



Mind the Gap on IceCube

Cosmic neutrino spectrum and muon anomalous magnetic moment

Toshihiko Ota



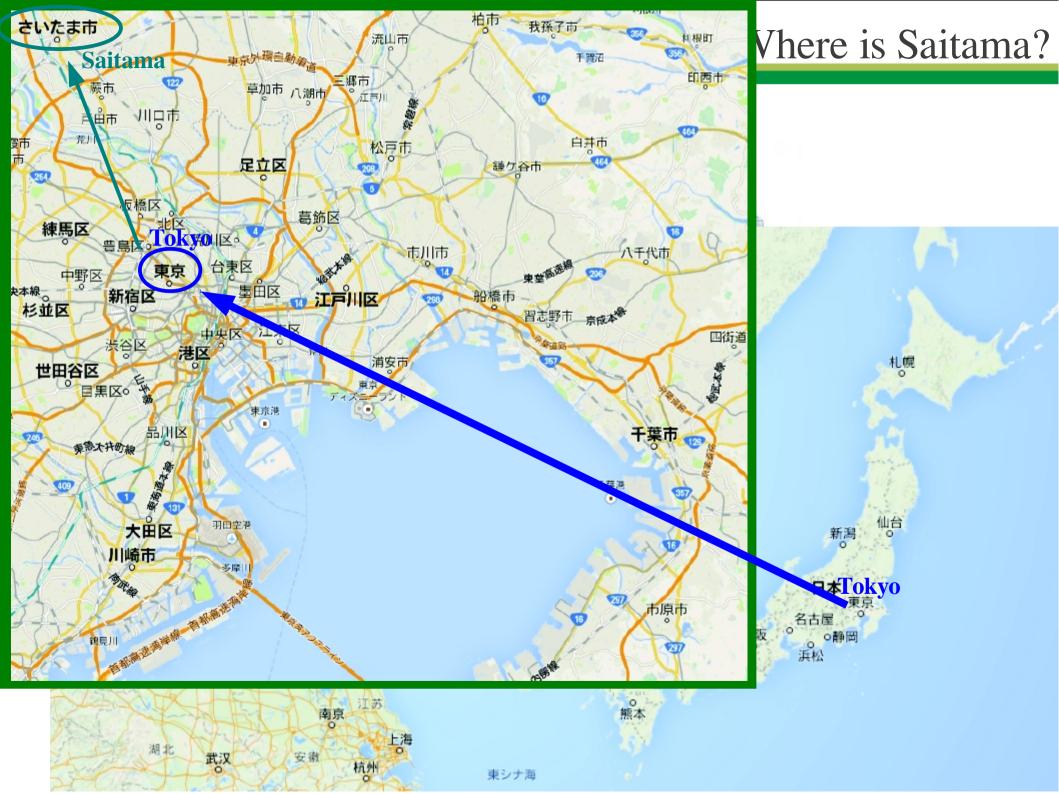
based on

T.Araki, Y.Konishi, F.Kaneko, TO, J.Sato, T.Shimomura

ArXiv.1409.4180v2









LHC Tests of Light Neutralino Dark Matter without Light Sfermions

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Abstract

We address the question how light the lightest MSSM neutralino can be as dark matter candidate in a scenario where all supersymmetric scalar particles are heavy. The hypothesis that the neutralino accounts for the observed dark matter density sets strong requirements on the supersymmetric spectrum, thus providing an handle for collider tests. In particular for a lightest neutralino below 100 GeV the relic density constraint translates into an upper bound on the Higgsino mass parameter μ in case all supersymmetric scalar particles are heavy. One can define a simplified model that highlights only the necessary features of the spectrum and their observable consequences at the LHC. Reinterpreting recent searches at the LHC we derive limits on the mass of the lightest neutralino that, in many cases, prove to be more constraining than dark matter experiments themselves.





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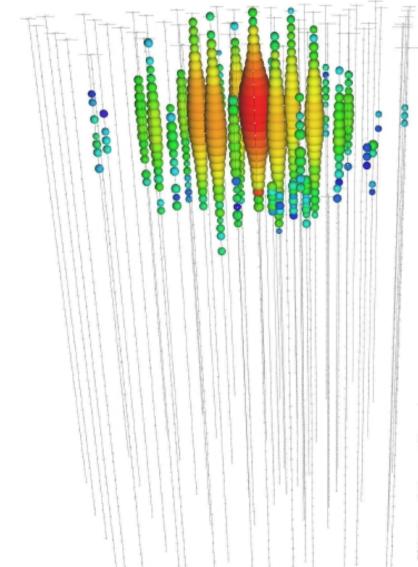
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• PeV cosmic neutrino spectrum IceCube collaboration PRL 113 (2014) 101101

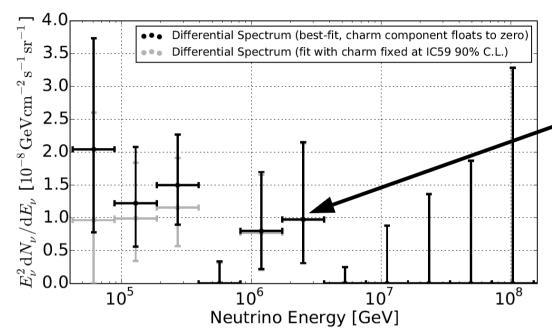






• PeV cosmic neutrino spectrum



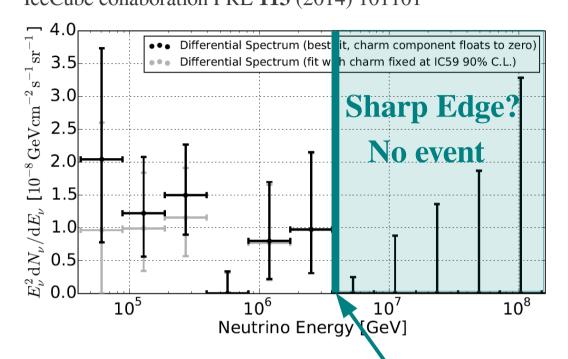


Interpreted to...





 PeV cosmic neutrino spectrum IceCube collaboration PRL 113 (2014) 101101



IceCube Edge?

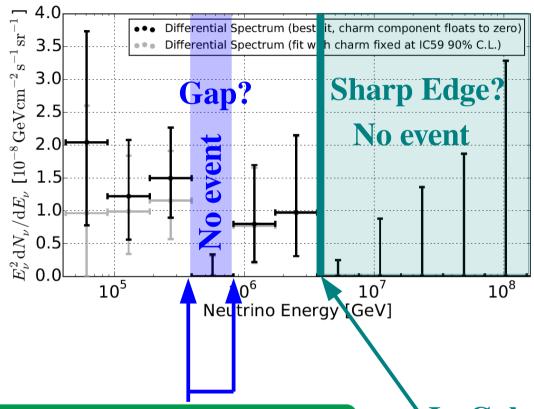
at 3 PeV may be astrophysical origin





• PeV cosmic neutrino spectrum





IceCube Gap

No event at 0.4-1 PeV

IceCube Edge?

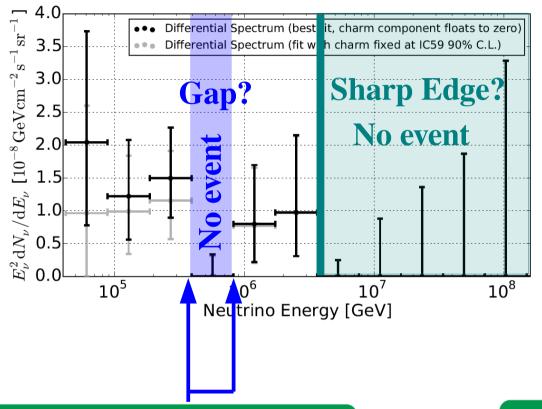
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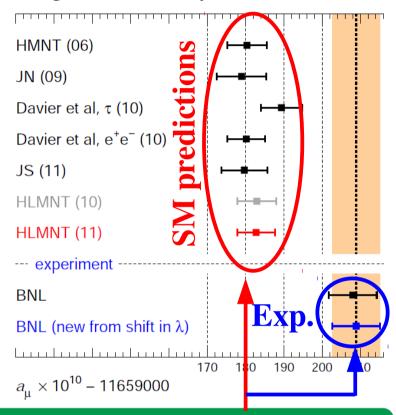
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IceCube collaboration PRL 113 (2014) 101101



●Muon g-2

Hagiwara et al., J.Phys. G38 (2011) 085003



IceCube Gap

No event at 0.4-1 PeV

$$g_{\mu}-2$$
 Gap

$$a_{\mu}^{\mathsf{Exp}} - a_{\mu}^{\mathsf{SM}} = (26.1 \pm 8.0) \cdot 10^{-10} \, (3.3\sigma)$$



New physics at the MeV scale

may explain both the gaps



- IceCube gap
 - Attenuation of cosmic neutrino by secret neutrino interaction
 - Gauged leptonic force $L_{\mu}-L_{\tau}$ as secret interaction
- 2 Muon anomalous magnetic moment
 - Gauged leptonic force as a contribution to g-2
 - Constraints from colliders and neutrino trident process
- A solution to the gaps
 - Reproduction of IceCube gap → distance to the neutrino source
 - → neutrino mass spectrum

If the IceCube Gap is explained by some New Physics (NP)...

埼玉大学 Cosmic neutrino and New Physics

If the IceCube Gap is explained by some New Physics (NP)...

NP at Source: PeV Dark matter decay

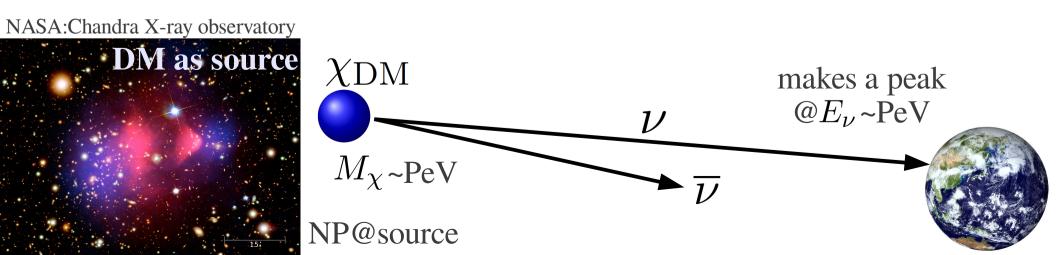
Feldstein Kusenko Matsumoto Yanagida, PRD88 (2013) 015004. Zabala PRD89 (2014) 123514.

Ibarra Tran Weniger Int.J.Mod.Phys. A28 (2013) 1330040.

Esmaili Serpico JCAP **1311** (2013) 054, Esmaili Kang Serpico, 1410.5979.

Ema Jinno Moroi PLB733(2014) 120, JHEP 1410 (2014) 150. Rott Kohri Park 1408.3799.

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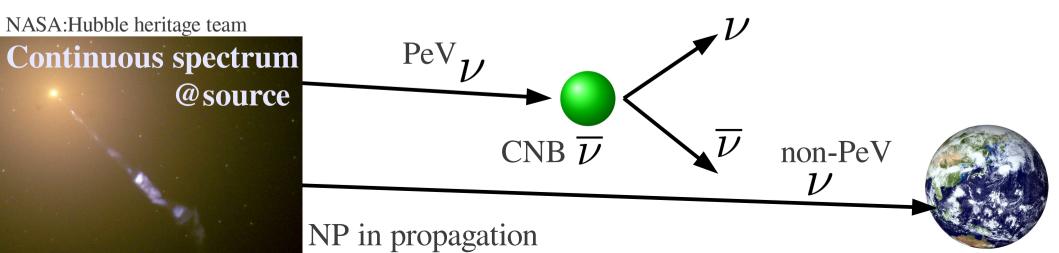
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NP in Propagation: Scattering with CNB with a MeV mediator

As an effective int.: Ng Beacom PR**D90** (2014) 065035, Ioka Murase PTEP **6** (2014) 061E01 With neutrino mass model: Ibe Kaneta PR**D90** (2014) 053011, Blum Hook Murase 1408.3799



埼玉大学 Cosmic neutrino and New Physics



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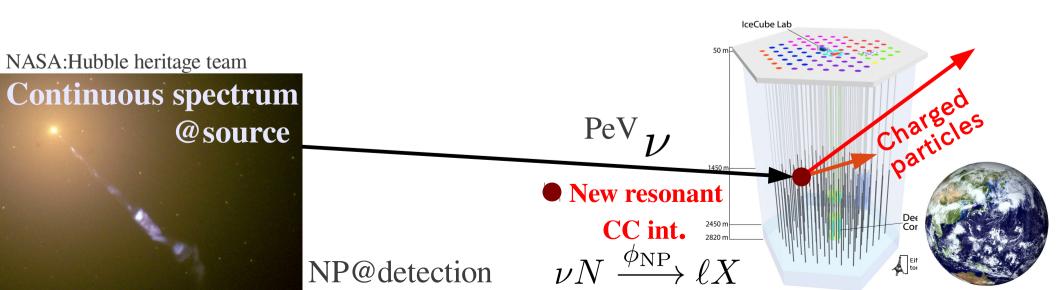
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● NP at Detection: CC int. mediated by a new TeV field Barger Keung PLB727 (2013) 190...









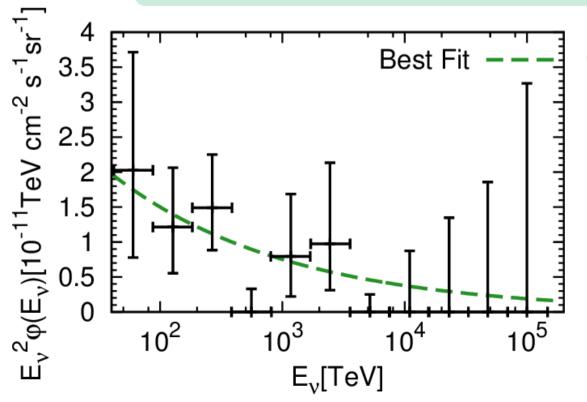
NP in propagation, namely Resonant scattering with CNB

• We set **3 assumptions** for cosmic neutrino sources



NP in propagation, namely Resonant scattering with CNB

- We set **3 assumptions** for cosmic neutrino sources
 - Continuous (power-law) spectrum

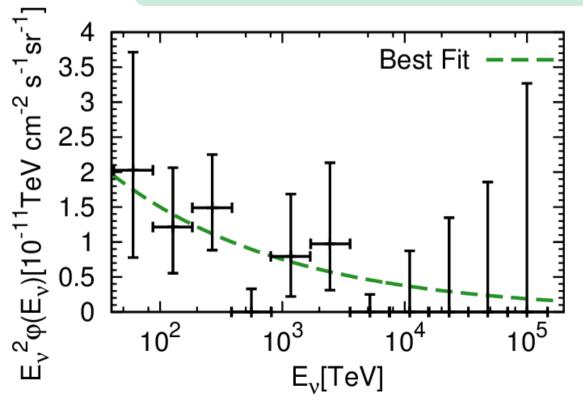


The spectrum shown with the green curve is reproduced, if there is no NP.



NP in propagation, namely Resonant scattering with CNB

- We set **3 assumptions** for cosmic neutrino sources
 - Continuous (power-law) spectrum
 - 2 Flavour ratio ~1:1:1 after leaving sources

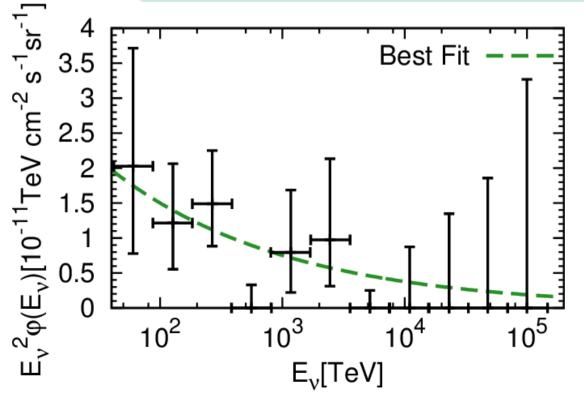


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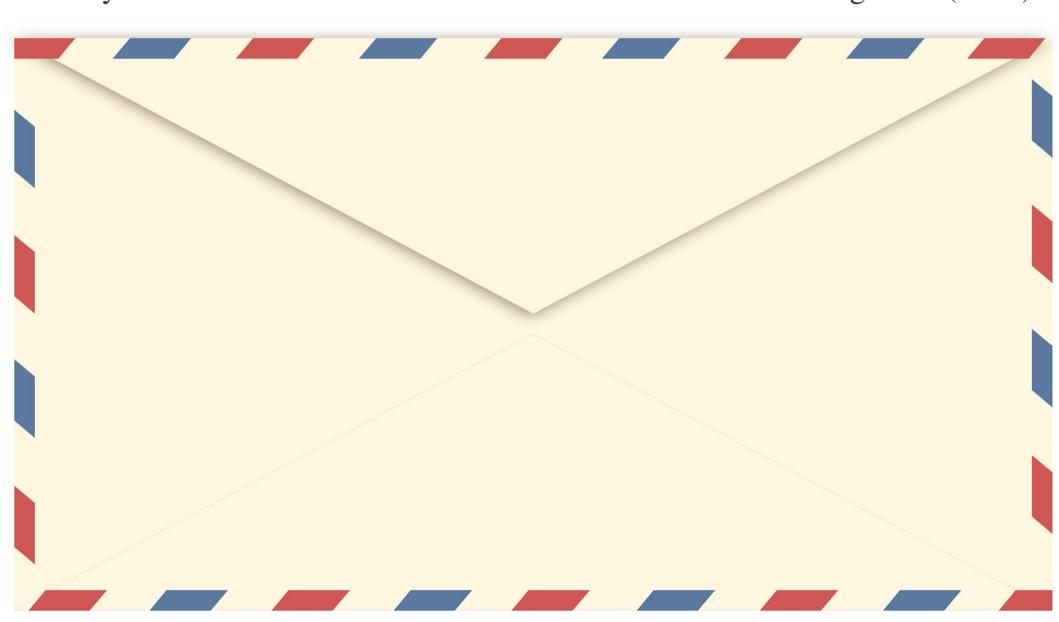
NP in propagation, namely Resonant scattering with CNB

- We set **3 assumptions** for cosmic neutrino sources
 - Continuous (power-law) spectrum
 - 2 Flavour ratio ~1:1:1 after leaving sources
 - 3 Sources distribute around a particular redshift z_{source}



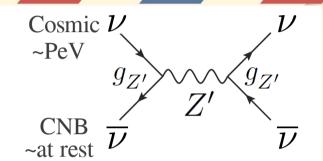
- The spectrum shown with the green curve is reproduced, if there is no NP.
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- 3 for simplicity.
 - \rightarrow z-dependence of source distribution e.g., The star-formation rate has a peak at $z = 1 \sim 2$.





Resonance condition

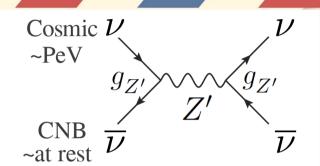
$$s \simeq 2E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNB}}} \stackrel{!}{=} M_{Z'}^2$$



Why CNB? $\rightarrow n_{\rm CNB} \gg n_{\rm Baryon}$ $n_{\rm CNB} = 56.8 \, [/{\rm cm}^3] \, {\rm for \ each \ dof}$

Resonance condition

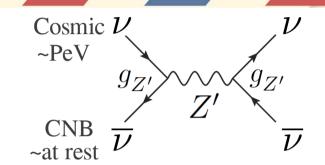
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NP @MeV scale

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$$\lambda \simeq 1/n_{\rm CNB}\sigma_{\rm @Res.}$$



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IceCube Gap requires

$$M_{Z'} \sim \text{MeV}, \quad \sigma_{\text{@Res.}} \gtrsim 10^{-30} \text{ [cm}^2\text{]}.$$

Let us calculate the cross-section in a particular model...

• Gauged $U(1)L_{\mu}-L_{\tau}$ force as **a benchmark model**

Charge assignments
$$Y(L_{\mu}) = +1, Y(L_{\tau}) = -1, Y(\mu_R) = +1, Y(\tau_R) = -1, Y(\text{others}) = 0.$$

$$\mathcal{L}_{L_{\mu}-L_{\tau}} = g_{Z'} \overline{L}_{\mu} \gamma^{\rho} L_{\mu} Z'_{\rho} - g_{Z'} \overline{L}_{\tau} \gamma^{\rho} L_{\tau} Z'_{\rho} + g_{Z'} \overline{\mu_R} \gamma^{\rho} \mu_R Z'_{\rho} - g_{Z'} \overline{\tau_R} \gamma^{\rho} \tau_R Z'_{\rho}$$

• Gauged $U(1)L_{\mu}-L_{\tau}$ force as **a benchmark model**

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$$+ g_{Z'} \overline{\mu_{R}} \gamma^{\rho} \mu_{R} Z'_{\rho} - g_{Z'} \overline{\tau_{R}} \gamma^{\rho} \tau_{R} Z'_{\rho}$$

$$= g_{ij} \overline{\nu}_{i} \gamma^{\rho} P_{L} \nu_{j} Z'_{\rho} + g_{Z'} \mathsf{diag}(0, 1, -1)_{\alpha\beta} \overline{\ell}_{\alpha} \gamma^{\rho} \ell_{\beta} Z'_{\rho}$$

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Neutrino secret int.

Coupling in mass eigenbasis

$$g_{ij} = g_{Z'}(U_{\mathrm{PMNS}}^{\dagger})_{i\alpha} \operatorname{diag}(0, 1, -1)_{\alpha\beta}(U_{\mathrm{PMNS}})_{\beta j}$$

Constrained! but... Contribute to muon g-2

We discuss it in Sec. 2

^{*} Cosmic neutrino is produced as a flavour eigenstate= a coherent sum of mass eigenstates. But the coherence is lost in its travel.

• Gauged $U(1)L_{\mu} - L_{\tau}$ force as **a benchmark model**

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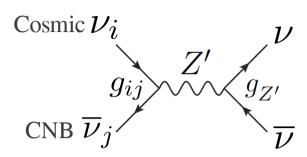
* Cosmic neutrino is produced as a flavour eigenstate= a coherent sum of mass eigenstates. But the coherence is lost in its travel.

- Motivated from...
 - (almost) Maximal mixing Choubey Rodejohann Eur. Phys. J C40 (2005) 259
 - Gauge anomaly free Foot Mod. Phys. A6 (1991) 527, He et al., PRD43 (1990) R22
- In this talk, we do not go into the details of the spontaneous breaking of the $L_{\mu}-L_{\tau}$ sym.

Model parameters

 g_{Z^\prime} and M_{Z^\prime}

$$\sigma(\nu_i \overline{\nu}_j \to \nu \overline{\nu}) = \frac{|g_{ij}|^2 g_{Z'}^2}{6\pi} \frac{s}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2} \qquad \sup_{\text{CNB}} \overline{\nu}_j \xrightarrow{Z'} \overline{\nu}_{\overline{\nu}}$$



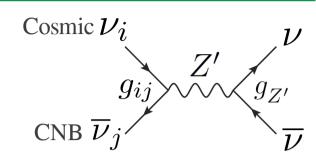
Decay rate

$$\Gamma_{Z'} = \frac{g_{Z'}^2 M_{Z'}}{12\pi}$$

• Cross-section@Resonance

$$\Gamma_{Z'} = \frac{g_{Z'}^2 M_{Z'}}{12\pi} \qquad \qquad \sigma_{\text{@Res.}} = \frac{4\pi |g_{ij}|^2}{M_{Z'}^2} \delta \left(1 - \frac{M_{Z'}^2}{s}\right)$$

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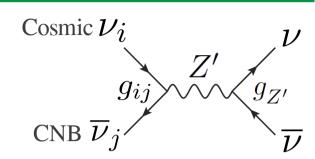
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For IceCube Gap

$$\Gamma_{Z'} = \frac{g_{Z'}^2 M_{Z'}}{12\pi} \qquad \sigma_{\text{@Res.}} = \frac{4\pi |g_{ij}|^2}{M_{Z'}^2} \delta \left(1 - \frac{M_{Z'}^2}{s}\right) \stackrel{\text{For IceCube Gap}}{=} 10^{-30} \text{ [cm}^2\text{]}$$

$$M_{Z'} \sim \text{MeV}$$

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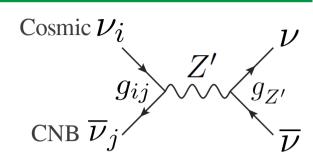
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$$M_{Z'} \sim \text{MeV} \xrightarrow{} g_{Z'} \simeq 10^{-3}$$

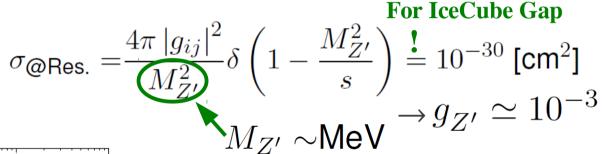
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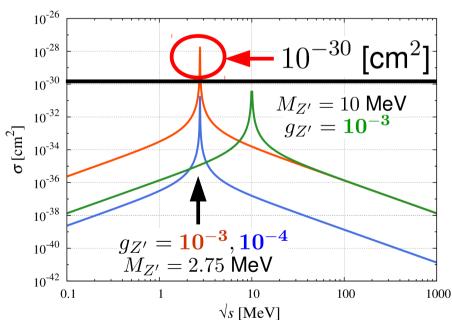


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Cross-section@Resonance





IceCube Gap requires

 $M_{Z'} \sim \text{MeV}, \quad g_{Z'} \gtrsim 10^{-4}.$

- The width might be **too narrow** for the **IceCube Gap (0.4-1PeV)**.
- We can ask the help to m_{ν} and $z \to \mathbf{Sec.}$

Before going into the details of the cosmic neutrino spectrum, let's check muon g-2.



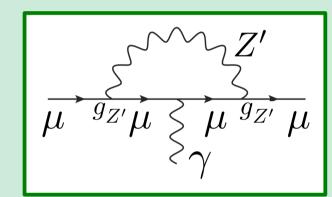
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•
$$M_{Z'} \gg m_{\mu} \rightarrow \Delta a_{\mu}^{Z'} = \frac{g_{Z'}^2}{12\pi^2} \frac{m_{\mu}^2}{M_{Z'}^2}$$
• $M_{Z'} \ll m_{\mu} \rightarrow \Delta a_{\mu}^{Z'} = \frac{g_{Z'}^2}{8\pi^2}$

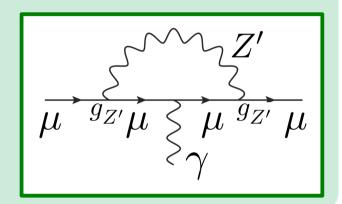
$$\bullet M_{Z'} \ll m_{\mu} \to \Delta a_{\mu}^{Z'} = \frac{g_{Z'}^2}{8\pi^2}$$



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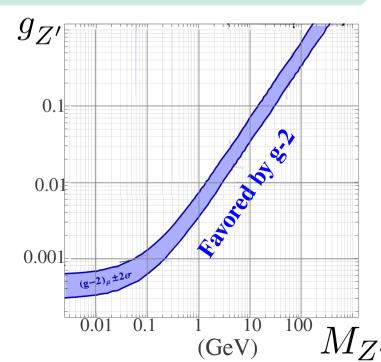
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$g_{\mu}-2$ Gap

$$a_{\mu}^{\mathsf{Exp}} - a_{\mu}^{\mathsf{SM}} = (26.1 \pm 8.0) \cdot 10^{-10} \, (3.3\sigma)$$

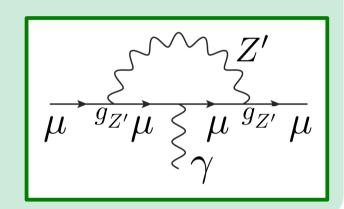
$$\rightarrow$$
 We need $\Delta a_{\mu}^{\rm NP} \simeq (20\text{-}30) \cdot 10^{-10}$



$$\mathscr{L}_{L_{\mu}-L_{\tau}} = g_{ij}\overline{\nu}_{i}\gamma^{\rho}\mathrm{P}_{L}\nu_{j}Z_{\rho}' + g_{Z'}\mathrm{diag}(0,1,-1)_{\alpha\beta}\overline{\ell}_{\alpha}\gamma^{\rho}\ell_{\beta}Z_{\rho}'$$

•
$$M_{Z'} \gg m_{\mu} \rightarrow \Delta a_{\mu}^{Z'} = \frac{(g_{Z'}^2)}{12\pi^2} \frac{m_{\mu}^2}{M_{Z'}^2}$$

$$\bullet M_{Z'} \ll m_{\mu} \rightarrow \Delta a_{\mu}^{Z'} = \frac{g_{Z'}^2}{8\pi^2}$$



$\overline{|g_{\mu}|}-2$ Gap

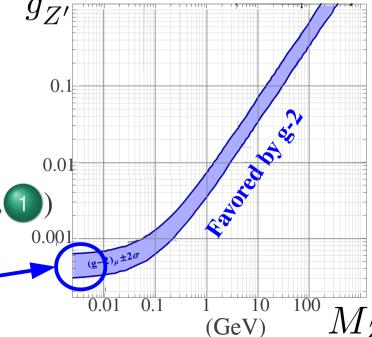
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• Let me remind (back-of-the envelope calc. in **Sec.** 1)

IceCube Gap requires

$$M_{Z'} \sim \text{MeV}, \quad g_{Z'} \gtrsim 10^{-4}.$$

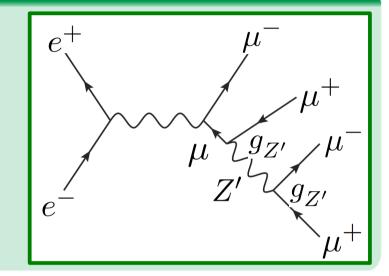


Collider bounds Harigaya et al., JHEP 1403 (2014) 105.

• Process: $e^+e^- \to 4\mu$ $PP(P\bar{P}) \to 4\mu/2\mu2\tau$

only constrain relatively heavy Z'

ightarrow LEP, LHC: $g_{Z'} \lesssim 0.1$ at $M_{Z'} \simeq 100$ GeV

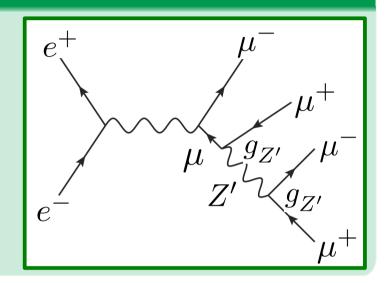


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Rare meson decays Lessa and Peres, PRD75 (2007) 094001

• Process: $\pi^+/K^+ \to \mu^+ \nu_\mu Z'$

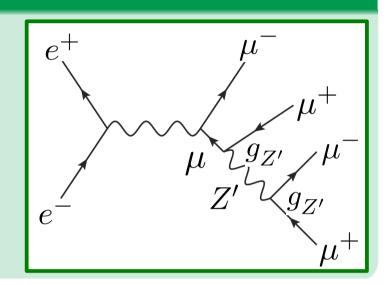
Bound from Kaon decay $\to g_{Z^\prime} \lesssim 0.01 {\rm at}\, M_{Z^\prime} {\sim} {\rm MeV}$

Collider bounds Harigaya et al., JHEP 1403 (2014) 105.

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Bound from Kaon decay $\to g_{Z^\prime} \lesssim 0.01 \, \mathrm{at} \, M_{Z^\prime} \text{~MeV}$

• The most relevant bound from lab. experiments is

Neutrino trident process in neutrino-nucleon scattering

Altmannshofer Gori Pospelov Yavin, PRL 113 (2014) 091801

■ Bounds from CMB, BBN, and also from SN1987A → References in Ng Beacom

埼玉大学 Constraints: Neutrino Trident Process



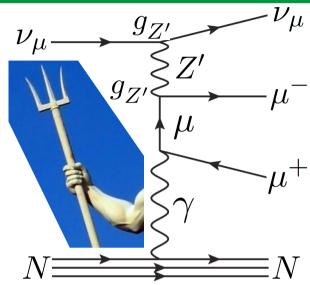
Neutrino trident process

in neutrino-nucleon scattering events

• Available data reported by CCFR in 1991!

37 events (± 12.4)

CCFR collaboration, PRL **66** (1991) 3117 **excavated recently** (only cited 18 times)



Altmannshofer et al., PRL 113 (2014) 091801

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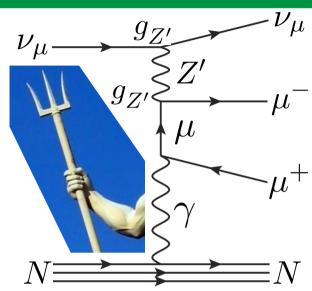
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ullet Expected **SM contribution** mediated by Z and W

45.3 events (± 2.3)

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Altmannshofer et al., PRL 113 (2014) 091801

Neutrino trident process

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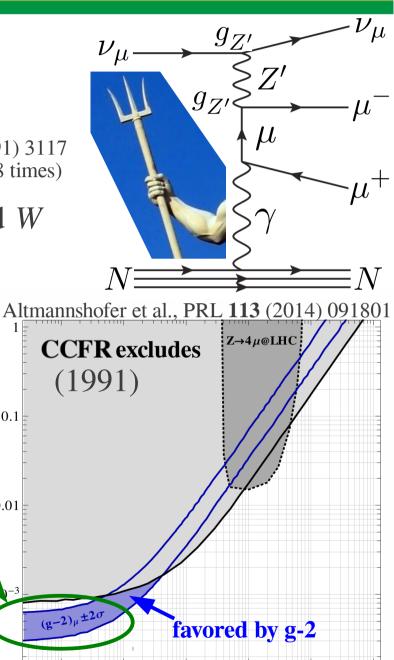
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Consistent \rightarrow constrains $g_{Z'}$ and $M_{Z'}$

$g_{\mu}-2$ favored - Trident excl.

 $M_{Z'} \lesssim 100$ MeV, $g_{Z'} \simeq \text{several} \cdot 10^{-4}$.



 $M_{Z'}({\rm GeV})$

 10^{2}

0.1

0.01

0.1

 $g_{Z^\prime{}_{0.01}}$

*The trident process must be recorded on the hard disks of

the near detectors in modern oscillation experiments. They should be opened!

Neutrino trident process

in neutrino-nucleon scattering events

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 - **37 events (±12.4)** CCFR collaboration, PRL **66** (1991) 3117 excavated recently (only cited 18 times)
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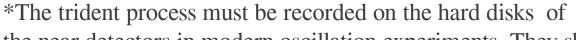
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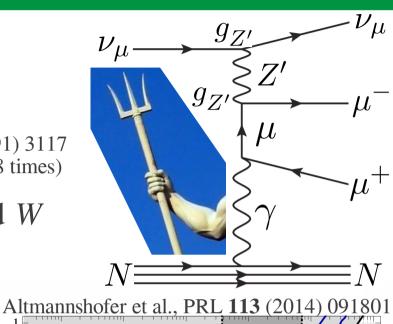
Coincide!

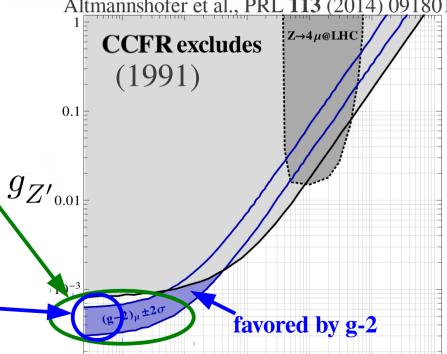
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 $M_{Z'}({\rm GeV})$

0.01

0.1

 10^{2}

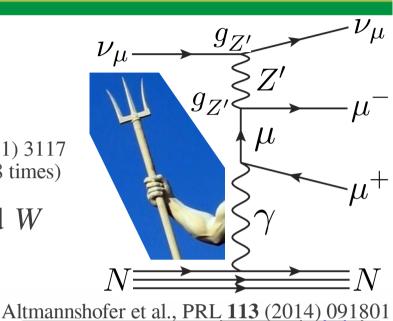
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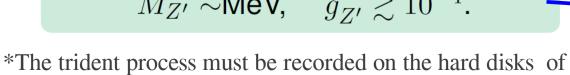
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favored by g-2

0.1

0.01 the near detectors in modern oscillation experiments. They should be opened!

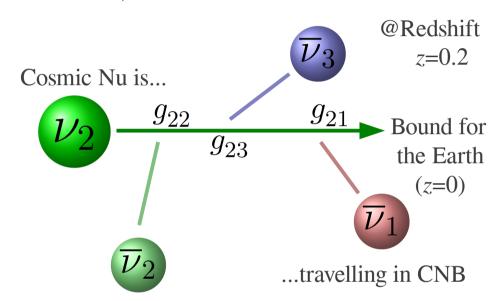
 $M_{Z'}({\rm GeV})$

 $\Delta a_{\mu}^{Z'} = 31.7 \cdot 10^{-10}$

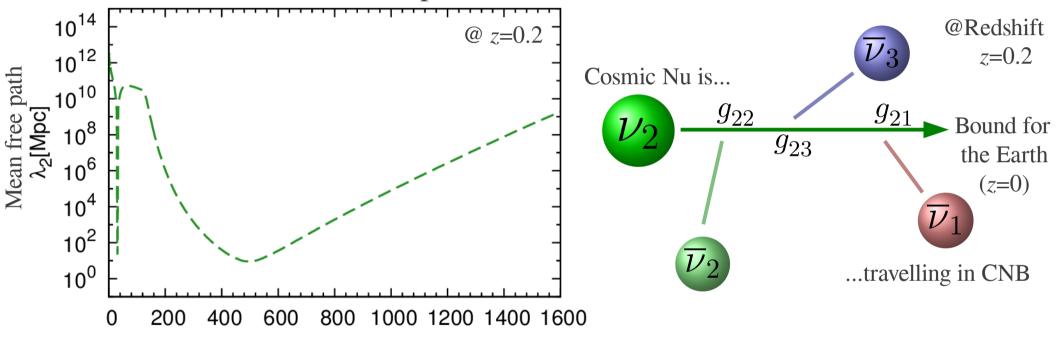


- IceCube gap
 - Attenuation of cosmic neutrino by secret neutrino interaction
 - Gauged leptonic force $L_{\mu}-L_{\tau}$ as secret interaction
- 2 Muon anomalous magnetic moment
 - Gauged leptonic force as a contribution to g-2
 - Constraints from colliders and neutrino trident process
- A solution to the gaps
 - Reproduction of IceCube gap → distance to the neutrino source
 - → neutrino mass spectrum

- $M_{Z'} = 2.75 \text{ MeV}, g_{Z'} = 5.0 \cdot 10^{-4} \rightarrow \text{Favored by } \textbf{g-2} \text{ and allowed by } \textbf{Trident}$
- We fix $m_{\nu_{\text{lightest}}} = 3.0 \cdot 10^{-3}$ eV and take IH $m_{\nu_3} \ll m_{\nu_1} < m_{\nu_2}$
- Let us calculate the mean free path (for 2^{nd} neutrino) at z=0.2.



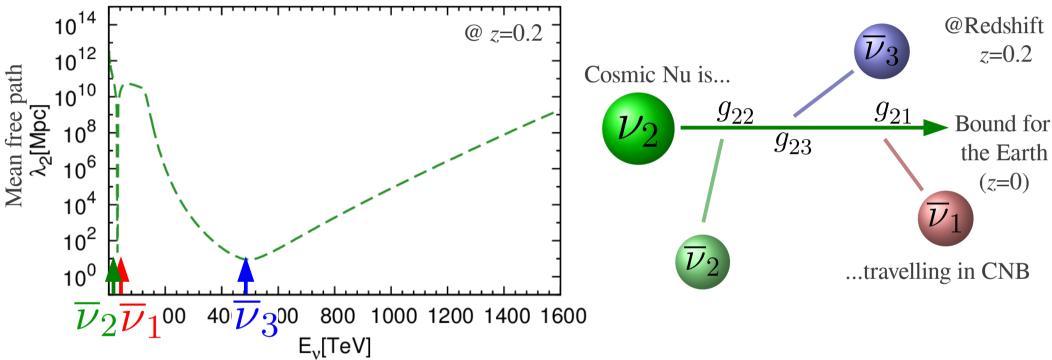
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Neutrino energy observed at IceCube

E_v[TeV]

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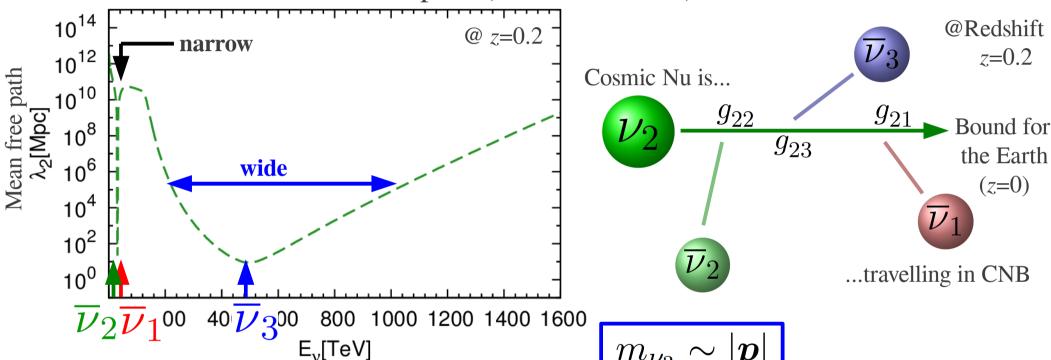
Neutrino energy observed at IceCube

Resonant condition w. CNB distribution

$$s \simeq 2\underline{E_{\nu_{i=2}}(1+z)}_{\text{Neutrino energy }@_{\mathcal{Z}}} \left[\sqrt{|\boldsymbol{p}|^2 + m_{\nu_{j}}^2} - |\boldsymbol{p}| \cos \theta \right] \stackrel{!}{=} M_{Z'}^2$$

 $|m{p}|$: CNB momentum follows Fermi-Dirac dist. $\lesssim (1+z)T_{\nu0} \sim 2.0\cdot 10^{-4}$ [eV] @z=0.2

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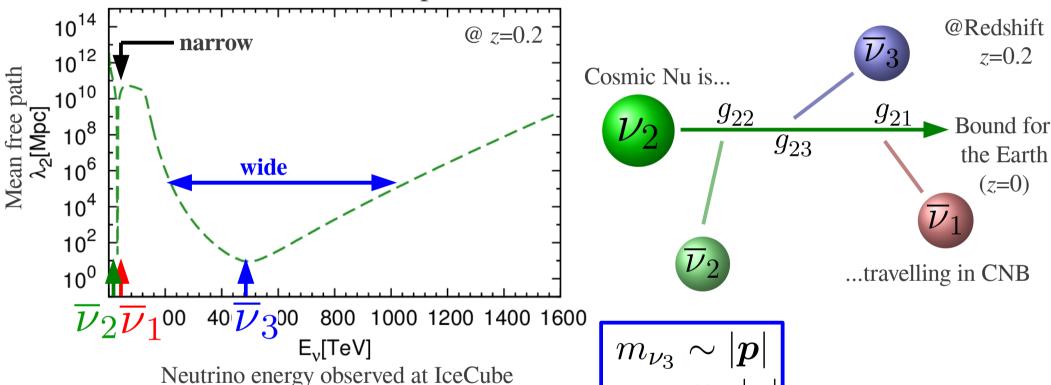
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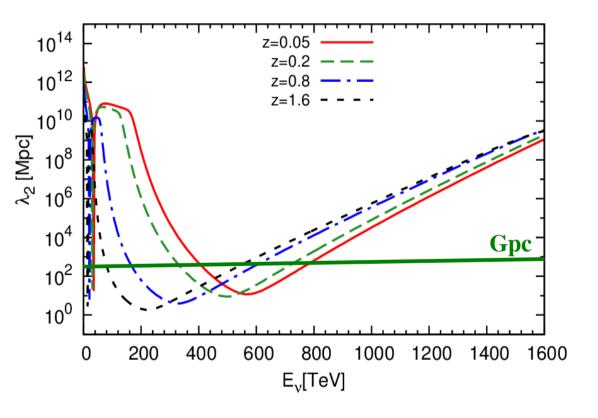
z shifts resonant E

Small $mNu \rightarrow$ wide width Large $z \rightarrow$ wide width

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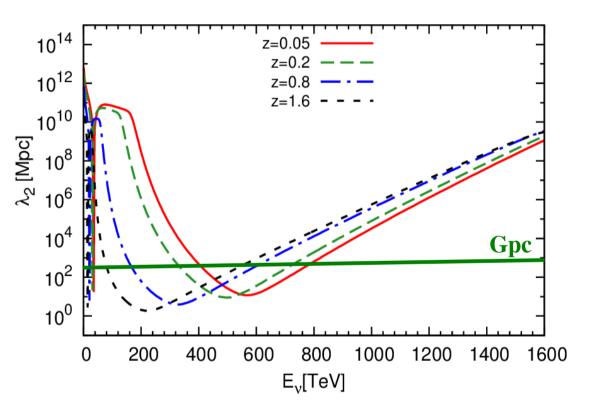
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- Let us have a closer look at *z* dependence of MFP

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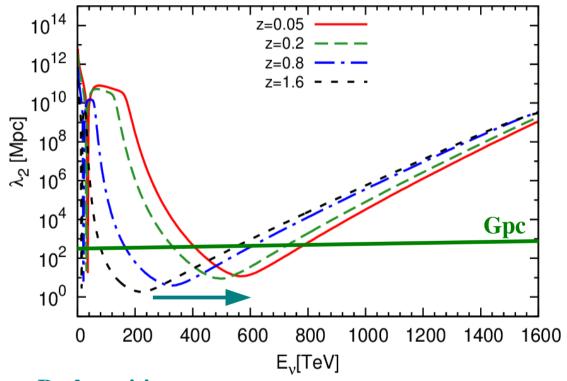
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- Cosmic neutrinos travel from $z_{\text{source to }z=0\text{ (Earth)}}$
- The resonance energy shifts along the travel path.

To keep the width of the gap appropriate, the source should not be so distant from the Earth.

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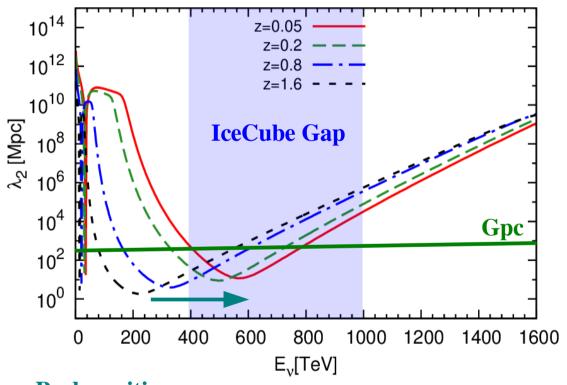
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Peak position moves

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Peak position moves

$$z_{\text{source}} \longrightarrow z = 0$$

• We set $z_{\text{source}}=0.2$ so that the IceCube Gap is reproduced.

In reality, sources of cosmic neutrinos are distributed following some distribution function (e.g., the star formation rate)

■ Mean free path → Spectrum

Following the approximation adopted in Ibe Kaneta PRD...

$$\varphi_i(E_{\nu}) = \varphi_i^{\text{original}}(E_{\nu}) \exp\left[-\int_0^{z_{\text{source}}} \frac{1}{\lambda_i(E_{\nu})}\right]$$

 $\frac{1}{\lambda_{i}(E_{\nu})} \frac{dL}{dz} dz = \frac{10^{14}}{10^{12}} \frac{\lambda_{10}^{12}}{10^{10}} \frac{\lambda_{10}^{12}}{\lambda_{3}^{10}} \frac{\lambda_{10}^{12}}{10^{10}} \frac{\lambda_{10}^{12}}{10^{$

200

The resulting gap does not depends on the initial flavour composition.

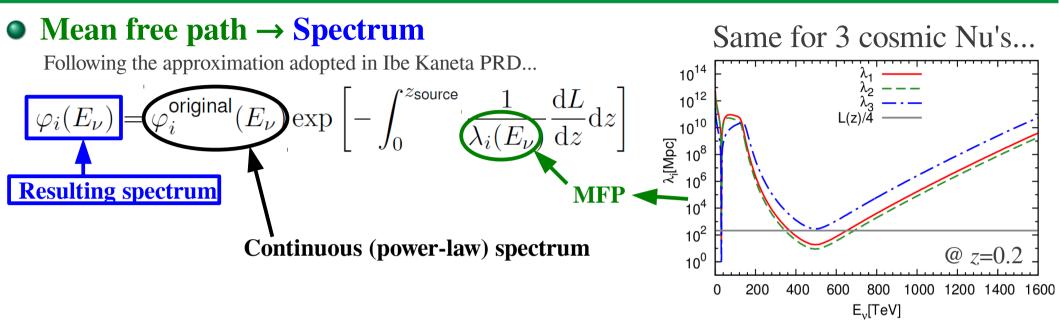
E,[TeV]

Same for 3 cosmic Nu's...

1000 1200 1400 1600







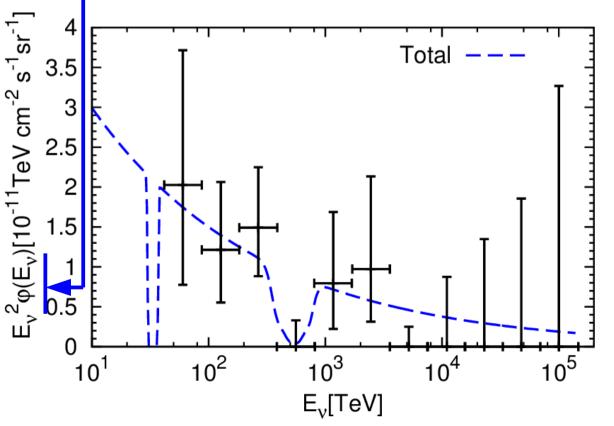
\bullet Mean free path \rightarrow Spectrum

Following the approximation adopted in Ibe Kaneta PRD...

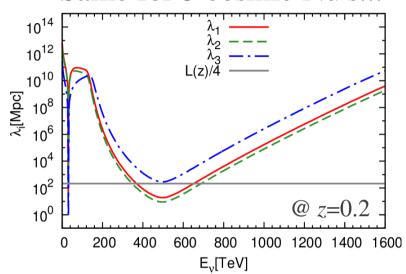
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Resulting spectrum

$$\varphi(E_{\nu}) = \sum_{i} \varphi_{i}(E_{\nu})$$
 assuming flavour universal $\varphi_{i}^{\text{original}}(E_{\nu})$



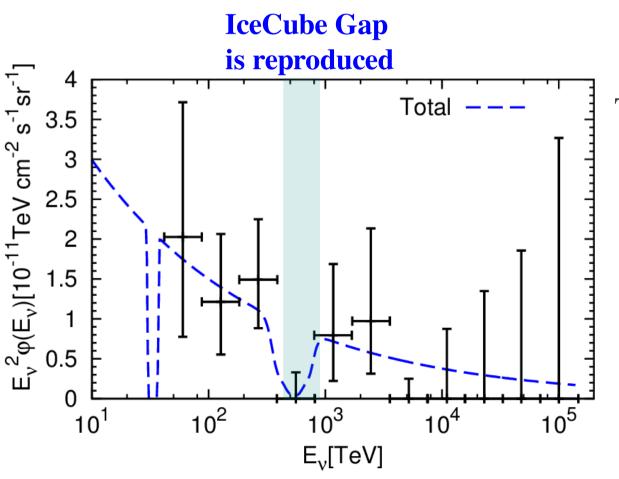
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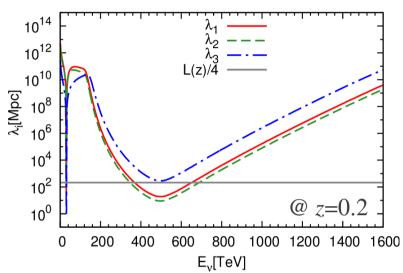
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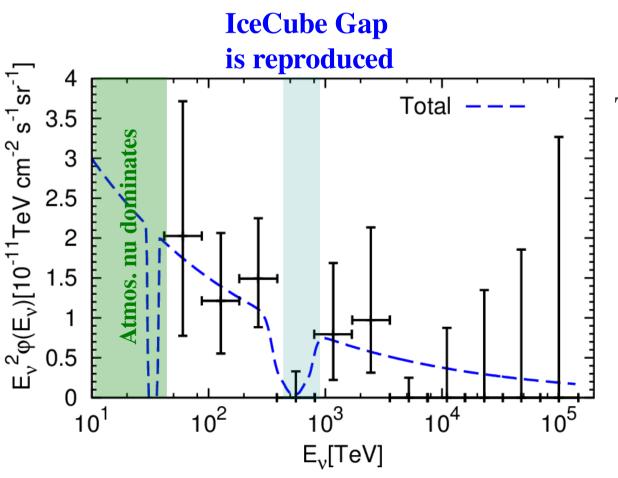
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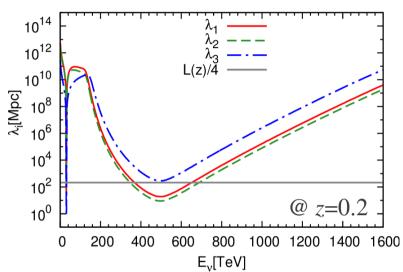
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\bullet Mean free path \rightarrow Spectrum

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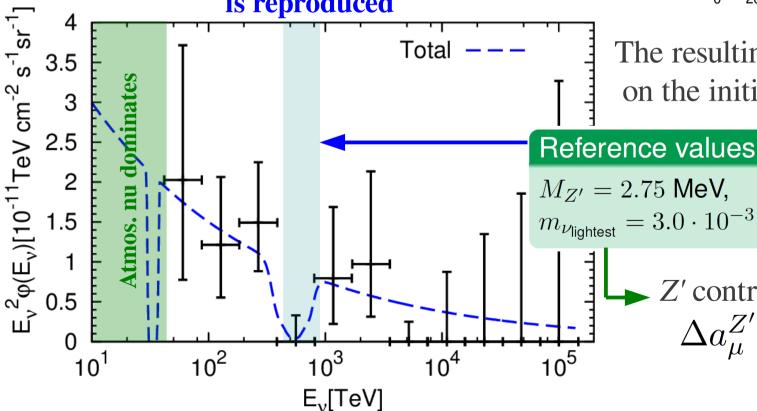
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10¹² λ_i[Mpc] 10⁴ 10^{2} @ z=0.2E,[TeV]

Same for 3 cosmic Nu's...



The resulting gap does not depends on the initial flavour composition.



 $M_{Z'} = 2.75 \text{ MeV}, \quad g_{Z'} = 5.0 \cdot 10^{-4},$

$$m_{\nu_{
m lightest}} = 3.0 \cdot 10^{-3} \ {
m eV}$$
 (IH) and $z_{
m source} = 0.2$.

Z' contribution to muon g-2 $\Delta a_{\mu}^{Z'} = 31.7 \cdot 10^{-10}$ g-2 Gap is filled



Summary and future prospects

We dig the cosmic neutrino spectrum to make a gap and swing around the surplus soil to fill the gap in muon g-2.

Reference values

$$\begin{split} M_{Z'} &= 2.75 \text{ MeV}, \quad g_{Z'} = 5.0 \cdot 10^{-4}, \\ m_{\nu_{\text{lightest}}} &= 3.0 \cdot 10^{-3} \text{ eV (IH) and } z_{\text{source}} = 0.2. \end{split}$$

- IceCube Gap is reproduced.
- $\Delta a_{\mu}^{Z'} = 31.7 \cdot 10^{-10}$

n g-2.

g-2 gap

This tool is called as "U(1) leptonic force Lmu-Ltau"

But we did not...

- ...take into account distribution of neutrino source.
- ...also take into account secondary neutrino effect.
- ...discussed details of the model.

This small try shows that the idea works!

More precise, detailed, and sophisticated study may be worth to be done.

Back up slides



Back up slides

Distribution: star formation rate z=1-2.

P gamma (photohadronic)→ IceCube edge

An astrophysical explanation:
P gamma (photohadronic) (Enu<Gap) + PP (Enu<Gap)

Z vs distance (pc and light year)

Lambda on the Mnu-vs-Enu plane @z=0 (with various mZ')

Normal hierarchy

Discussion on Spontaneous breaking of Lmu-Ltau → some reference...? Check Gori et al...or Harigaya et al...Spires



