision measurements
- Typically, in specific realizations advertised as natural one has 1-0.1% fine-tuning, after experimental constraints are taken into account

Fermionic top partners T

- Limits depending on dominant decay
- Constraints on T → bW channel (typically 50% branching ratio in models without T-parity) and on T → t+MET (expected in models with T-parity)
- Current limits on mass around 400 500 GeV
- Naturalness under stress, but not completely dead yet...



Scalar top partners  $\tilde{t}$ 

- In generic SUSY  $m_{ ilde{t}}\gtrsim 1~{
  m TeV} 
  ightarrow$  serious fine-tuning problem
- But, for  $m_{\tilde{t}} \ll m_{\tilde{q}}$  and  $m_{\tilde{t}} \ll m_{\tilde{g}}$  limits become much weaker
- Currently only theorist-level robust limit on stops,  $m_{\tilde{t}} \gtrsim 150 250$  GeV, depending on decay mode and LSP mass Papucci et al [1110.6926]
- Related limits on direct sbottom production from ATLAS [1112.3832]
- Reasonanble fine-tuning still possible if stops and sbottom are only colored superpartner below TeV



- Naturalness window still half open
- But no experimental hint of a larger framework just around the corner

Dominant attitude in theory:

- Hierarchy problem may or may not be relevant
- Model building now dominated by LHC data, not theory prejudice





#### Hierarchy problem and Higgs physics





stolen from R. Rattazzi

- The SM Higgs with mass  $m_h \sim 125~{\rm GeV}$  has many decay channels that are potentially observable at the LHC and Tevatron
  - Now:  $H \rightarrow ZZ^*$ ,  $H \rightarrow \gamma\gamma$ , and  $H \rightarrow b\bar{b}$
  - Shortly:  $H \rightarrow WW^*$
  - $\sim 1$  year perspective:  $H o au^+ au^-$
- Also different production channels can be isolated
  - Now: gluon fusion and (maybe) vector boson fusion
  - Longer Perspective: W/Z and  $t\bar{t}$  associated production
- Rich Higgs physics available in near future
- If new physics exists, Higgs interactions likely to be modified
- If new physics restores naturalness, Higgs interactions are necessarily modified
- Measuring Higgs rates at the LHC may be the shortest route to new physics!

# Higgs Observations



- Significant background, but great mass resolution
- Both ATLAS and CMS observe an excess near  $m_h \sim 125$  GeV, ATLAS centered at 126 and CMS centered at 125
- In both case the best fit cross section at the peak exceeds the SM value, though the latter is well within uncertainties
- CMS also observes an excess in inclusive  $\gamma\gamma jj$  channel dominated by VBF production mode, corresponding to cross section well exceeding the SM one (though, again, uncertainties are still large)





- Very low background, great mass resolution
- ATLAS has 3 events at  $m_{4/} pprox$  124 GeV
- CMS has 2 events at  $m_{4/} \approx 126 \text{ GeV}$



	Signal	WW	$WZ/ZZ/W\gamma$	tī	tW/tb/tqb	$Z/\gamma^*$ + jets	W + jets	Total Bkg.	Obs.
$5 m_H = 125 \text{ GeV}$	$25 \pm 7$	$110 \pm 12$	$12 \pm 3$	$7 \pm 2$	$5 \pm 2$	$13 \pm 8$	$27 \pm 16$	$173 \pm 22$	174
$\dot{o} m_H = 240 \text{ GeV}$	$60 \pm 17$	$432\pm49$	$24 \pm 3$	$68 \pm 15$	$39 \pm 9$	$8 \pm 2$	$36 \pm 24$	$607 \pm 63$	629
$\underline{5} m_H = 125 \text{ GeV}$	6 ± 2	$18 \pm 3$	6 ± 3	7 ± 2	4 ± 2	6 ± 1	$5 \pm 3$	45 ± 7	56
$-m_H = 240 \text{ GeV}$	$23 \pm 9$	$99 \pm 22$	$8 \pm 1$	$73 \pm 27$	$35 \pm 19$	$6 \pm 2$	$7 \pm 7$	$229 \pm 55$	232
$\underline{\mathbf{v}}_{H} = 125 \text{ GeV}$	$0.4 \pm 0.2$	$0.3 \pm 0.2$	negl.	$0.2 \pm 0.1$	negl.	$0.0 \pm 0.1$	negl.	$0.5 \pm 0.2$	0
$\dot{\bigtriangledown} m_H = 240 \text{ GeV}$	$2.5 \pm 0.6$	$1.1 \pm 0.7$	$0.1 \pm 0.1$	$2.6 \pm 1.3$	$0.3 \pm 0.3$	negl.	$0.1\pm0.1$	$4.2 \pm 1.7$	2

- Significant background, poor mass resolution, better for exclusion than discovery
- No clear excess here, which begins to feel weird
- Bad luck, background misestimation, or something interesting going on?



- Low mass range excluded by Tevatron and LHC except for 122-127 GeV range
- Even lower mass range excluded by LEP,
- High mass range excluded by LHC, or highly disfavored by EWPT



- Slightly too much signal in  $\gamma\gamma$  channel
- Slightly too little signal in WW and ZZ channels
- Overall good consistency with SM Higgs predictions



- Broad  $\sim$  3 sigma excess in low mass range, mostly originating from  $\sim$  40 excess events in  $W/Z + H \rightarrow b\bar{b}$  production mode in CDF
- Consistent with 120-140 GeV Higgs



- Points to somewhat enhanced rate in VH production channel, the heavier Higgs, the larger cross section boost is needed
- Doesn't strongly favor any mass between 120 and 135 GeV

### Announcement

# The Higgs boson has been discovered at has the mass near 125 GeV

Next Level Is it the SM Higgs? Higgs Theory

#### Higgs effective theory

Define effective Higgs Lagrangian at  $\mu\approx m_h\sim 125\,GeV.$  Couplings relevant for current LHC data

$$\mathcal{L}_{eff} = c_{V} \frac{2m_{W}^{2}}{v} h W_{\mu}^{+} W_{\mu}^{-} + c_{V} \frac{m_{Z}^{2}}{v} h Z_{\mu} Z_{\mu} - c_{b} \frac{m_{b}}{v} h \bar{b} b - c_{\tau} \frac{m_{\tau}}{v} h \bar{\tau} \tau$$
$$+ c_{g} \frac{\alpha_{s}}{12\pi v} h G_{\mu\nu}^{a} G_{\mu\nu}^{a} + c_{\gamma} \frac{\alpha}{\pi v} h A_{\mu\nu} A_{\mu\nu}$$

 Only one theoretical prejudice: custodial isospin requires same Higgs coupling to W and Z. Otherwise for c<sub>W</sub> ≠ c<sub>Z</sub>:

$$\Delta T = (c_W^2 - c_Z^2) \alpha_{EM}^{-1} \frac{\Lambda^2}{16\pi^2 v^2}$$

For  $\Lambda \sim 1$  TeV  $c_W$  and  $c_Z$  have to within 1%

- For time being already assume that  $c_{ au} = c_b$  (until better au au data arrive)
- $\bullet$  Top already integrated out, contributing to  $c_g$  and  $c_\gamma$
- SM predicts  $c_V=c_b=c_g=1$  and  $c_\gamma=2/9$
- Any of the couplings can be modified in specific scenarios beyond the SM
- All LHC Higgs rates can be easily expressed as functions of the c<sub>i</sub> couplings

The decay widths of the Higgs relative to the SM predictions are modified approximately as,

$$\frac{\Gamma(h \to b\bar{b})}{\Gamma_{SM}(h \to b\bar{b})} \simeq |c_b|^2$$

$$\frac{\Gamma(h \to WW^*)}{\Gamma_{SM}(h \to WW^*)} = \frac{\Gamma(h \to ZZ^*)}{\Gamma_{SM}(h \to ZZ^*)} \simeq |c_V|^2$$

$$\frac{\Gamma(h \to gg)}{\Gamma_{SM}(h \to gg)} \simeq |c_g|^2$$

$$\frac{\Gamma(h \to \gamma\gamma)}{\Gamma_{SM}(h \to \gamma\gamma)} \simeq \left|\frac{\hat{c}_{\gamma}}{\hat{c}_{\gamma,SM}}\right|^2 \qquad (1)$$

where, taking into account W loop and assuming  $m_h \approx 125$ ,  $\hat{c}_\gamma \approx c_\gamma - c_V$ , and  $\hat{c}_{\gamma,SM} \approx -0.8$ 



For  $m_h \sim 125~{
m GeV}$  total Higgs width scales as

$$rac{\Gamma(h)}{\Gamma_{SM}(h)}\simeq 0.65c_b^2+0.25c_V^2+0.1c_g^2$$

Assuming  $H \rightarrow bb$  dominates Higgs widths

$$R_{V} \equiv \frac{\sigma(pp \to h)\mathrm{Br}(h \to ZZ^{*})}{\sigma_{SM}(pp \to h)\mathrm{Br}_{SM}(h \to ZZ^{*})} \simeq \left|\frac{c_{g}c_{V}}{c_{b}}\right|^{2},$$

$$R_{\gamma} \equiv \frac{\sigma(pp \to h)\mathrm{Br}(h \to \gamma\gamma)}{\sigma_{SM}(pp \to h)\mathrm{Br}_{SM}(h \to \gamma\gamma)} \simeq \left|\frac{c_{g}\hat{c}_{\gamma}}{\hat{c}_{\gamma,SM}c_{b}}\right|^{2},$$

$$R_{\gamma,VBF} \equiv \frac{\sigma(pp \to hjj)\mathrm{Br}(h \to \gamma\gamma)}{\sigma_{SM}(pp \to hjj)\mathrm{Br}_{SM}(h \to \gamma\gamma)} \simeq \left|\frac{c_{V}\hat{c}_{\gamma}}{\hat{c}_{\gamma,SM}c_{b}}\right|^{2}.$$

$$R_{\mathrm{Tev}} \equiv \frac{\sigma(p\bar{p} \to Vh)\mathrm{Br}(h \to b\bar{b})}{\sigma_{SM}(p\bar{p} \to Vh)\mathrm{Br}_{SM}(h \to b\bar{b})} \simeq c_{V}^{2}, \qquad (2)$$

$$\mathcal{L}_{eff} = c_{V} \frac{2m_{W}^{2}}{v} h W_{\mu}^{+} W_{\mu}^{-} + c_{V} \frac{m_{Z}^{2}}{v} h Z_{\mu} Z_{\mu} - c_{b} \frac{m_{b}}{v} h \bar{b} b - c_{b} \frac{m_{\tau}}{v} h \bar{\tau} \tau$$
$$+ c_{g} \frac{\alpha_{s}}{12\pi v} h G_{\mu\nu}^{a} G_{\mu\nu}^{a} + c_{\gamma} \frac{\alpha}{\pi v} h A_{\mu\nu} A_{\mu\nu}$$

- We will find the region of effective theory parameter space favored by 2011 LHC Higgs data
- Interesting to check whether the current LHC data are consistent with the SM Higgs
- Also interesting, whether they favor or disfavor any particular BSM scenario
- Of course at this stage one cannot make very strong statements about Higgs couplimgs (some of you don't even think Higgs has been discovered)
- Consider it a warm-up exercise, in preparation for serious signals
- Recently Carmi [1202.3144], Azatov [1202.3415] and Espinosa [1202.3697]

#### Illegal ATLAS/CMS combination



Carmi et al [1202.3144], Moriond updates not included, bands are 1 sigma

#### Fits assuming $m_h = 125$ GeV

Preliminary



- Only dimension-5 Higgs couplings allowed to vary
- On this plane Tevatron never within 1 sigma band
- Combined =  $\chi^2 < 6.25$  (full 4-parameter  $\Delta \chi^2$  plots in preparation)

#### Fits assuming $m_h = 125$ GeV

Preliminary



- Composite Higgs inspired parametrization
- Couplings to fermions and gauge boson allowed to vary independently

#### Fits assuming $m_h = 125$ GeV

Preliminary



 Higgs coupling to EW gauge bosons, and dimension 5 effective Higgs coupling to gluons allowed to vary

• Top partner models relation 
$$c_{\gamma} = 2c_g/9$$

#### Scalar partner toy model

- Very toy "natural" model: just one scalar top partner (this is not SUSY, where at least two scalar partners are needed)
- Top partner interactions with Higgs to cancel top quadratic divergences

$$-\left(yHQt^{c}+\mathrm{h.c.}\right)-|\tilde{t}|^{2}\left(M^{2}+2y^{2}|H|^{2}\right)$$

- Only one free parameter: top partner mass  $m_{\tilde{t}}^2 = M^2 + y^2 v^2$
- New contributions to effective dimension 5 Higgs interactions

$$rac{c_{
m g}}{c_{
m g,SM}} = rac{c_{\gamma}}{c_{\gamma,{
m SM}}} \simeq 1 + rac{m_t^2}{2m_t^2}$$



#### Fermion partner model

- For fermionic top partner, non-renormalizable interactions with Higgs needed to cancel top quadratic divergence
- Simple model inspired by T-parity conserving Little Higgs

 $-(y\sin(|H|/f)Qt^{c} + h.c.) - yf\cos(|H|/f)TT^{c}$ 

- Again only one free parameter: top partner mass  $m_T = yf \cos(v/\sqrt{2}f)$
- New contributions to effective dimension 5 Higgs interactions

$$rac{c_g}{c_{g,\mathrm{SM}}} = rac{c_\gamma}{c_{\gamma,\mathrm{SM}}} \simeq 1 - rac{m_t^2}{m_T^2}$$





- Beginning of a beautiful friendship
- More Higgs data from LHC may favor/disfavor particular BSM scenarios...
- ...or just confirm the SM again

### What If ?



- Current combined Higgs data allow, while Tevatron and VBF  $\gamma\gamma$  channel in CMS favor increased Higgs coupling to WW and ZZ
- What if indeed  $c_V > 1$ ?

- If SM Higgs doublet mixes with a singlet or another doublet, then always  $c_V = \cos \alpha < 1$ . Thus enhancement impossible in typical SUSY models.
- For Higgs being a pseudo-Goldstone boson of any compact coset (Little Higgs and composite Higgs), also  $c_V = \cos(v/f) < 1$ . Again, enhancement of  $c_V$  impossible
- Low et al [0907.5413]: sum rule proving  $c_V > 1$  implies charge-2 Higgs
- AA et al [1202.1532] : stronger sum rule (assuming custodial symmetry)

$$1 - c_V^2 \approx rac{v^2}{6\pi} \int_0^\infty rac{ds}{s} \left( 2\sigma_{I=0}^{
m tot}(s) + 3\sigma_{I=1}^{
m tot}(s) - 5\sigma_{I=2}^{
m tot}(s) 
ight).$$

•  $c_V > 1$  implies enhancement of isospin 2 channel of WW scattering

#### Simplest realization of isospin 2 enhancement

- Quintuplet of weakly coupled scalars  $Q = (Q^{--}, Q^{-}, Q^{0}, Q^{+}, Q^{++})$
- · Coupled to electroweak gauge bosons in custodially invariant way

$$\frac{g_{Q}}{v}\left\{\sqrt{\frac{2}{3}}Q^{0}\left(m_{W}^{2}W_{\mu}^{+}W_{\mu}^{-}-m_{Z}^{2}Z_{\mu}^{2}\right)+\left(Q^{++}m_{W}^{2}W_{\mu}^{-}W_{\mu}^{-}+\sqrt{2}Q^{+}m_{W}m_{Z}W_{\mu}^{-}Z_{\mu}+hc\right)\right\}$$

Sum rule fulfilled for

$$g_Q^2 = \frac{6}{5} \left( c_V^2 - 1 \right)$$

- What is special about  $g_Q^2 = 6/5(c_V^2 1)$  ?
- Quintuplet, much like Higgs, contributes to WW scattering but, unlike Higgs, it has *opposite* couplings to W and Z
- For generic ab 
  ightarrow cd process in the limit g' 
  ightarrow 0

 $A(s,t,u)\delta^{ab}\delta^{cd} + A(t,s,u)\delta^{ac}\delta^{bd} + A(u,t,s)\delta^{ad}\delta^{bc}$ 

For example  $A_{W^+W^- \rightarrow ZZ} = A(s, t, u)$ ,  $A_{W^+W^+ \rightarrow W^+W^+} = A(t, s, u) + A(u, t, s)$ , etc

• Isospin singlet and quintuplet contribute as Alboteanu et al [0806.4145]

$$A(s,t,u) = \frac{s}{v^2} \left( 1 - c_V^2 \frac{s}{s - m_h^2} \right) + \frac{g_Q^2}{v^2} \left( \frac{s^2}{3(s - m_Q^2)} - \frac{t^2}{2(t - m_Q^2)} - \frac{u^2}{2(u - m_Q^2)} \right)$$

• For  $s \gg m_{h,Q}^2$ 

$$A(s,t,u) \approx rac{s}{v^2} \left(1-c_V^2+rac{5g_Q^2}{6}
ight)$$

Higgs overshoots unitarization, but for  $g_Q^2 = 6/5(c_V^2 - 1)$  quintuplet restores unitary behavior as long as  $m_Q$  is not too large

- Quinituplet can be part of renormalizable Higgs sector provided one allows for higher-than-doublet representations under  $SU(2)_W$
- Minimal model: scalar Φ in (3,3) representation under global SU(2) × SU(2) (complex triplet + real triplet under SU(2)<sub>W</sub>)
- $\bullet~$  Under custodial isospin  $\Phi~$  decomposes as singlet +~ triplet +~ quintuplet

$$\begin{pmatrix} \frac{v}{2\sqrt{2}} + \frac{1}{\sqrt{3}}h - \frac{1}{\sqrt{6}}Q^{0} + \frac{i}{\sqrt{2}}\pi^{0} & -\frac{1}{\sqrt{2}}(Q^{+} + i\pi^{+}) & -Q^{++} \\ -\frac{1}{\sqrt{2}}(Q^{-} + i\pi^{-}) & \frac{v}{2\sqrt{2}} + \frac{1}{\sqrt{3}}h + \sqrt{\frac{2}{3}}Q^{0} & -\frac{1}{\sqrt{2}}(Q^{+} - i\pi^{+}) \\ -Q^{--} & -\frac{1}{\sqrt{2}}(Q^{-} - i\pi^{-}) & \frac{v}{2\sqrt{2}} + \frac{1}{\sqrt{3}}h - \frac{1}{\sqrt{6}}Q^{0} - \frac{i}{\sqrt{2}}\pi^{0} \end{pmatrix}$$

corresponding to  $c_V=\sqrt{8/3}$  and  $g_Q=\sqrt{2}.$ 

- Smaller  $c_V$  can be obtained when  $\Phi$  mixes with EW singlet, or doublet (Georgi,Machacek [(1985)])
- More general Higgs representations under  $SU(2) \times SU(2)$  studied in Low,Lykken [1005.0872]

#### Possible effect on Higgs

• Custodial invariant coupling of Higgs and quintuplet:

$$\mathcal{L}_{hQQ} = -2g_{hQQ}m_Q^2rac{h}{v}\left(|Q^{++}|^2+|Q^{+}|^2+rac{1}{2}(Q^0)^2
ight).$$

Minimal renormalizable model:  $g_{hQQ} = \sqrt{\frac{2}{3}} \frac{m_h^2 + 2m_Q^2}{m_Q^2}$ 

 $\bullet\,$  Shifts effective Higgs coupling to  $\gamma\gamma$  by

$$\delta c_{\gamma} pprox rac{5}{24} g_{hQQ}$$

• Thus, generic prediction of increased Higgs couplings to WW and ZZ, and decreased effective Higgs coupling to photons

- The puzzle of electroweak symmetry breaking is about to be solved
- Hints from the LHC and other experiments consistently point to weakly coupled electroweak symmetry breaking with a light Higgs boson
- Measuring Higgs coupling may soon give us strong hints favoring or disfavoring particular models beyond the Standard Model
- If data clearly points to  $c_V > 1$ , all hands on board to search for 5 more Higgs bosons!
- At least this year is going to be exciting...