The Enigma of the Highest Energy Particles in the Universe

- 1. Observations
- 2. Anisotropies and Mass Composition
- 3. Secondary Neutrinos
- 4. Tests of Lorentz Symmetry
- 5. Cosmic Rays and the Structured Universe



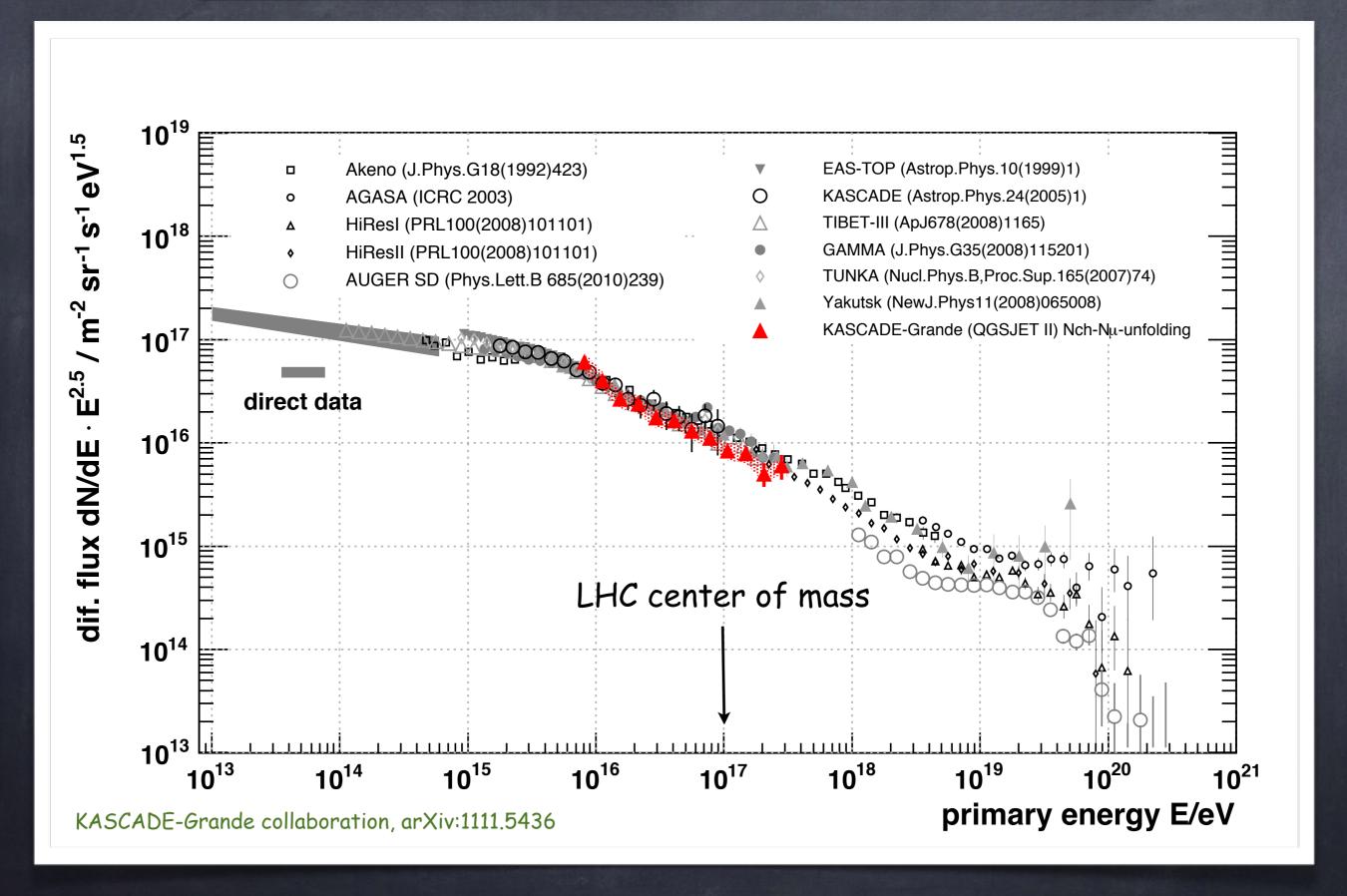
Bundesministerium für Bildung und Forschung



Günter Sigl

II. Institut theoretische Physik, Universität Hamburg

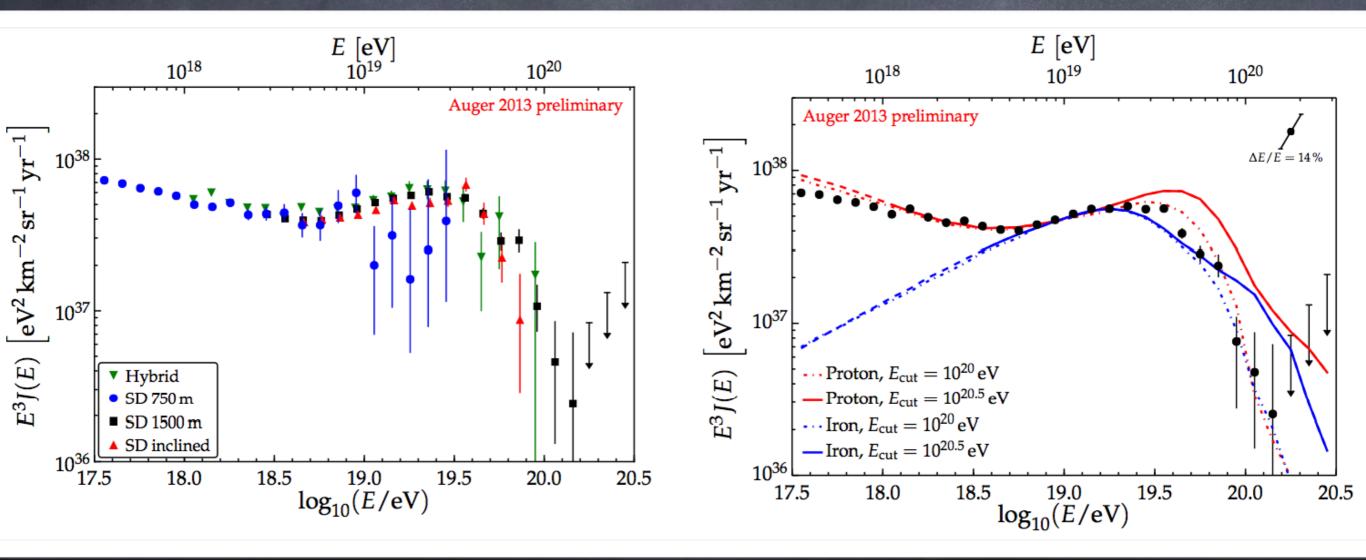
The All Particle Cosmic Ray Spectrum

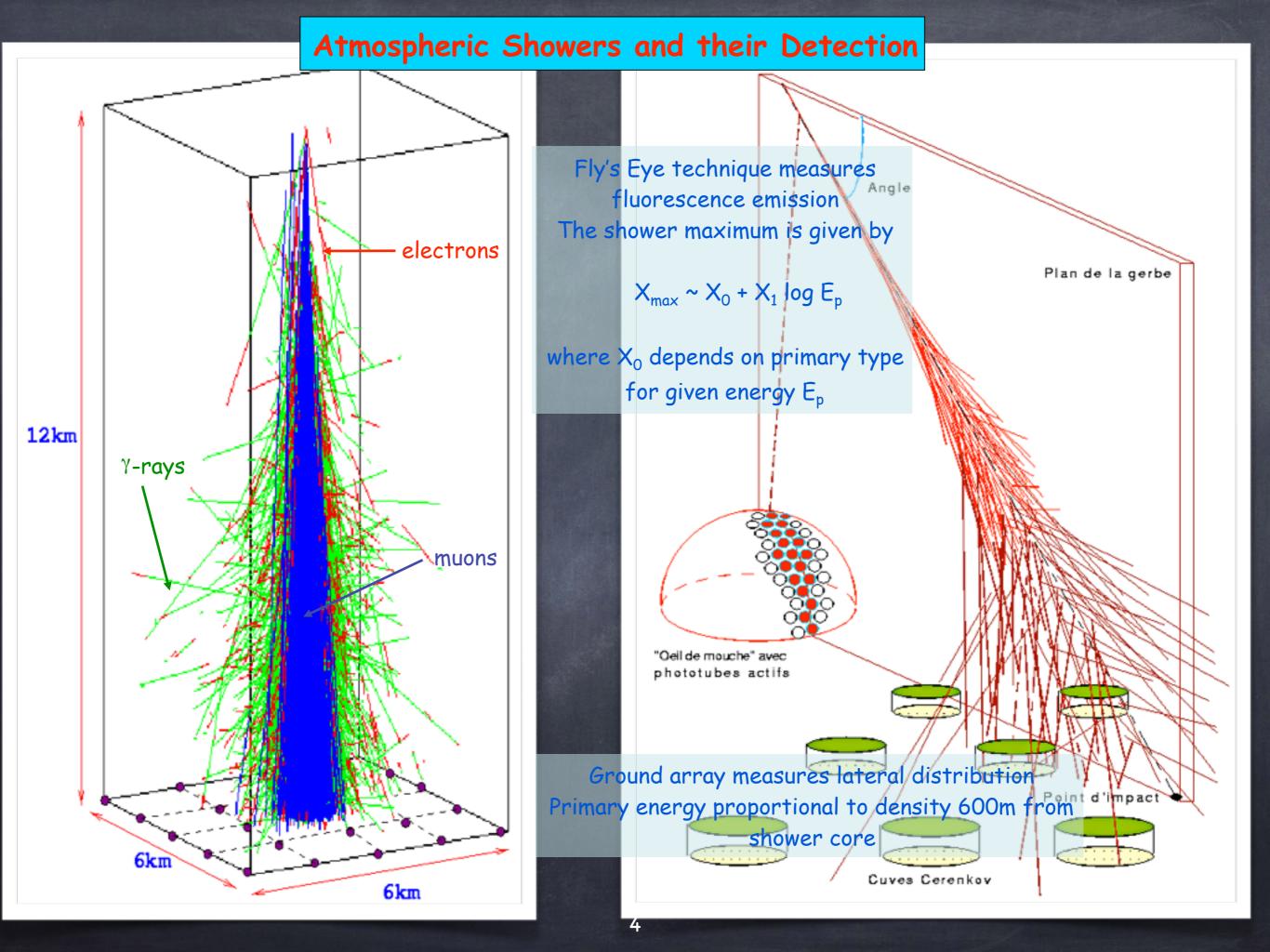


Pierre Auger Spectra

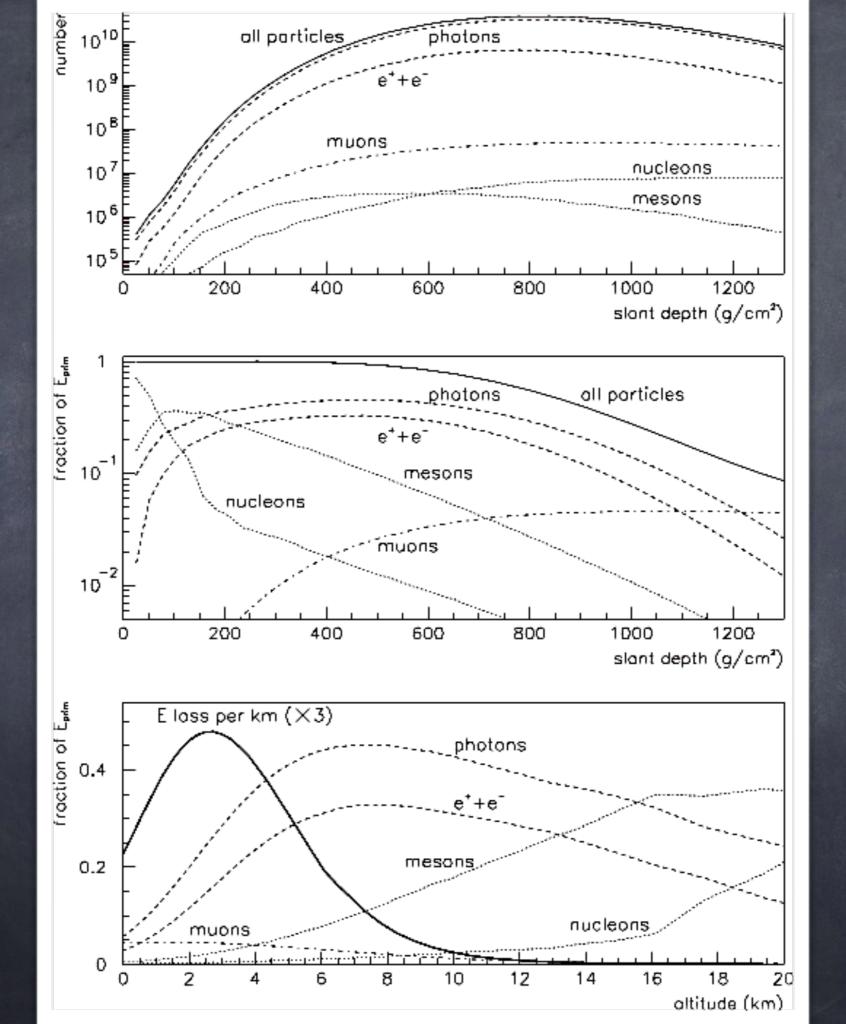
Auger exposure = 31440 km² sr yr, 8259 events above 10¹⁹ eV up to December 2012

Pierre Auger Collaboration, PRL 101, 061101 (2008) and Phys.Lett.B 685 (2010) 239 and ICRC 2013, arXiv:1307.5059, higlight talk Letessier-Selvon

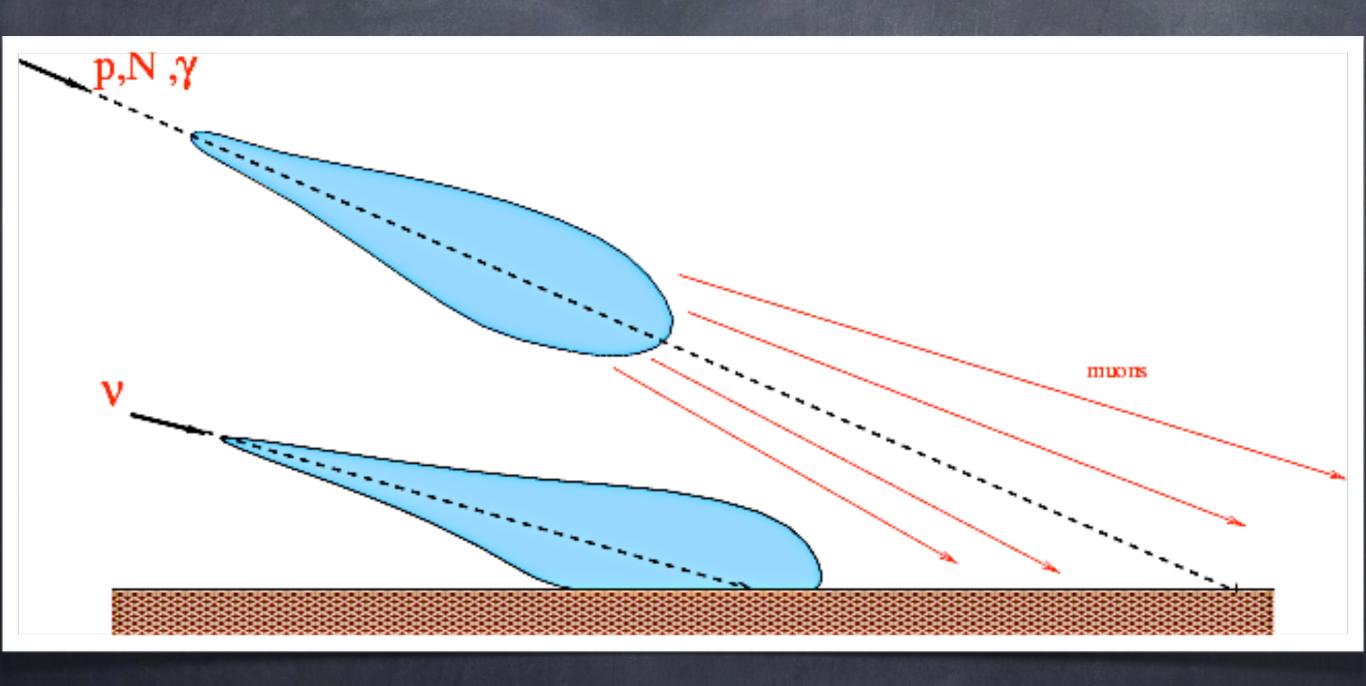


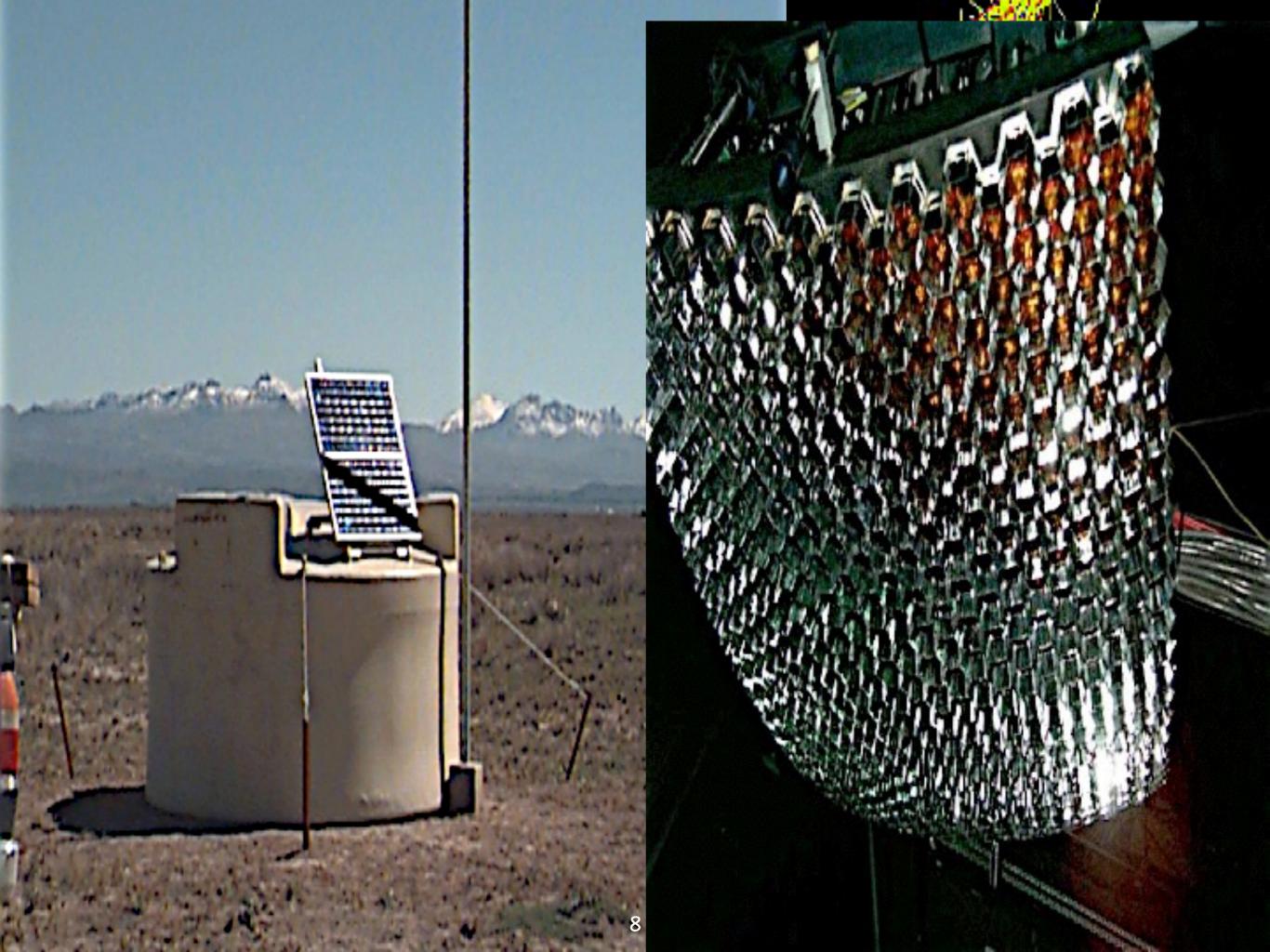


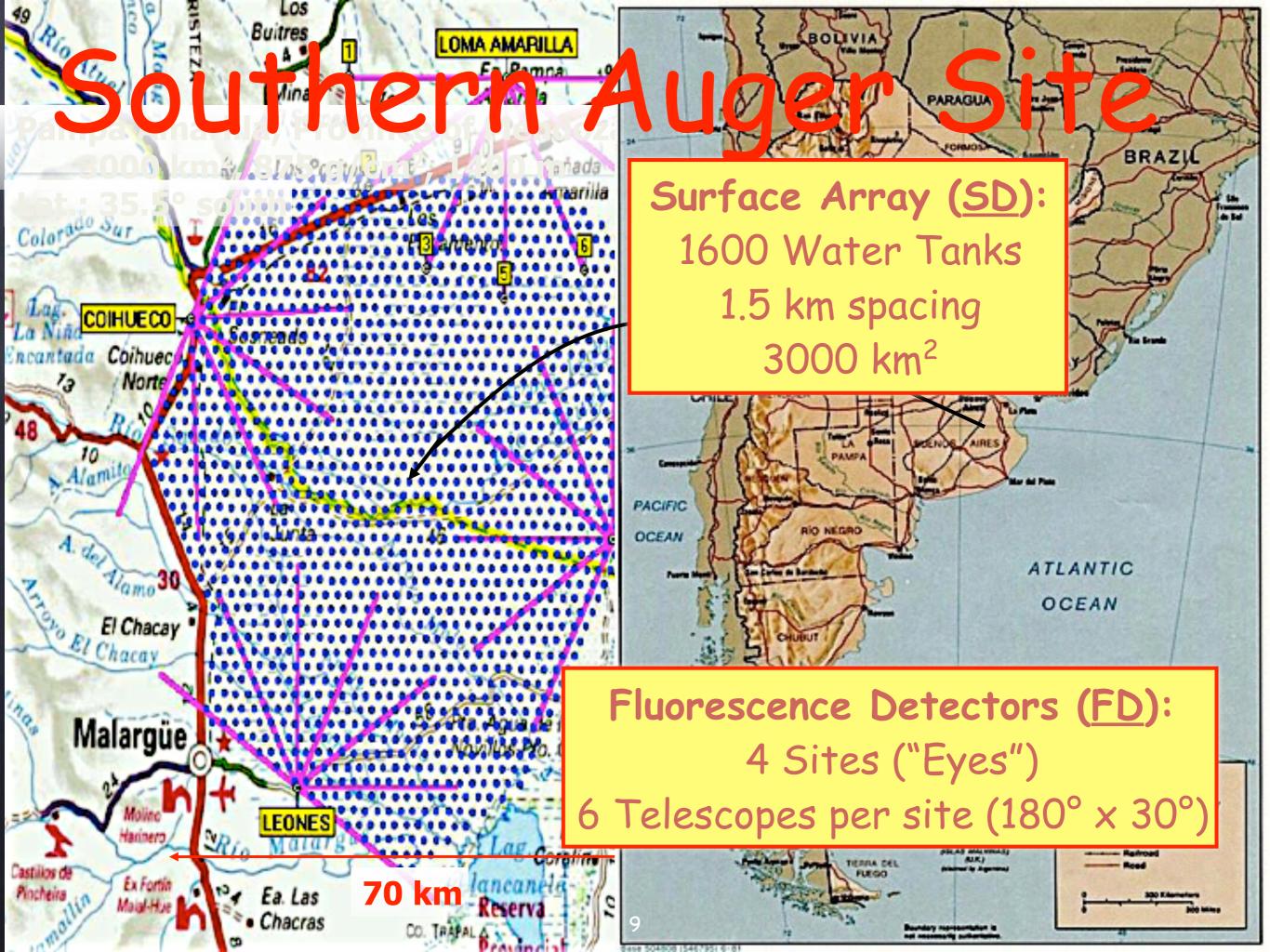
hadronic cascade nuclei, nucleo nadronic ms π,κ mesons long decay muons tail photons e.m. tail electrons positrons electromagnetic cascade



Cosmic ray versus neutrino induced air showers







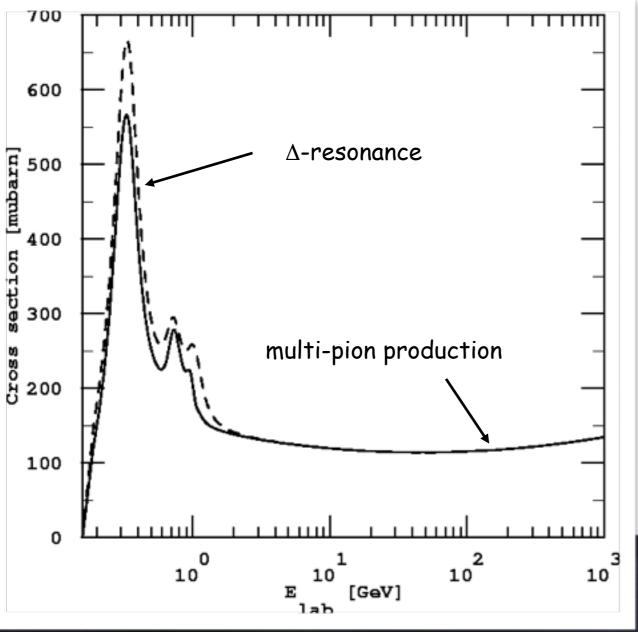
The Ultra-High Energy Cosmic Ray Mystery consists of (at least) Four Interrelated Challenges

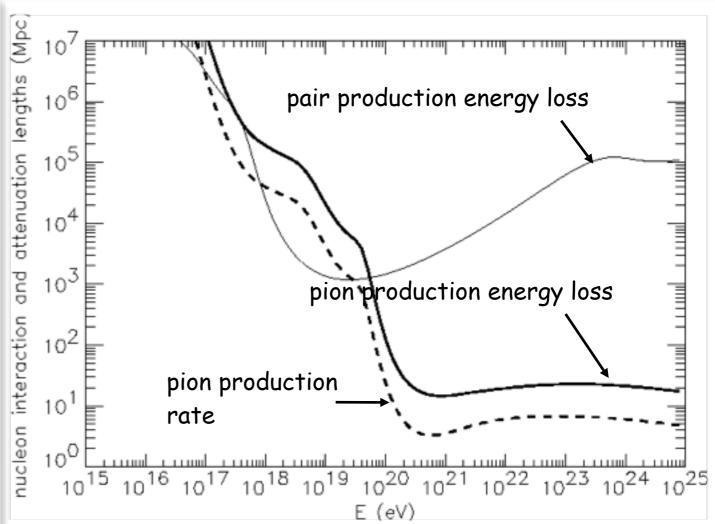
- 1.) electromagnetically or strongly interacting particles above 10^{20} eV loose energy within less than about 50 Mpc.
 - 2.) in most conventional scenarios exceptionally powerful acceleration sources within that distance are needed.
- 3.) The observed distribution does not yet reveal unambiguously the sources, although there are hints of correlations with local large scale structure
 - 4.) The observed mass composition may become heavy toward highest energies, but no completely clear picture yet between experiments and air shower models

The Greisen-Zatsepin-Kuzmin (GZK) effect

Nucleons can produce pions on the cosmic microwave background

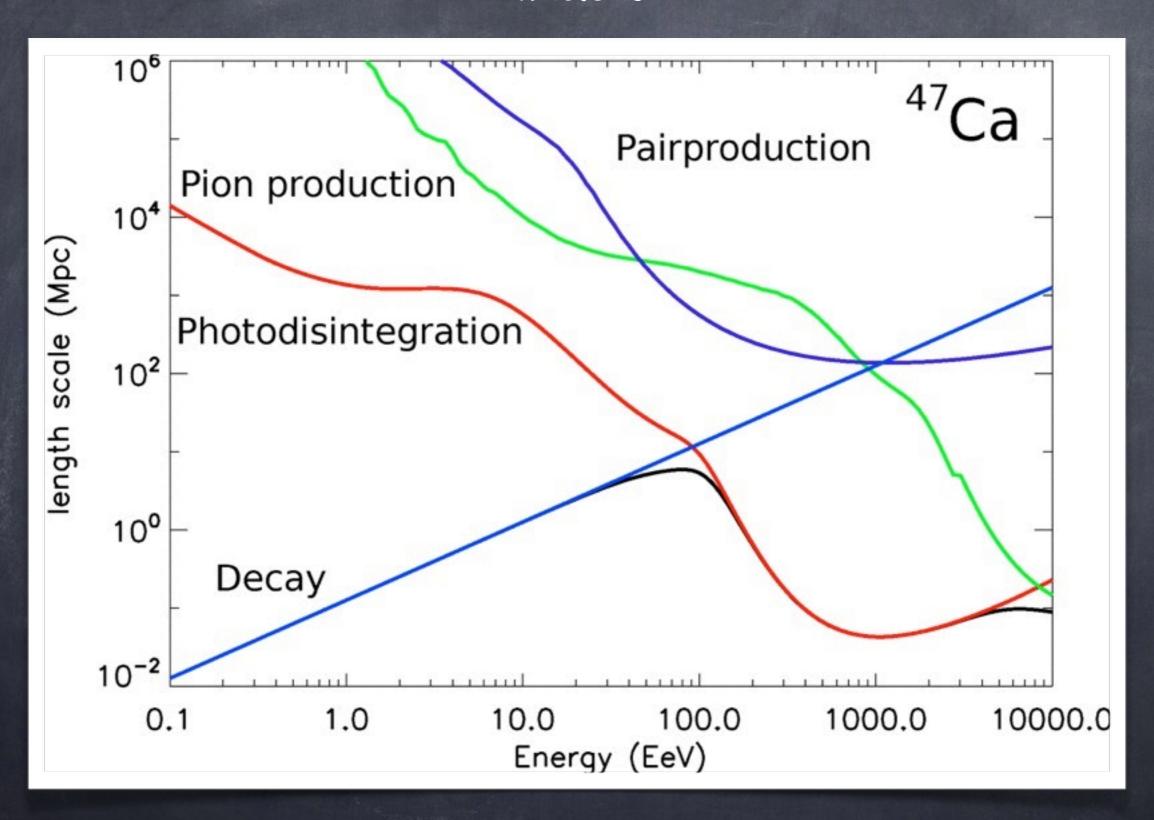




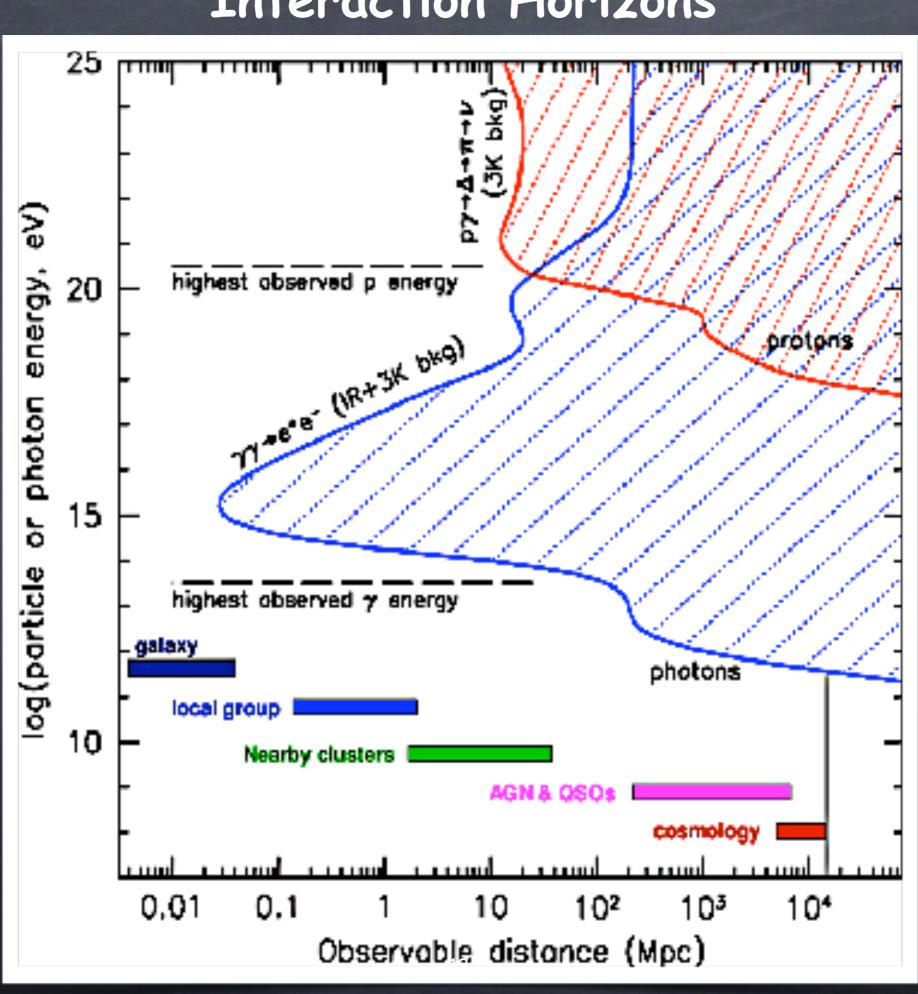


sources must be in cosmological backyard Only Lorentz symmetry breaking at Γ >10¹¹ could avoid this conclusion.

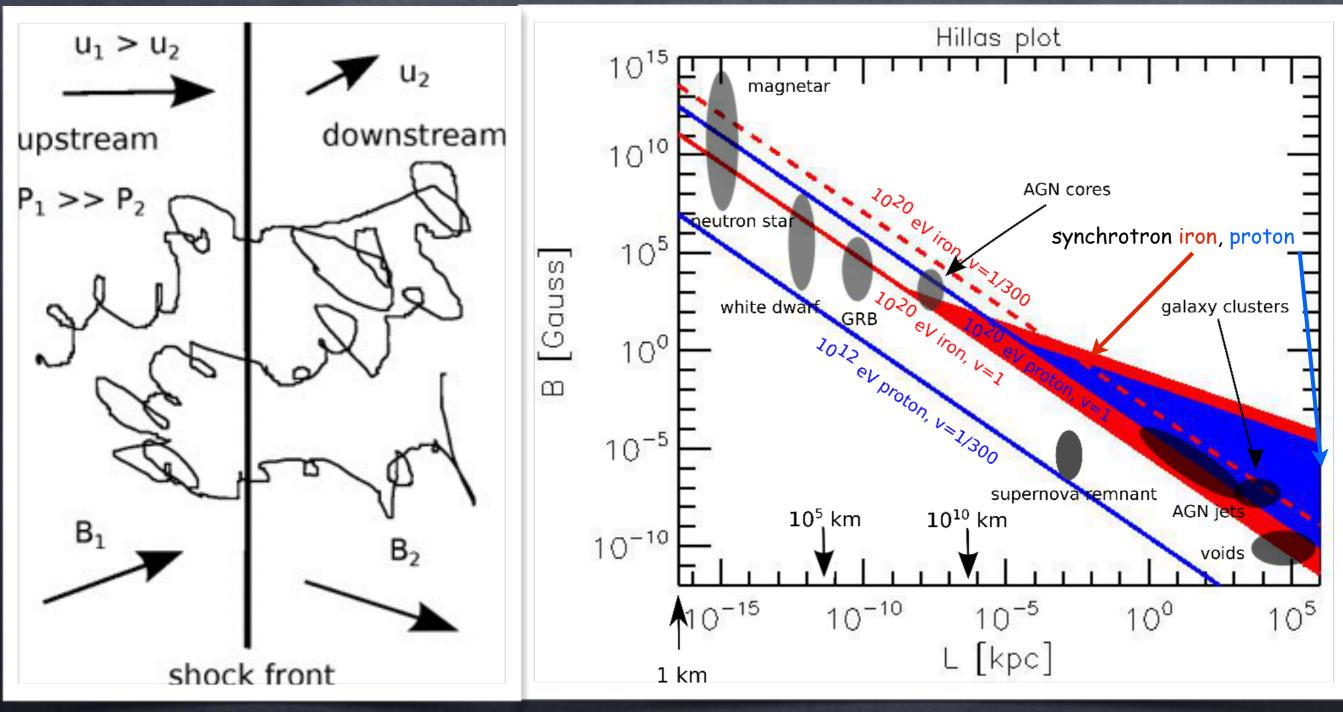
Length scales for relevant processes of a typical heavy nucleus



Interaction Horizons



1st Order Fermi Shock Acceleration



Fractional energy gain per shock crossing $\sim u_1$ - u_2 on a time scale r_L/u_2 .

Together with downstream losses this leads to a spectrum E^{-q} with q > 2 typically. Confinement, gyroradius < shock size, and energy loss times define maximal energy

Some general Requirements for Sources

Accelerating particles of charge eZ to energy E_{max} requires induction $\epsilon > E_{max}/eZ$. With $Z_0 \sim 100\Omega$ the vacuum impedance, this requires dissipation of minimum power of

$$L_{\min} \sim \frac{\epsilon^2}{Z_0} \simeq 10^{45} Z^{-2} \left(\frac{E_{\max}}{10^{20} \text{ eV}}\right)^2 \text{ erg s}^{-1}$$

This "Poynting" luminosity can also be obtained from $L_{min} \sim (BR)^2$ where BR is given by the "Hillas criterium":

$$BR > 3 \times 10^{17} \,\Gamma^{-1} \left(\frac{E_{\text{max}}/Z}{10^{20} \,\text{eV}}\right) \,\text{Gauss cm}$$

where Γ is a possible beaming factor.

If most of this goes into electromagnetic channel, only AGNs and maybe gamma-ray bursts could be consistent with this.

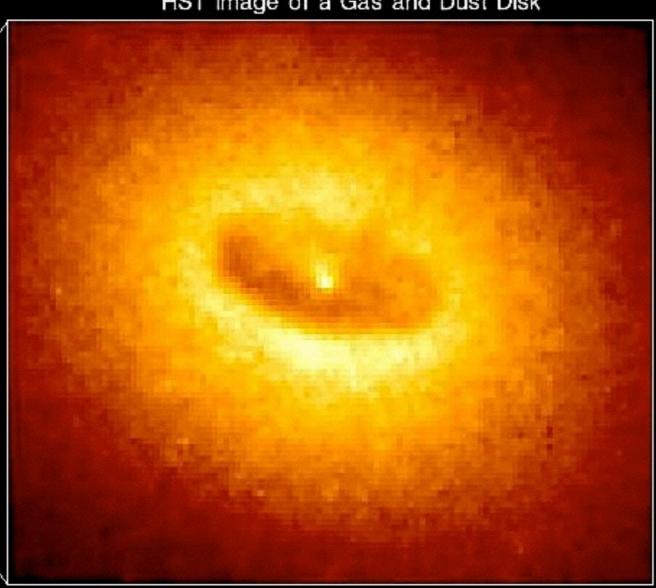
Core of Galaxy NGC 4261

Hubble Space Telescope

Wide Field / Planetary Camera

Ground-Based Optical/Radio Image

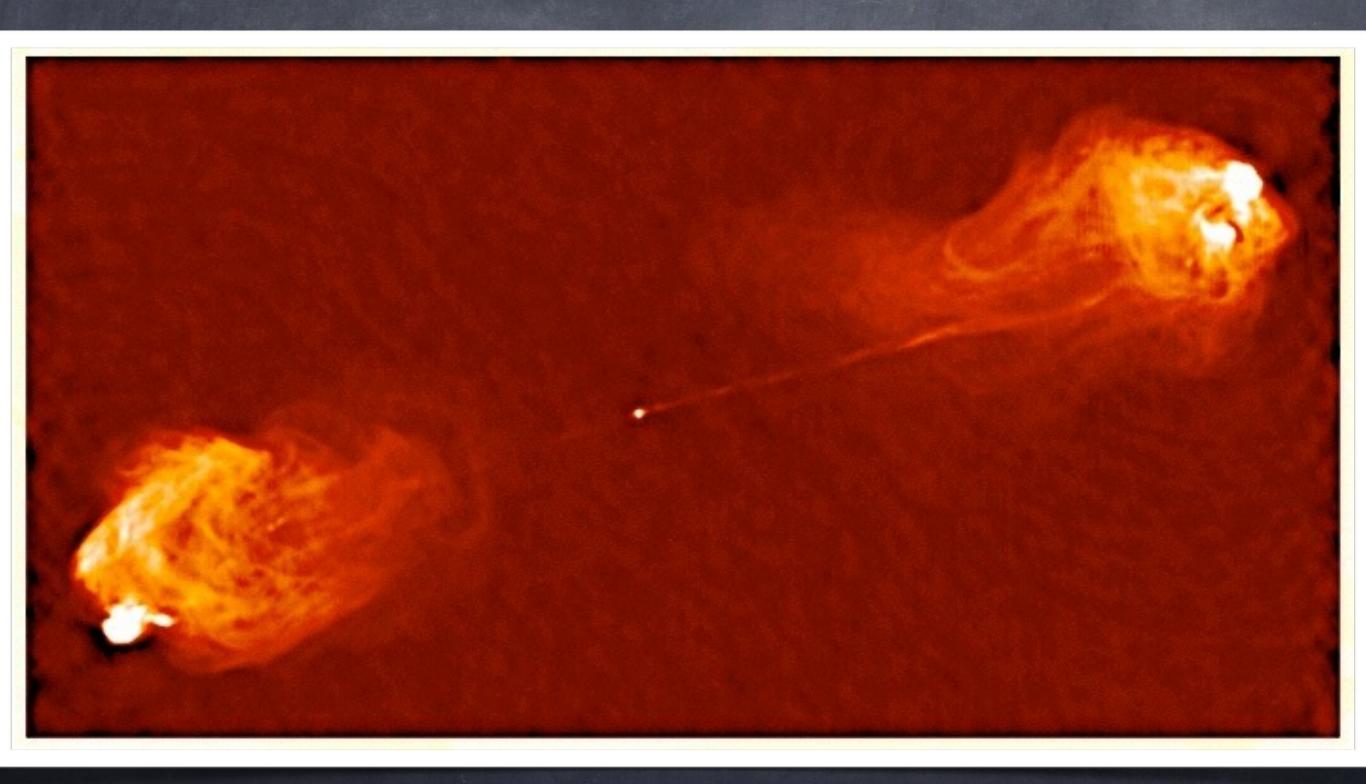
HST Image of a Gas and Dust Disk



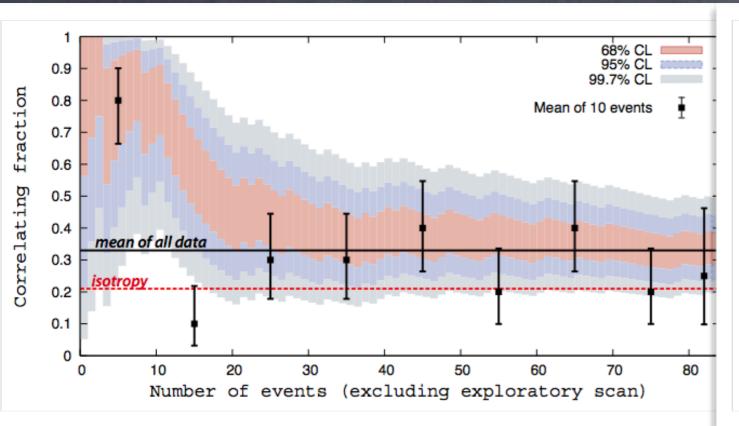
380 Arc Seconds 88,000 LIGHTYEARS

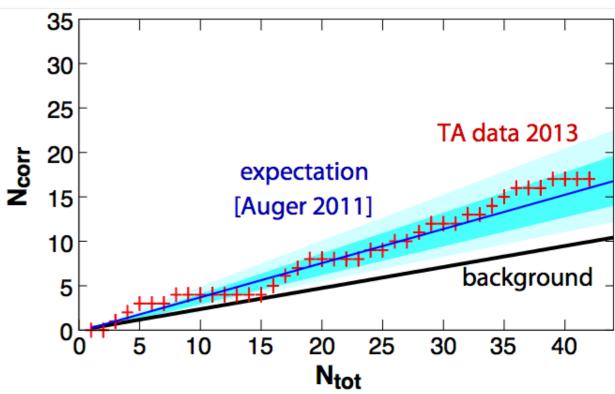
1.7 Arc Seconds 400 LIGHT-YEARS

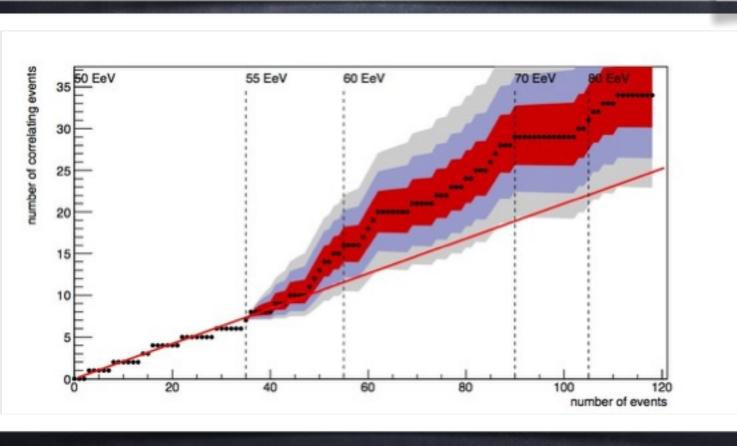
Or Cygnus A



Status of Large Scale UHECR Anisotropy







Kampert and Tinyakov, arXiv:1405.0575

All Sky View from Pierre Auger and Telescope Array

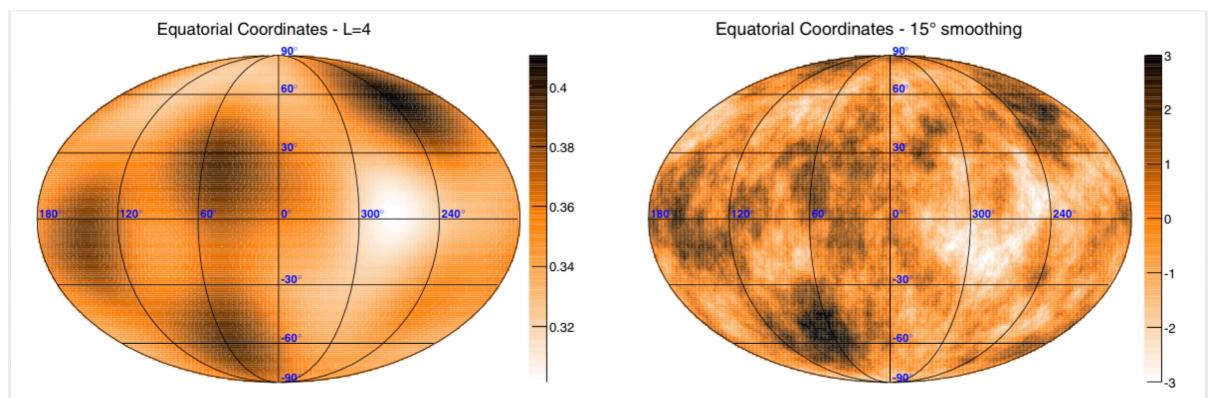
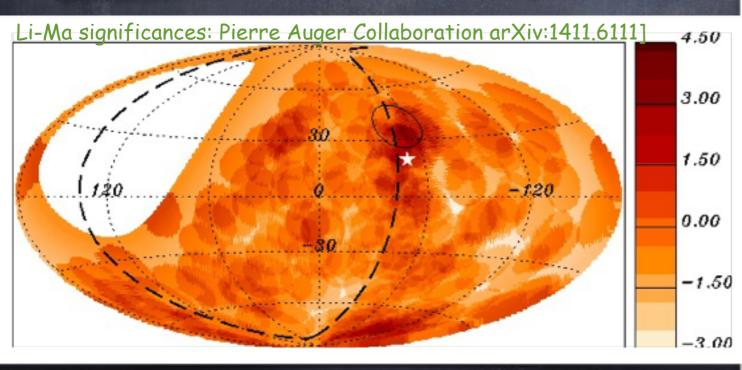
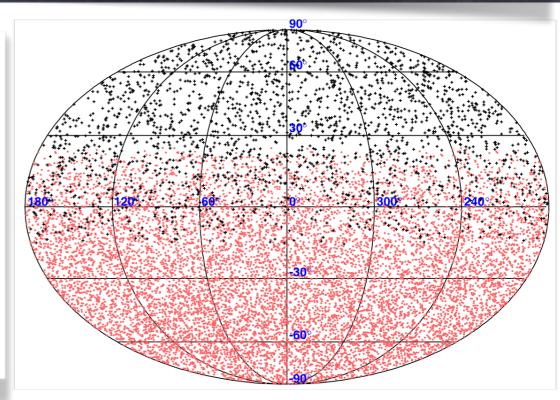


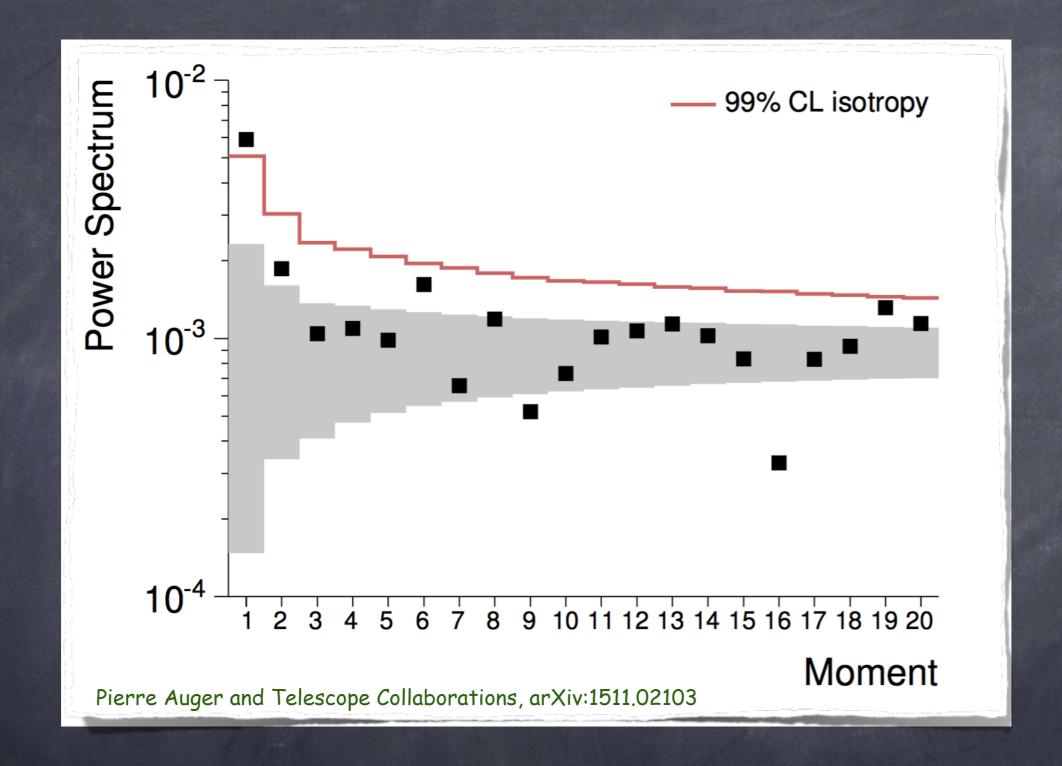
Figure 9. Left: flux sky map in km⁻² yr⁻¹ sr⁻¹ units, using a multipolar expansion up to $\ell = 4$. Right: significance sky map smoothed out at a 15° angular scale.

(A color version of this figure is available in the online journal.)

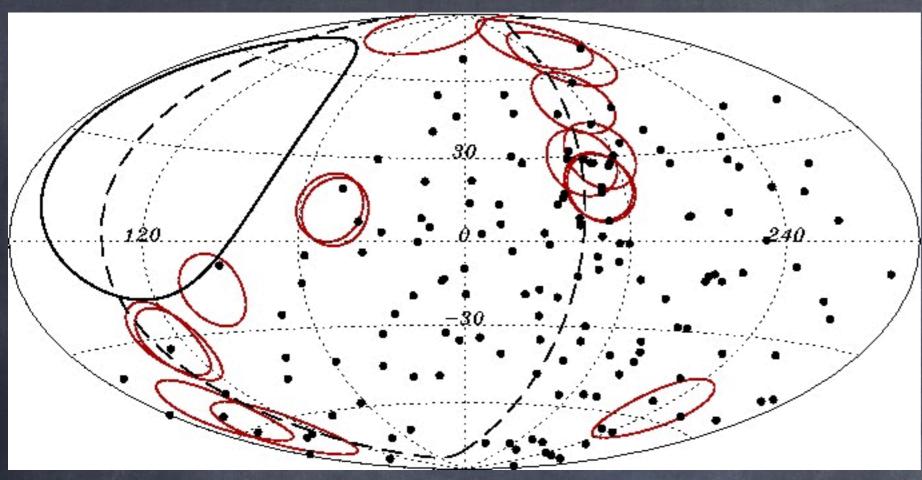
Pierre Auger and Telescope Collaborations, ApJ 794 (2014) 172 [arXiv:1409.3128]







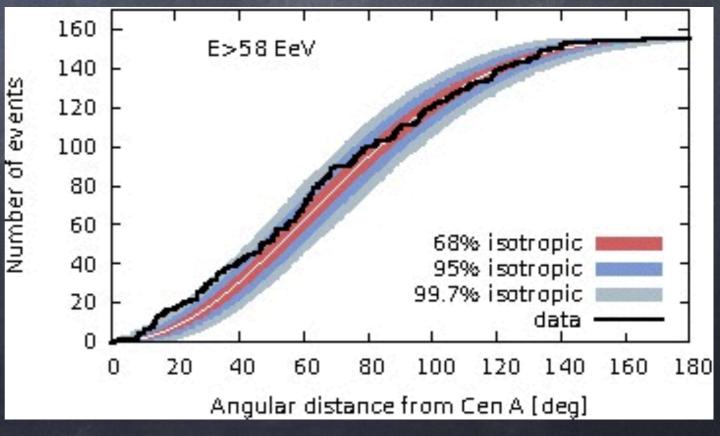
Centaurus A is a UHECR source candidate



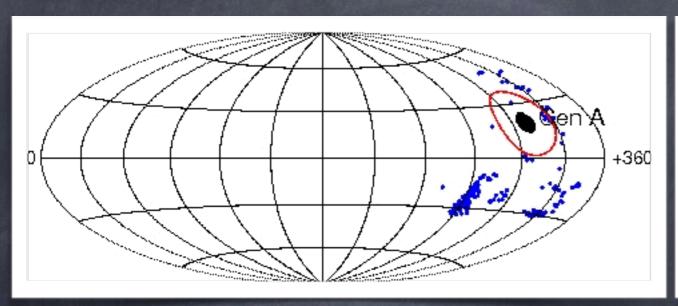
Pierre Auger Collaboration, arXiv:1411.6111

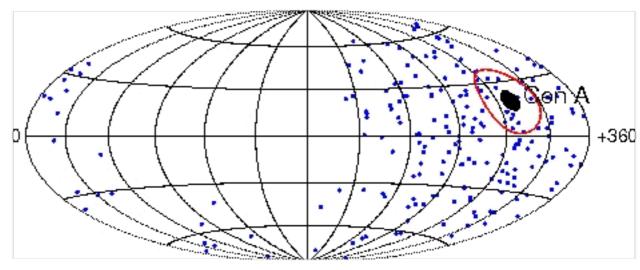
Events above 58 EeV against radio galaxies brighter than 10^{40} erg/s within 90 Mpc (red circles)

Pierre Auger sees a slight excess in the direction of Centaurus A above 58 EeV

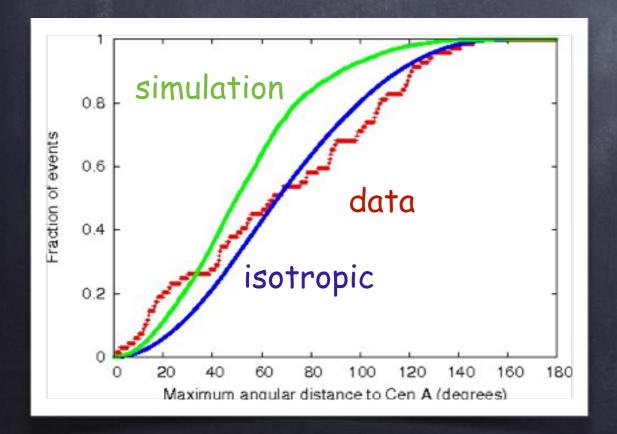


But even for iron primaries Centaurus A can not be the only UHECR source





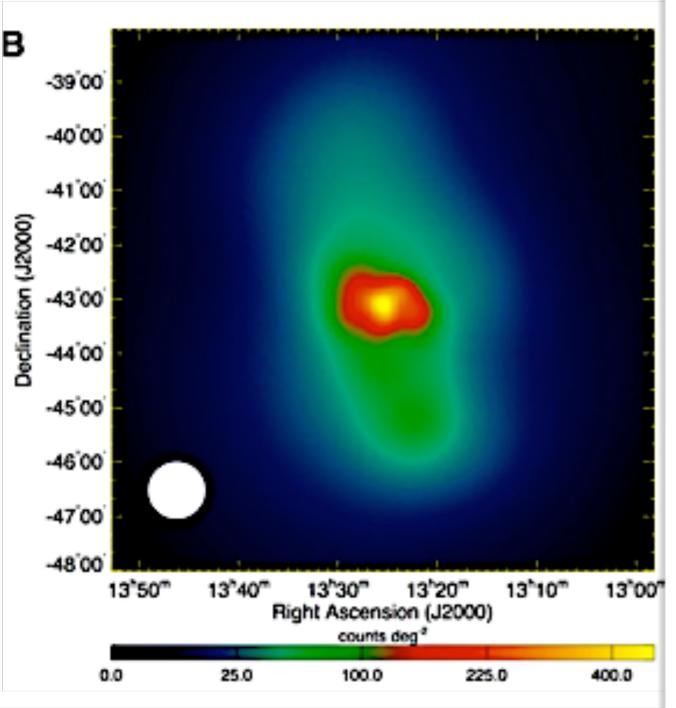
Iron Image of Cen A in the Prouza-Smida Galactic magnetic field model

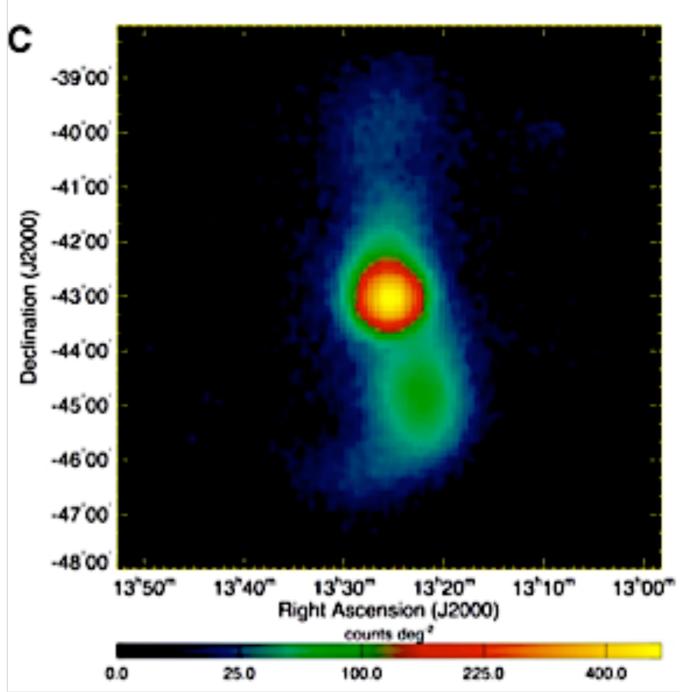


Including an extreme choice for the turbulent Galactic field component with strength 10 μ G, coherence length 50 pc, 10 kpc halo extension

Giacinti, Kachelriess, Semikoz, Sigl, Astropart. Phys. 35 (2011) 192

Lobes of Centaurus A seen by Fermi-LAT

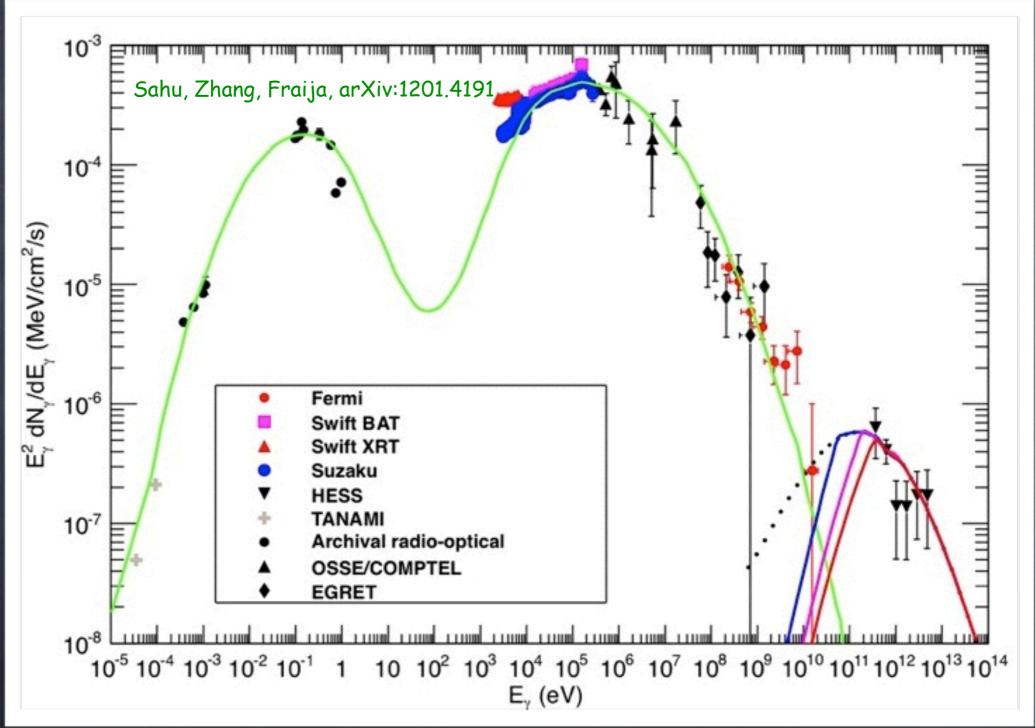




> 200 MeV y-rays

Radio observations

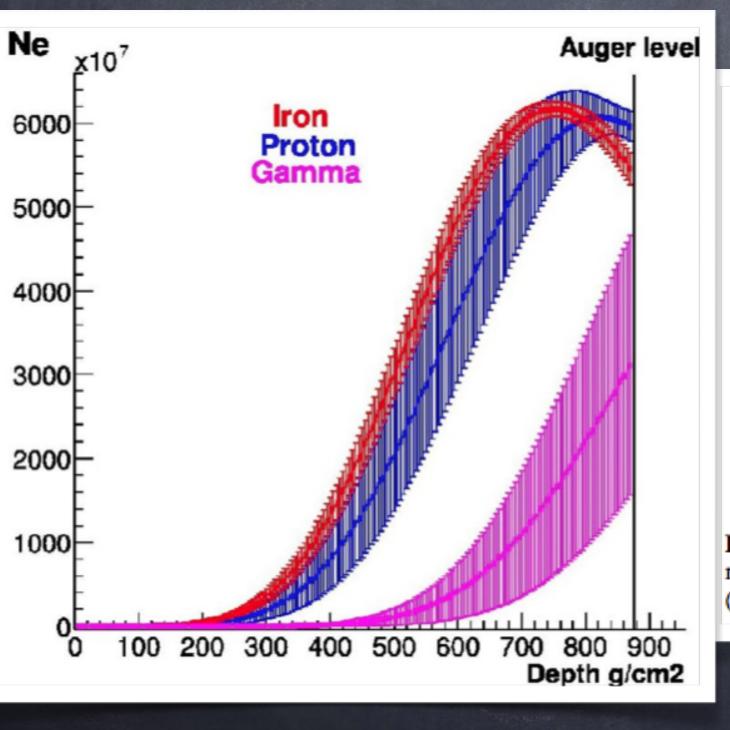
Centaurus A as Multimessenger Source: A Mixed hadronic+leptonic Model



Low energy bump = synchrotron
high energy bump = synchrotron self-Compton
TeV-y-rays: py interactions of shock-accelerated protons

Mass Composition

Depth of shower maximum X_{max} and its distribution contain information on primary mass composition



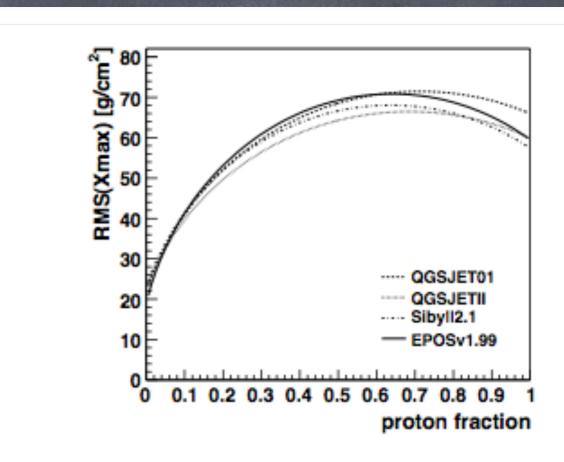
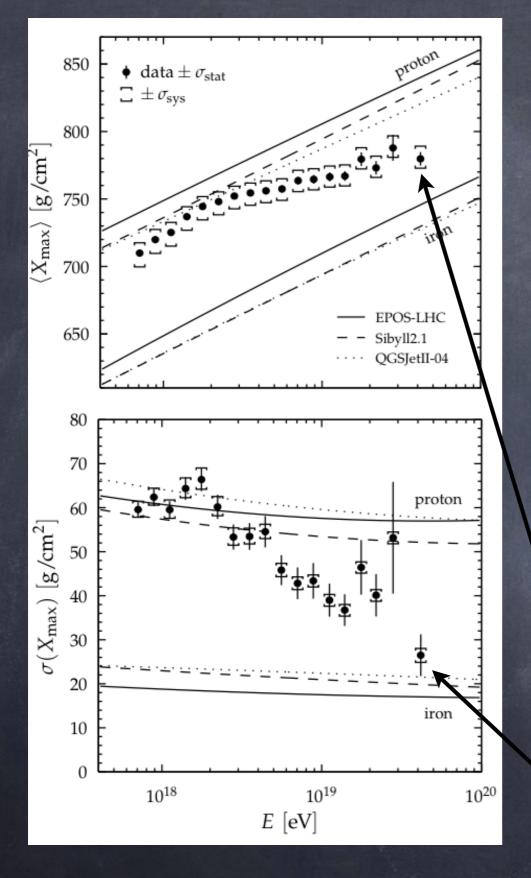
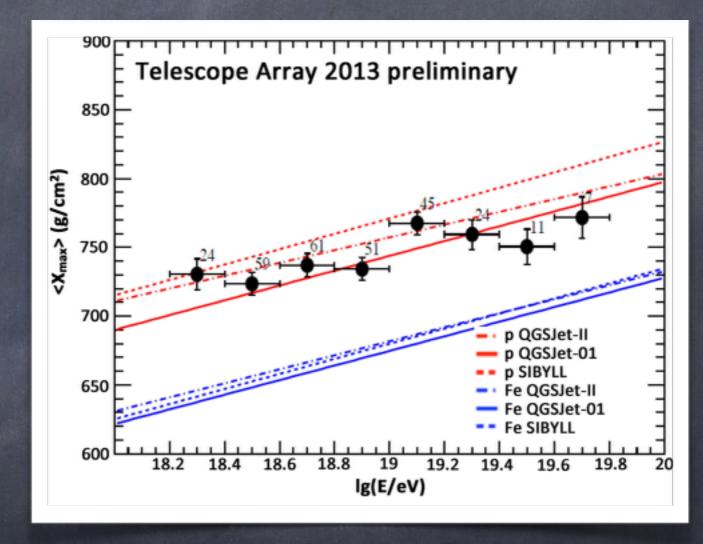


FIGURE 1. RMS(X_{max}) from different hadronic interaction models [23] and a two-component p/Fe composition model ($E = 10^{18} \text{ eV}$).

Pierre Auger data suggest a heavier composition toward highest energies:



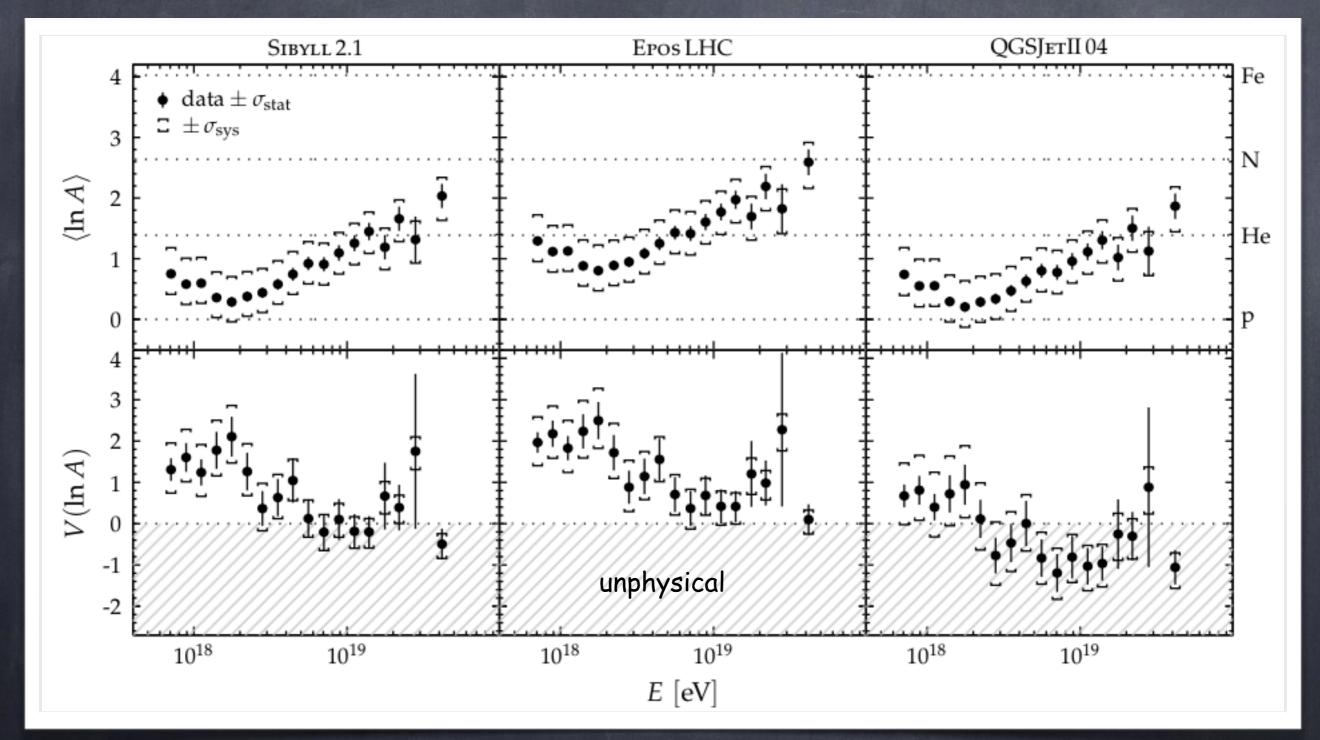
but not confirmed on the northern hemisphere by HiRes and Telescope Array which are consistent with protons



potential tension with air shower simulations and some hadronic interaction models because a mixed composition would predict larger RMS(X_{max})

Pierre Auger Collaboration, arXiv:1409.4809

combined measurement of X_{max} and its fluctuation $\sigma(X_{max})$ can be translated to distribution of atomic mass A within a given hadronic interaction model



An attempt to reconstruct individual elements

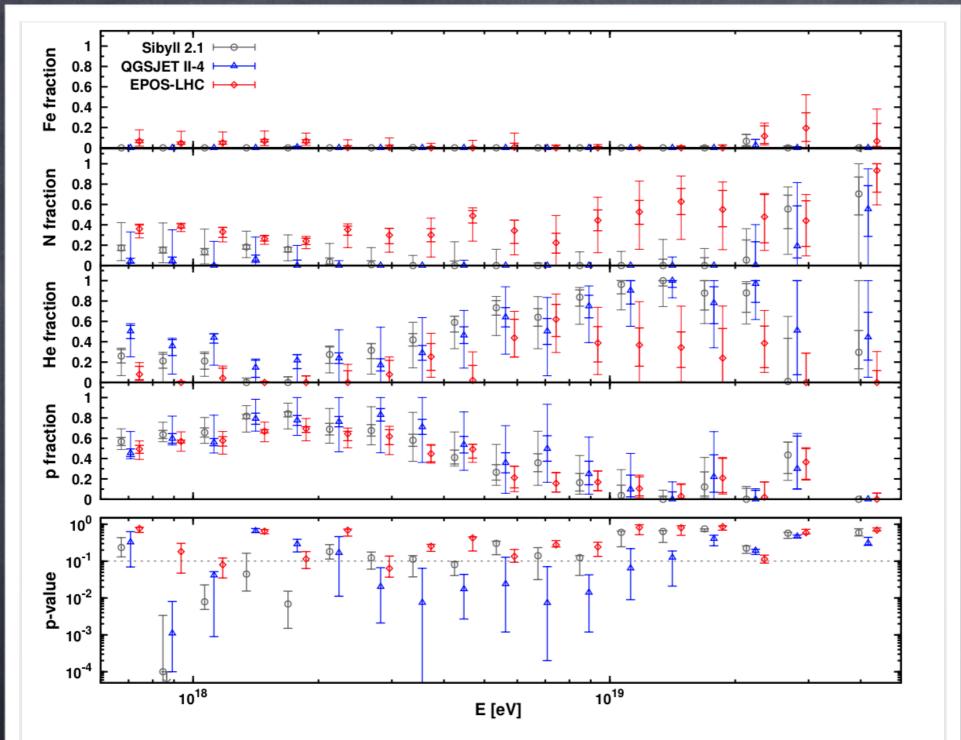
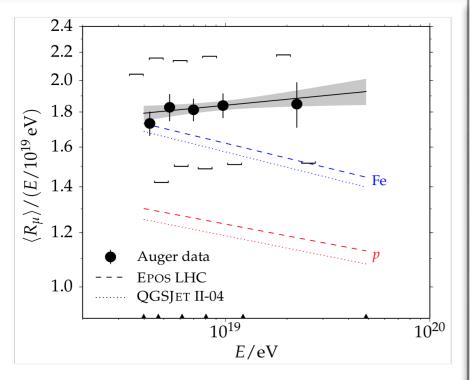
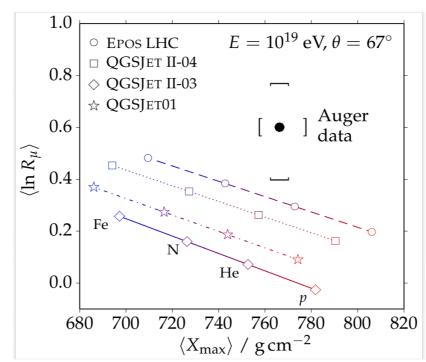
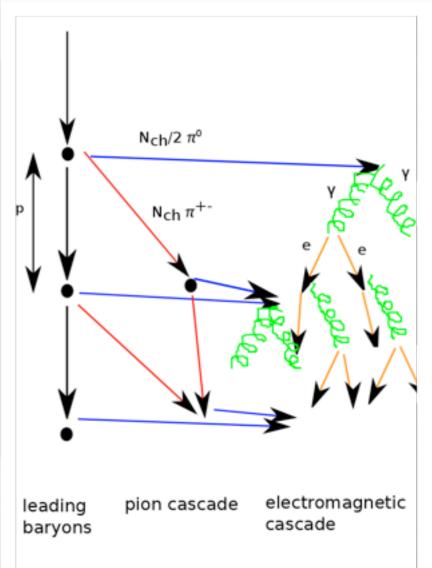


FIG. 4: Fitted fraction and quality for the scenario of a complex mixture of protons, helium nuclei, nitrogen nuclei, and iron nuclei. The upper panels show the species fractions and the lower panel shows the p-values.

Muon number measured at 1000 m from shower core systematically higher than predicted







Pierre Auger Collaboration, arXiv:1408.1421

The muon number scales as

$$N_{\mu} \propto E_{\mathrm{had}} \propto (1 - f_{\pi^0})^N$$
,

with the fraction going into the electromagnetic channel $f_{\pi^0} \simeq \frac{1}{3}$ and the number of generations N strongly constrained by X_{max} . Larger N_{μ} thus requires smaller f_{π^0} !

KASCADE data suggest a heavy composition below $\sim 10^{18}$ eV possibly becoming lighter around 10^{18} eV

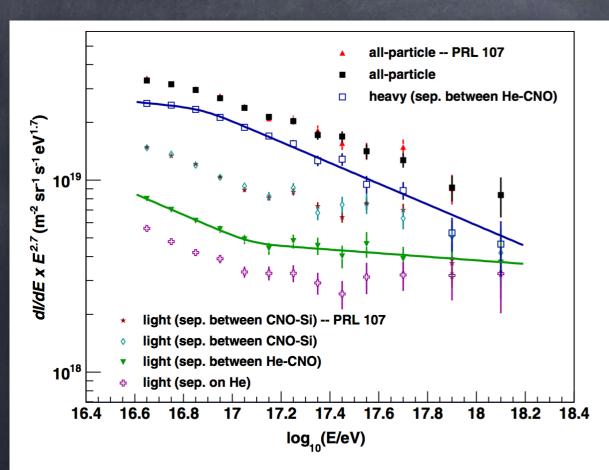
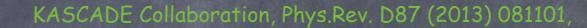
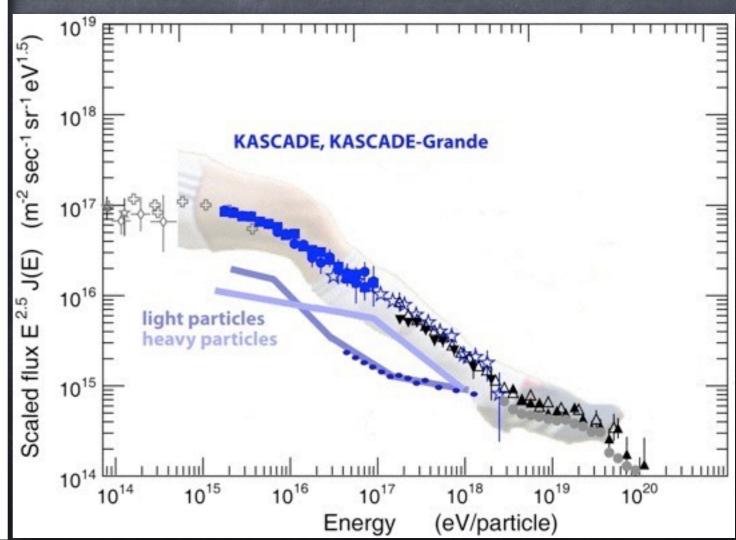
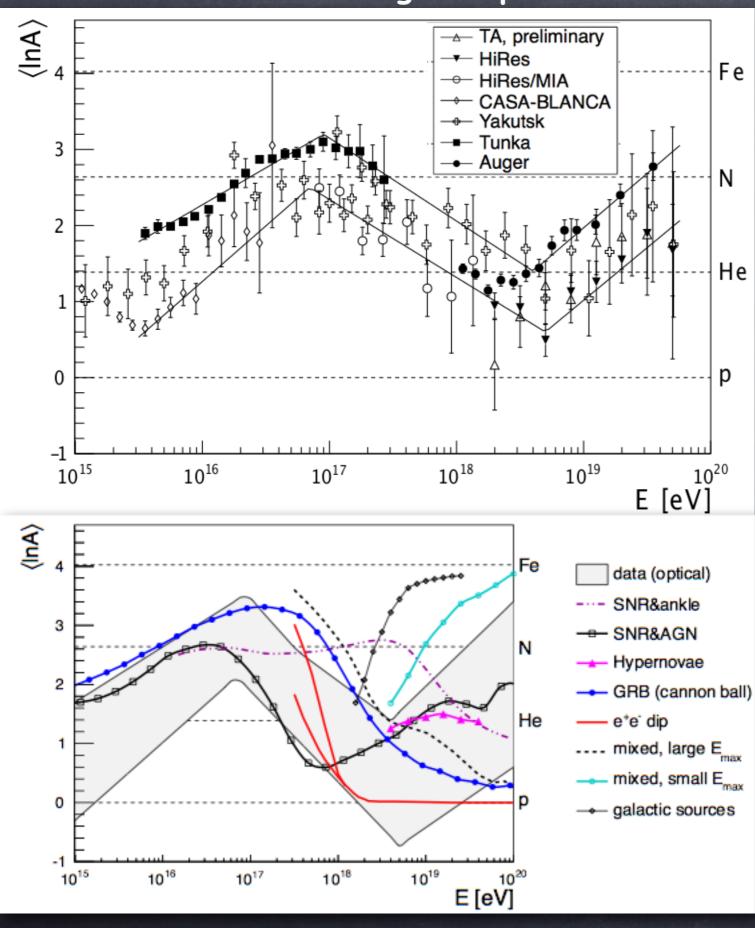


FIG. 4 (color online). The all-particle and electron-rich spectra from the analysis [8] in comparison to the results of this analysis with higher statistics. In addition to the light and heavy spectrum based on the separation between He and CNO, the light spectrum based on the separation on He is also shown. The error bars show the statistical uncertainties.



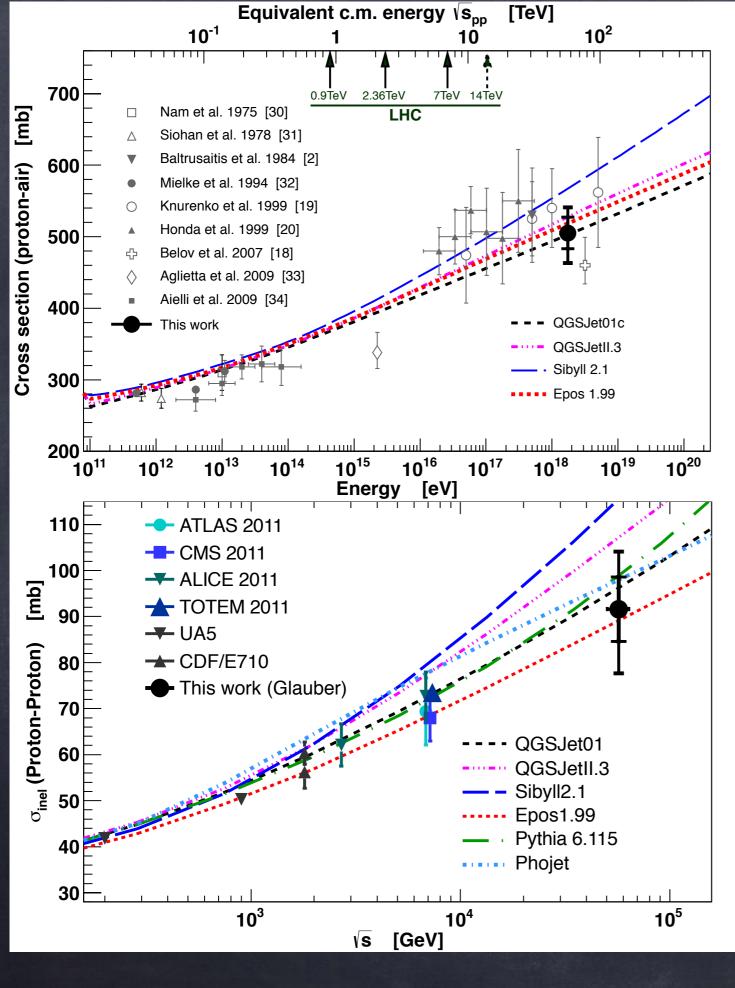


The global picture for the mass composition

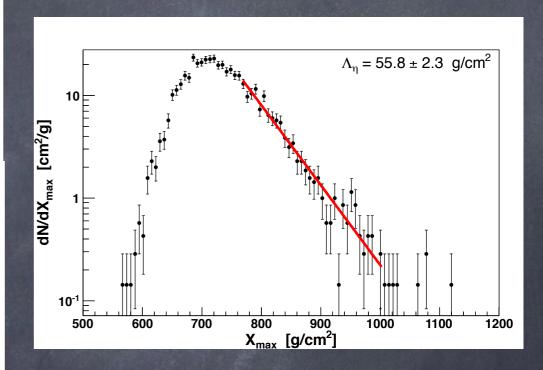


K.-H.Kampert and M.Unger, Astropart.Phys. 35 (2012) 660

Indications of "Peters cycles" for galactic and extragalactic sources whose maximal energies are proportional to the charge Z and extend up to $\sim 10^{17}$ and 10^{20} eV, respectively



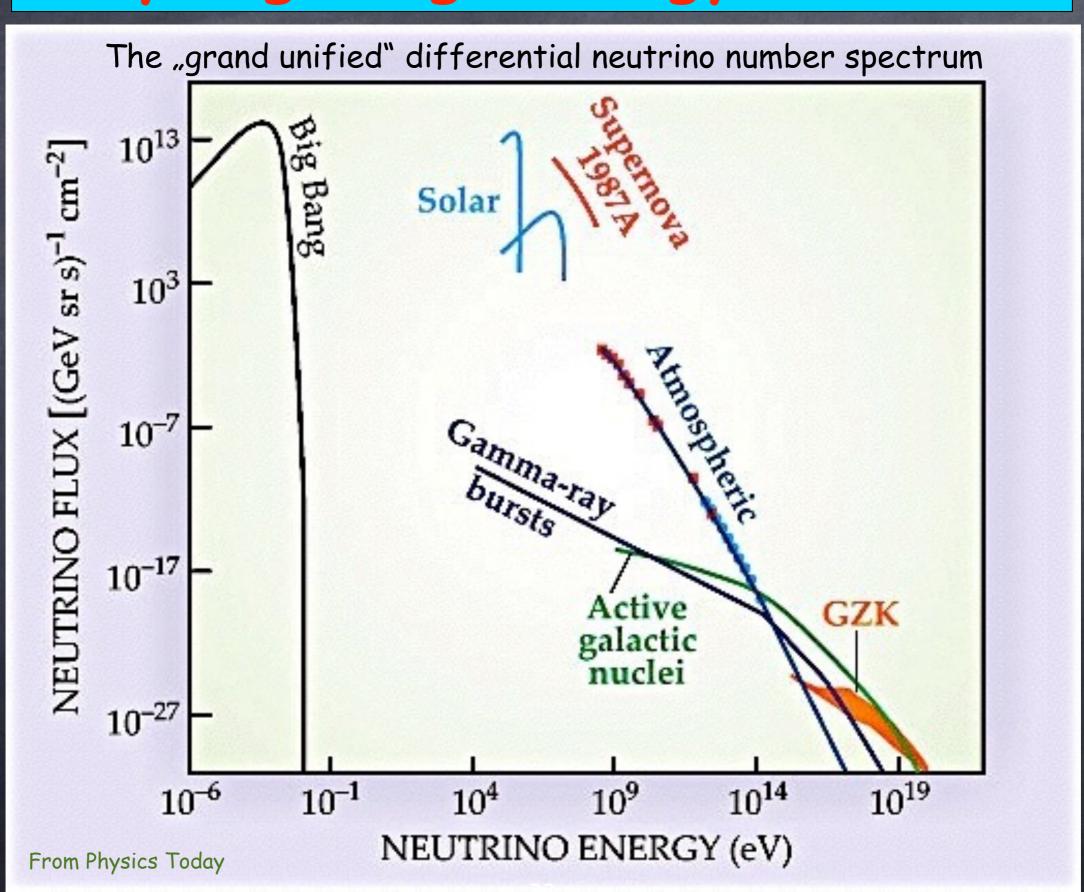
p-air cross section derived from exponential tail of depth of shower maxima



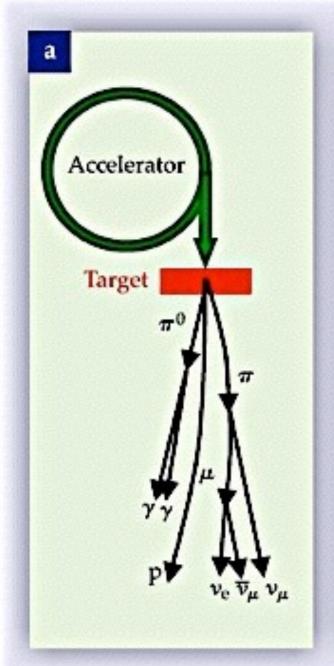
pp cross section derived from Glauber model

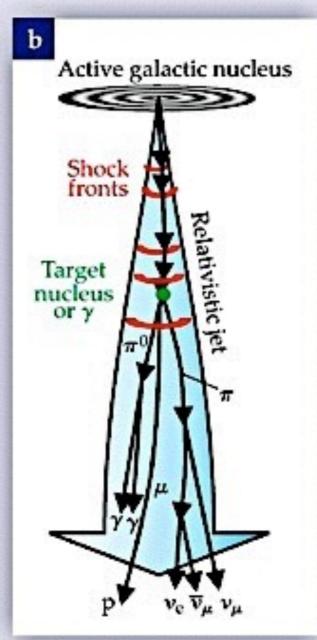
Pierre Auger Collaboration, PRL 109, 062002 (2012)

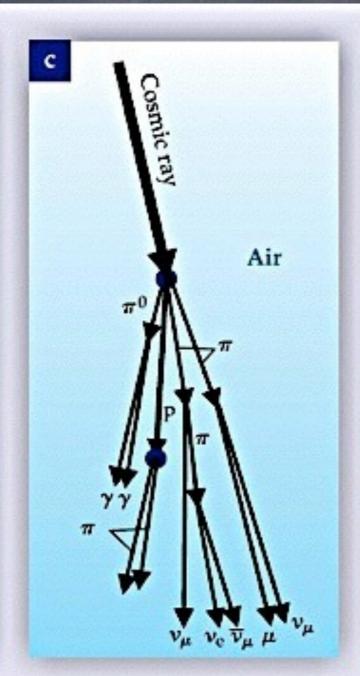
Very High High Energy Neutrinos

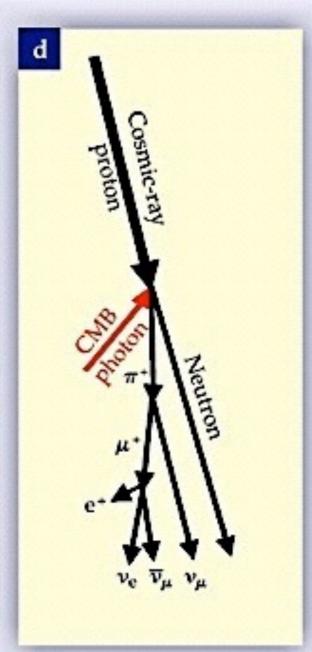


Summary of neutrino production modes









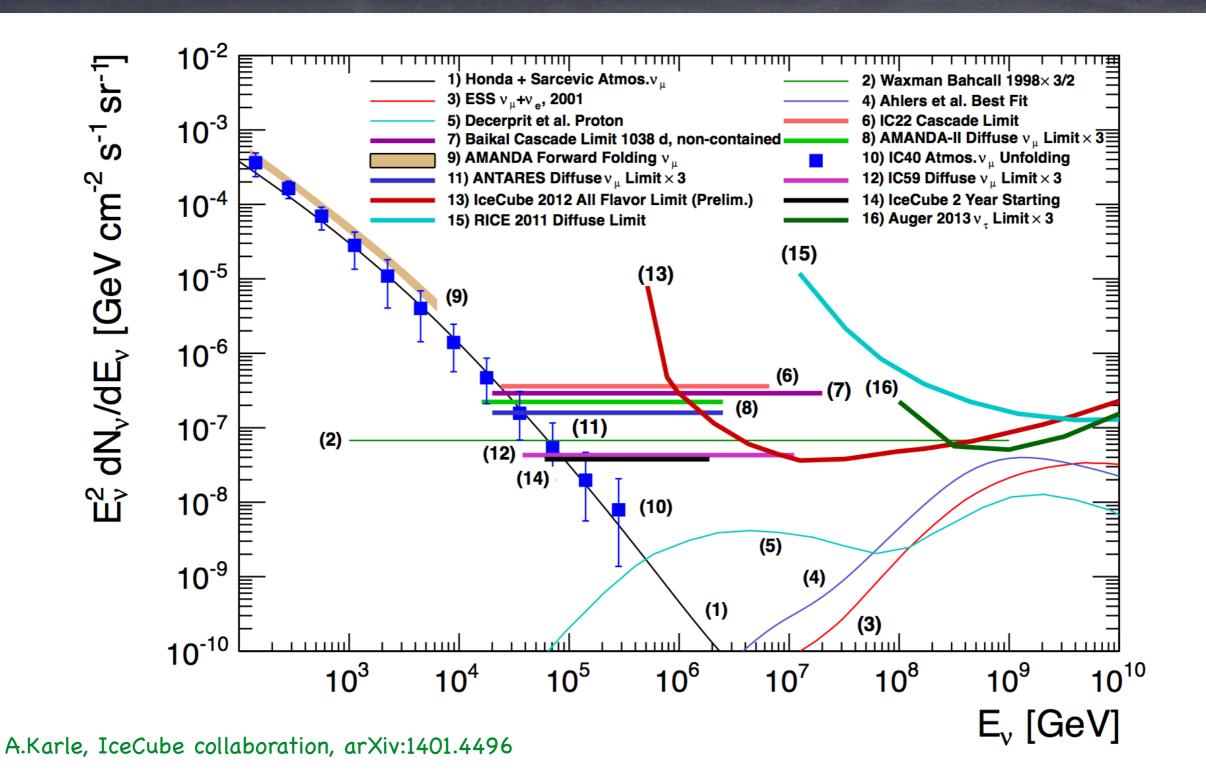
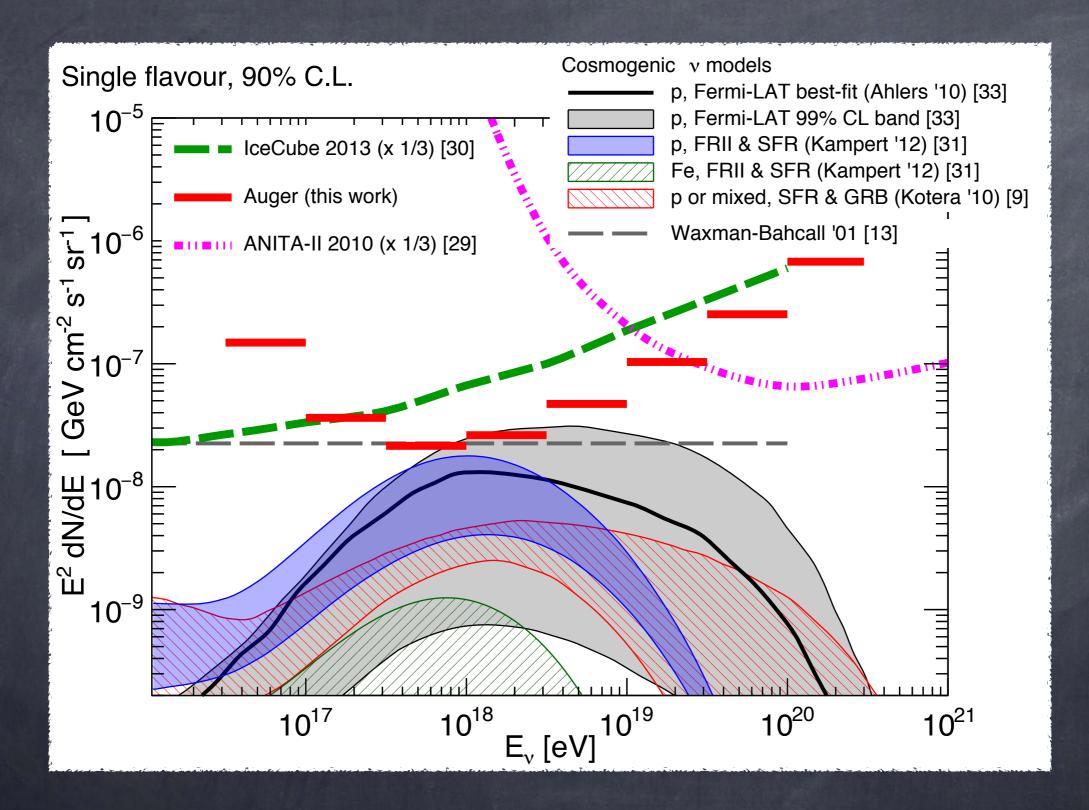


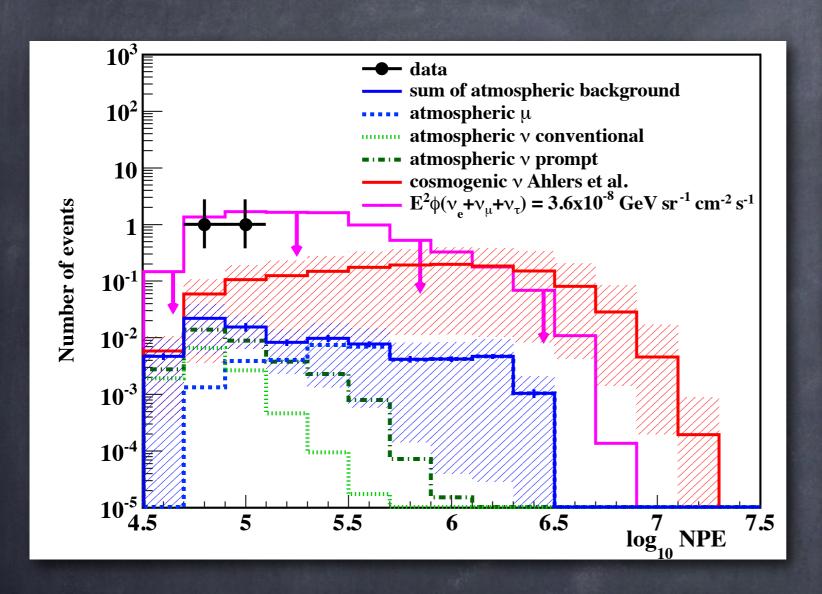
Figure 7: An overview is presented of observed atmospheric neutrino fluxes, upper limits to diffuse fluxes and models. The IceCube 2012 differential upper limit (11) turn up sharply at 1PeV because of observed PeV events. The best fit diffuse flux using starting events in IceCube (12) forms evidence for a diffuse astrophysical flux up to PeV energies above the atmospheric neutrino spectrum extending to a few 100 TeV.

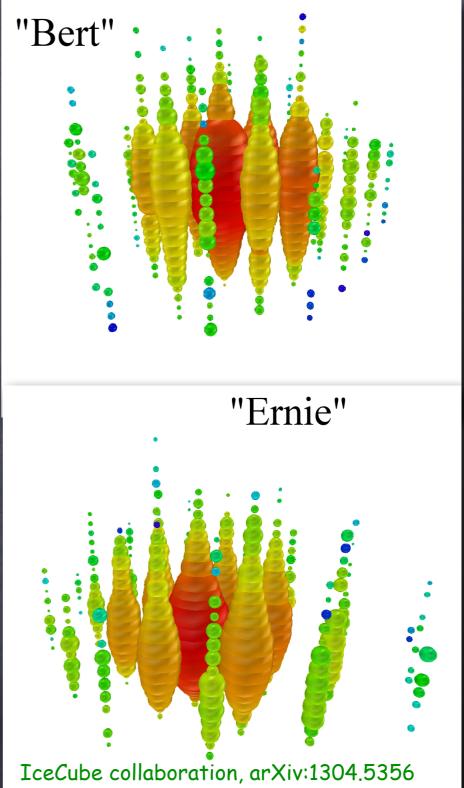


Pierre Auger collaboration, arXiv:1504.05397

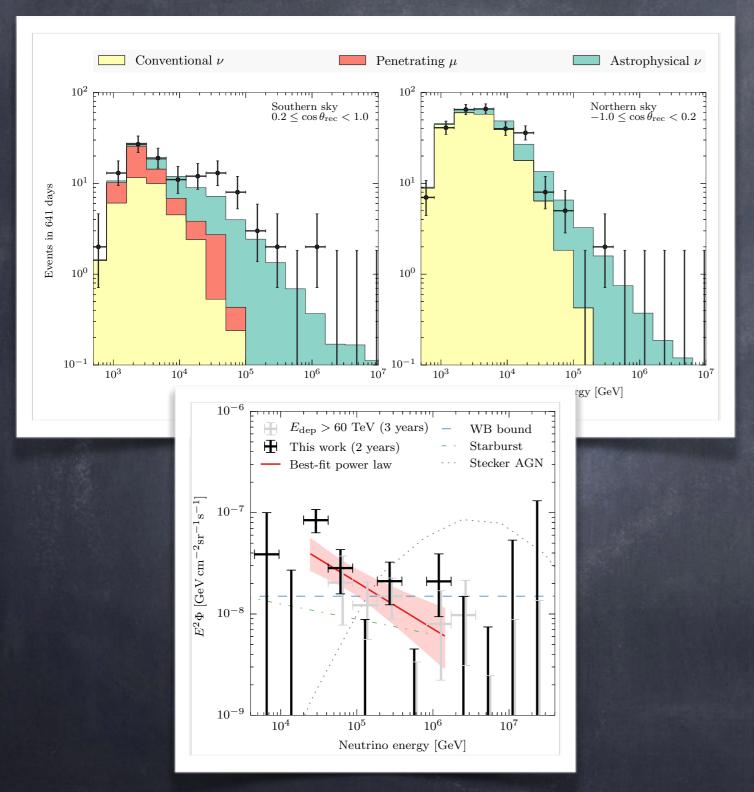
But now two PeV energy candidate neutrinos observed

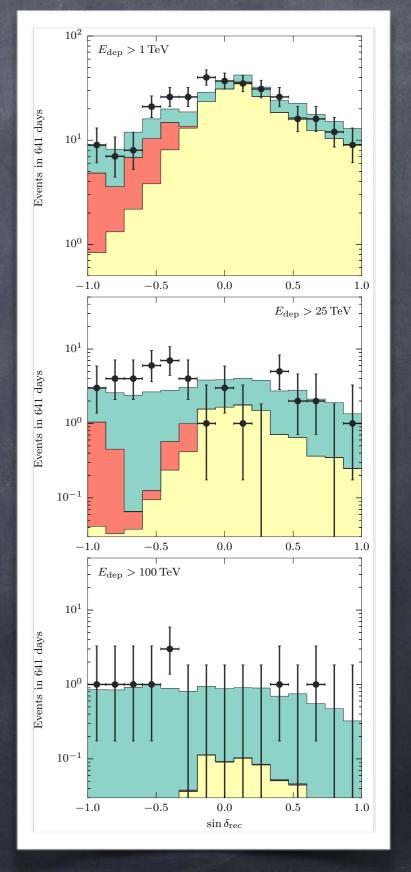
by IceCube





IceCube observed 283 cascade and 105 track events from the Southern sky above 1 TeV deposited energy:



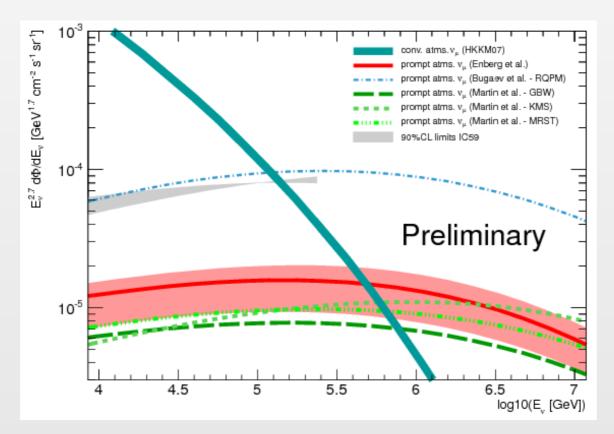


The "background": atmospheric ν

- * To assess the entity of an IceCube diffuse signal of purely astrophysical origin an accurate estimation of the background is mandatory.
- * Atmospheric neutrinos are a source of background:

Cosmic Rays + Atmospheric Nuclei \rightarrow hadrons \rightarrow neutrinos

- * Two contributing mechanisms, following two different power-law regimes:
 - conventional ν flux from the decay of π^{\pm} and K^{\pm}
 - prompt ν flux from charmed and havier hadrons (D, B)



Garzelli, Moch, Sigl, JHEP 1510 (2015) 115

Transition point: still subject of investigation ([IceCube collab., [arXiv:1302.0127]]).

The QCD core of the Z-moments for prompt fluxes:

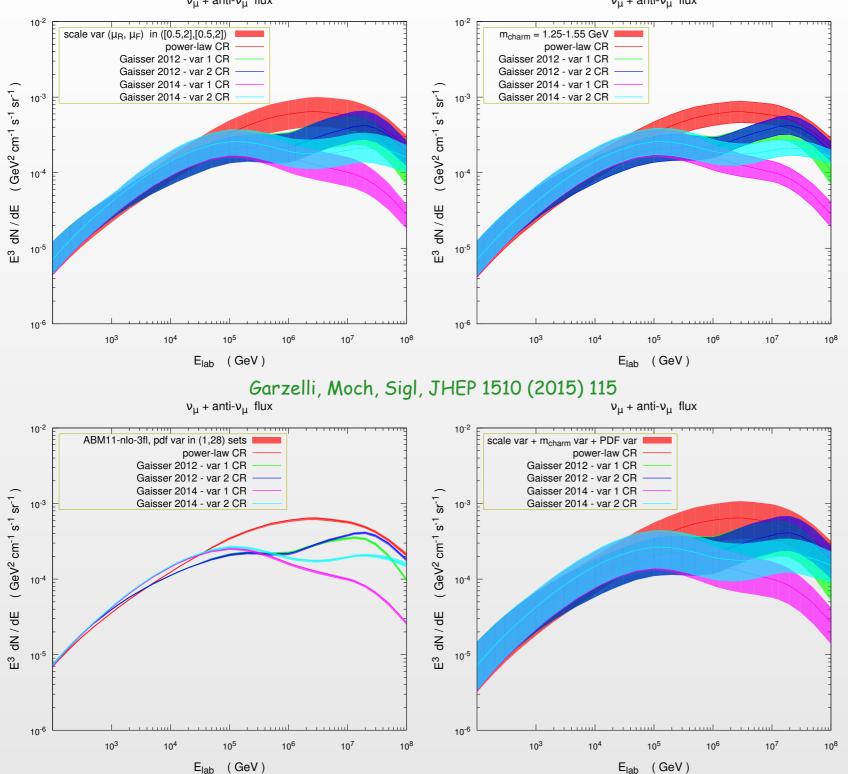
 $d\sigma(pp \rightarrow charmed\ hadrons)/dx_E$

$$Z_{ph}(E_h) = \int_0^1 \frac{dx_E}{x_E} \frac{\phi_p(E_h/x_E, 0)}{\phi_p(E_h, 0)} \frac{\lambda_p(E_h)}{\lambda_p(E_h/x_E)} \frac{A_{air}}{\sigma_{p-Air}^{tot, inel}(E_h)} \frac{d\sigma_{pp \to c\bar{c} \to h+X}}{dx_E} (E_h/x_E)$$

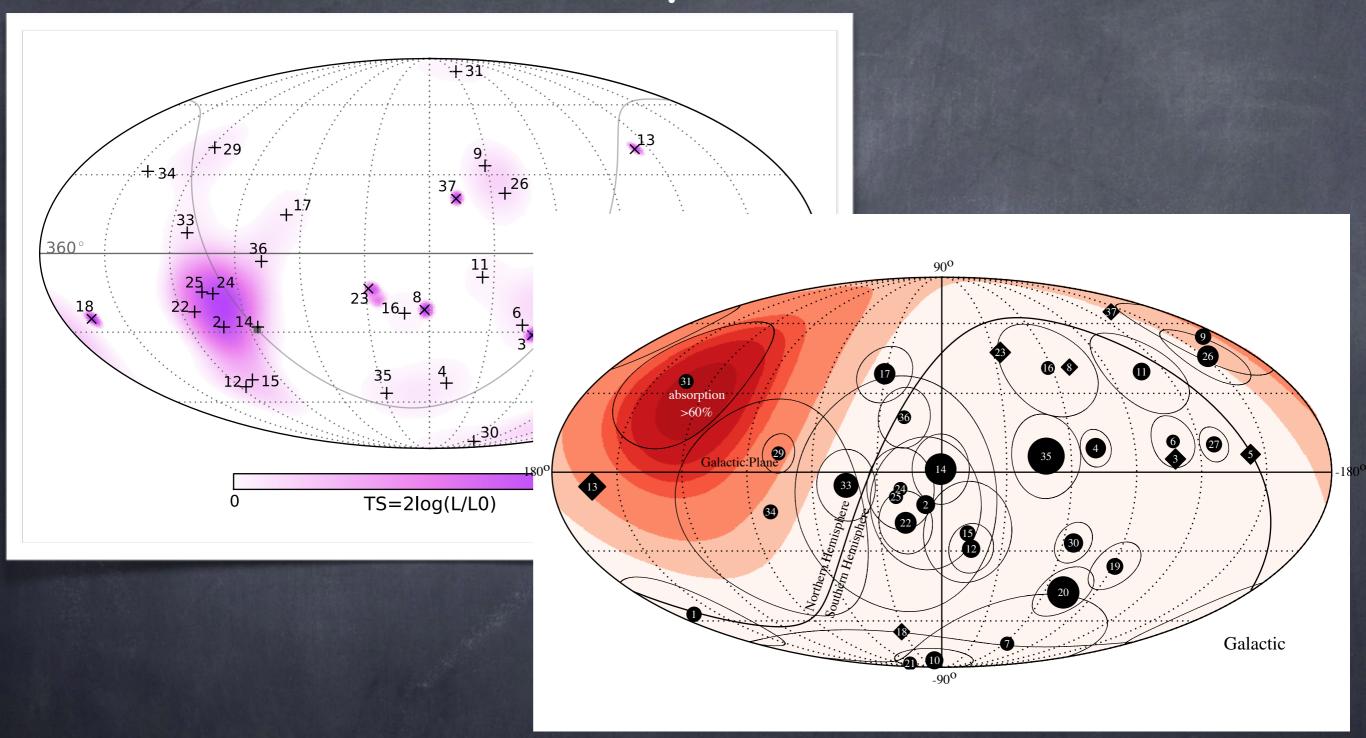
We used a (NLO QCD + Parton Shower + Hadronization) approach, with central scale, PDF and m_{charm} choices driven by previous considerations (see LO/NLO/NNLO plots) and variations in the following intervals:

- central scale $(\mu_R, \mu_F) = \mu_0 = \sqrt{p_{T,charm}^2 + 4m_{charm}^2}$, with independent variations of $\mu_R \in (0.5, 2)\mu_0$ and $\mu_F \in (0.5, 2)\mu_0$
- m_{charm}^{pole} = 1.40 GeV, with variation in [1.25,1.55] GeV
- PDFs:
 - * ABM11-NLO-3fl full set (central + 28 variations)
 - * CT10-nlo-3fl (central)
 - * NNPDF3.0-3fl (central)

$(\nu_{\mu} + \bar{\nu}_{\mu})$ fluxes: (scale + mass + PDF) variation ν_{μ} + anti- ν_{μ} flux ν_{μ} + anti- ν_{μ} flux summary

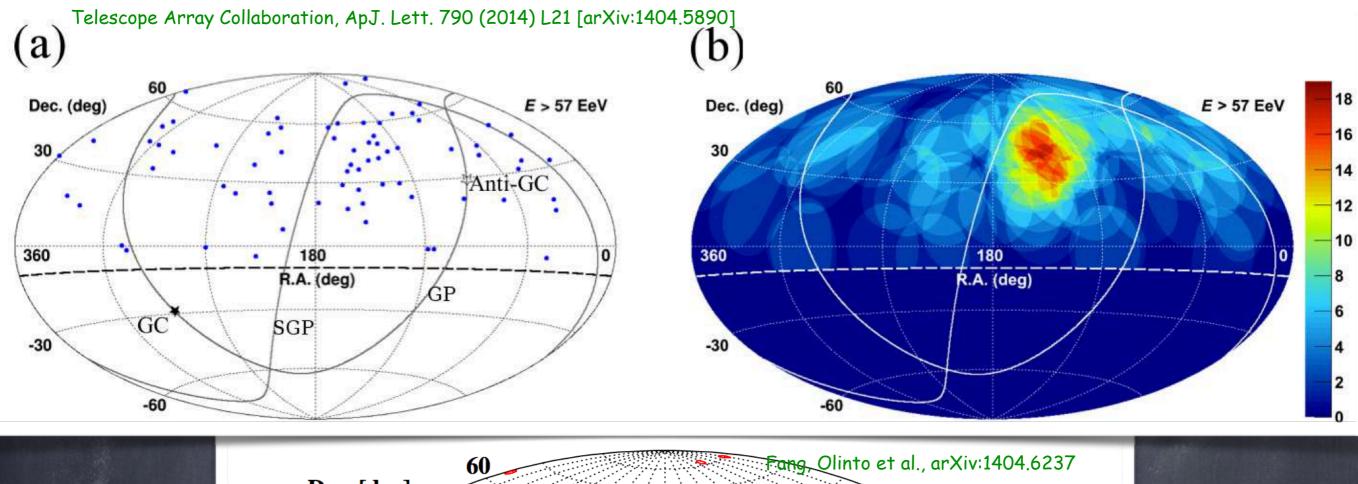


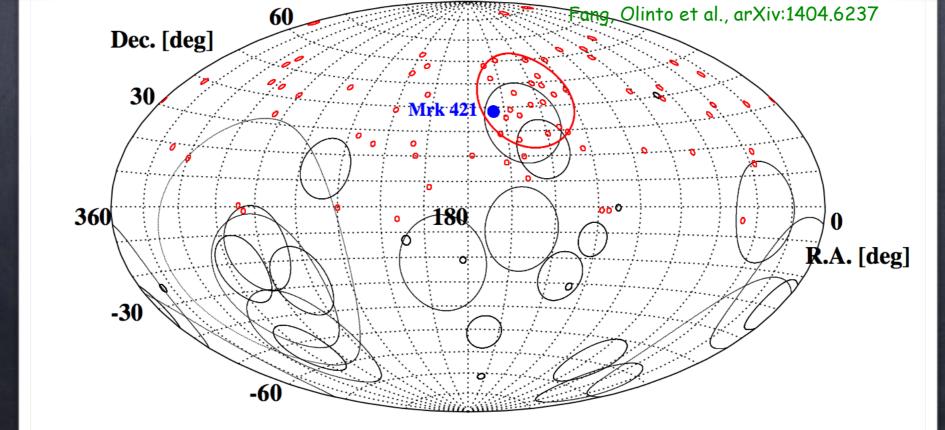
Neutrino sky distribution



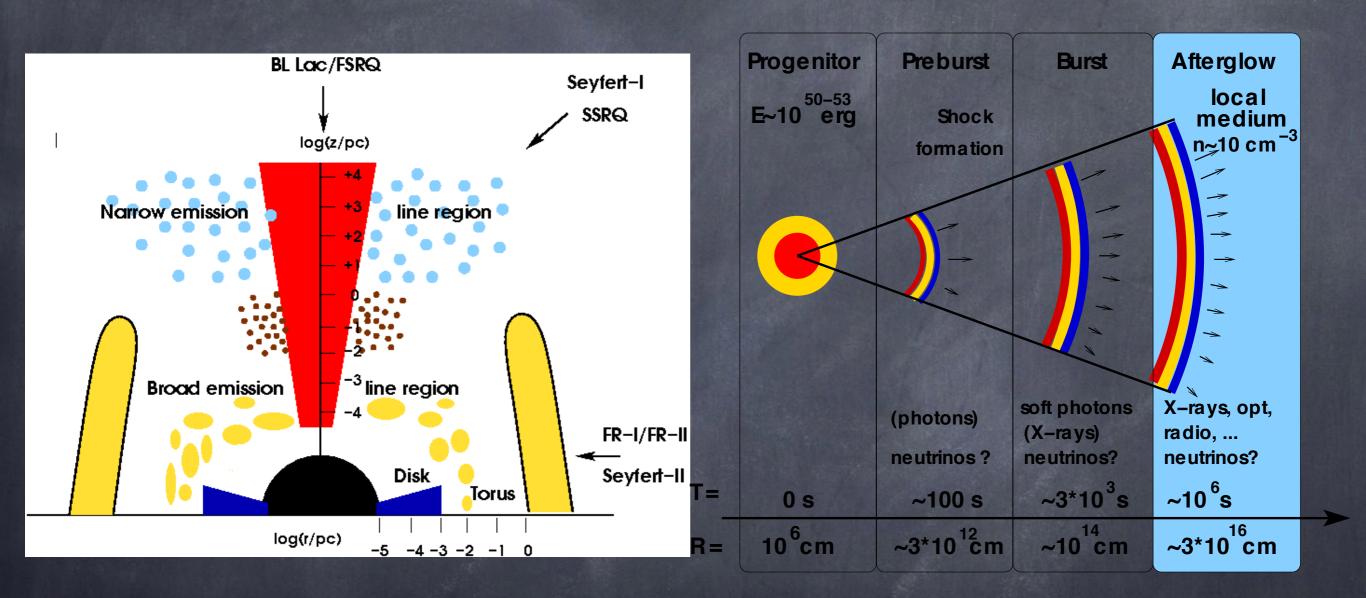
IceCube collaboration, Phys.Rev. Lett. 113 (2014) 101101 [arXiv:1405.5303]

A possible Correlation of IceCube Neutrinos with the Cosmic Ray Excess seen by Telescope Array?





Discrete Extragalactic High Energy Neutrino Sources



active galaxies

gamma ray bursts

Figures from J. Becker-Tjus, Phys.Rep. 458 (2008) 173

Neutrino Fluxes from Gamma-Ray Bursts

GRBs are optically thick to charged cosmic rays and nuclei are disintegrated => only neutrons escape and contribute to the UHECR flux by decaying back into protons

Diffuse neutrino flux from GRBs can thus be linked to UHECR flux (if it is dominantly produced by GRBs)

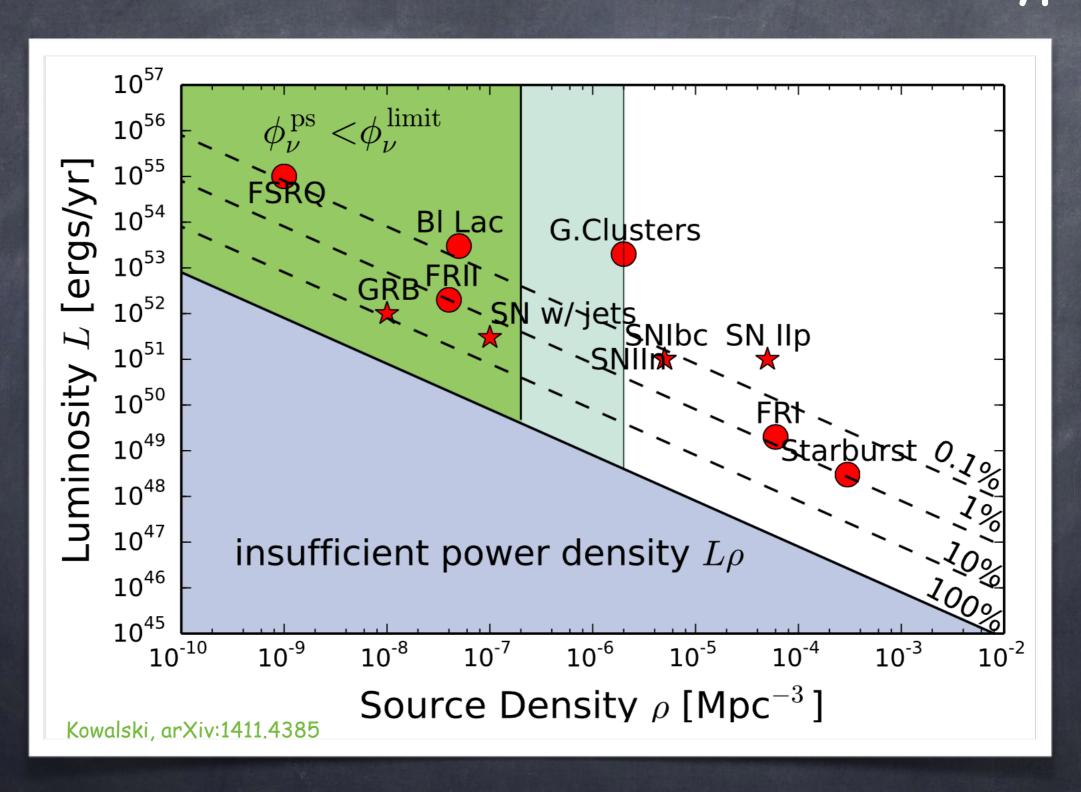
$$\Phi_{\nu}(E_{\nu}) \sim \frac{1}{\eta_{\nu}} \Phi_{p} \left(\frac{E}{\eta_{\nu}}\right) ,$$

where $\eta_{\nu} \simeq 0.1$ is average neutrino energy in units of the parent proton energy.

Above ~ 10^{17} eV neutrino spectrum is steepened by one power of E $_{\rm v}$ because pions/muons interact before decaying

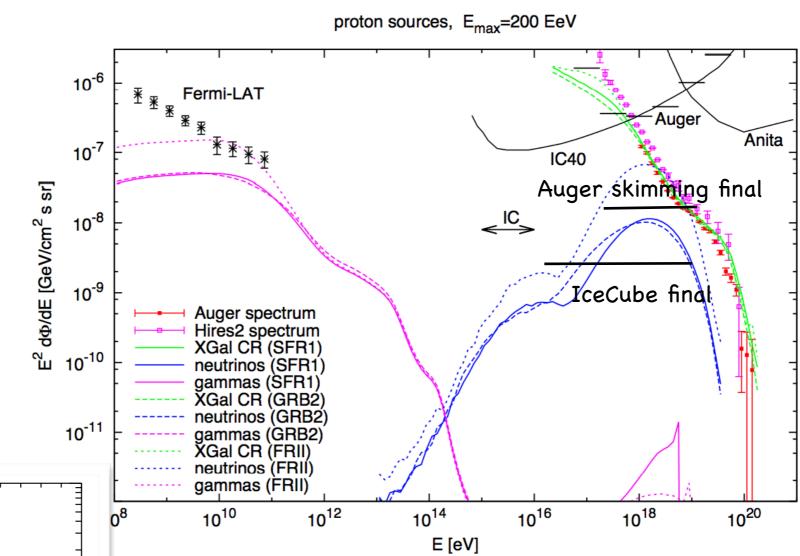
Correlation studies with GRBs now constrain the GRB contribution to observed diffuse neutrino flux to < 1%, see IceCube collaboration arXiv:1601.06484; also implies subdominant contribution of GRBs to ultra-high energy cosmic rays

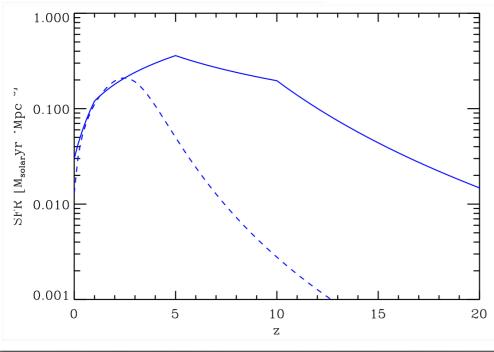
But a combination of the measured diffuse flux with upper limits on individual sources constrains source type



Cosmogenic Neutrinos: Maximal Fluxes for Pure Proton Injection insufficient to explain IceCube neutrinos

- Including secondary photons
- strong source evolution is here constrained by Fermi-LAT results





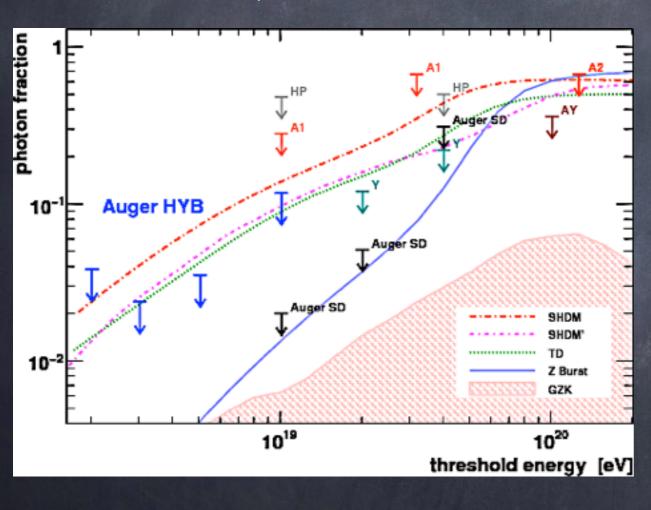
In scenario with $E_{\rm max}=200\,{\rm EeV}$ for different source evolution models (SFR1, GRB2 source spectral index is $\alpha=2.4$ for the SFR1 and GRB2 models, while $\alpha=2.2$ for Indicated are the propagated proton spectrum, the resulting (all flavor) neutrino luxes. The photon background measured by Fermi-LAT [10] is indicated, besides the ν bounds included in figure 1.

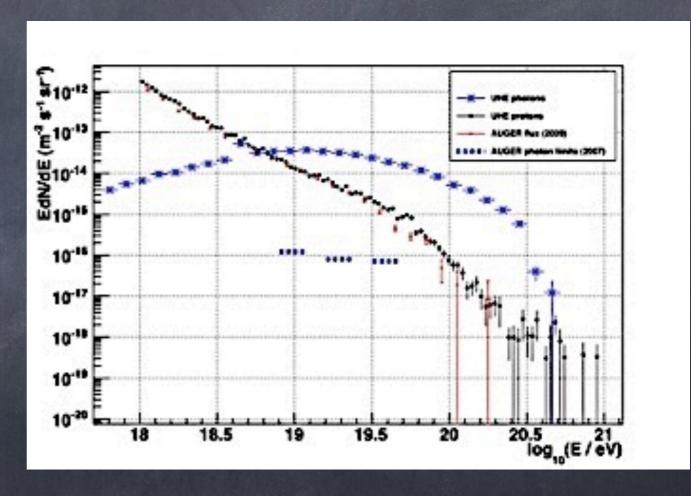
Lorentz Symmetry Violation in the Electromagnetic Sector

The idea:

Experimental upper limits on UHE photon fraction

Contradict predictions if pair production is absent





Pierre Auger Collaboration, Astropart. Phys. 31 (2009) 399 Maccione, Liberati, Sigl, PRL 105 (2010) 021101

Lorentz Symmetry Violation in the Photon Sector

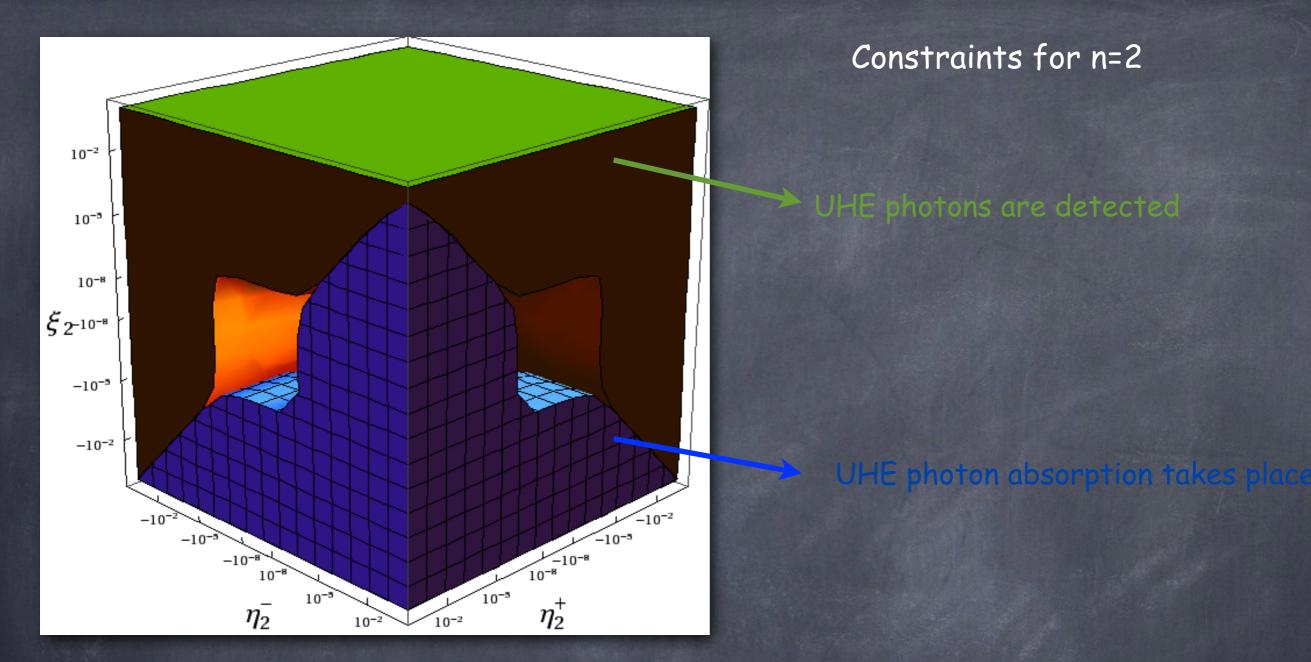
For a photon dispersion relation

$$\omega_{\pm}^2 = k^2 + \xi_n^{\pm} k^2 \left(\frac{k}{M_{\rm Pl}}\right)^n, n \ge 1,$$

pair production may become inhibited, increasing GZK photon fluxes above observed upper limits: In the absence of LIV for electrons/positrons for n=1 (CPT-odd terms) this yields:

$$\xi_1 \le 10^{-12}$$

Even for n=2 (CPT-even) one has sensitivity to $\xi_2 \sim 10^{-6}$ Such strong limits may indicate that Lorentz invariance violations are completely absent!

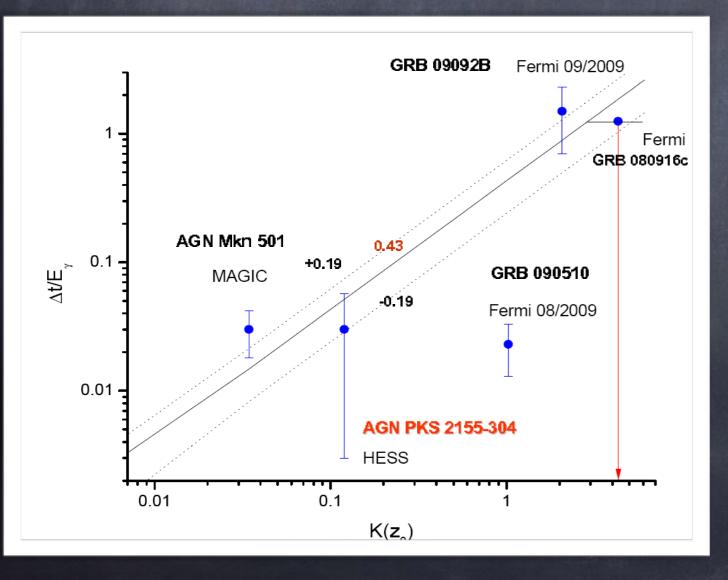


Such strong limits suggest that Lorentz invariance violations are completely absent!

The modified dispersion relation also leads to energy dependent group velocity $V=\delta E/\delta p$ and thus to an energy-dependent time delay over a distance d:

$$\Delta t = -\xi \, d \frac{E}{M_{\rm Pl}} \simeq -\xi \left(\frac{d}{100 \,{\rm Mpc}} \right) \left(\frac{E}{{\rm TeV}} \right) \,{\rm sec}$$

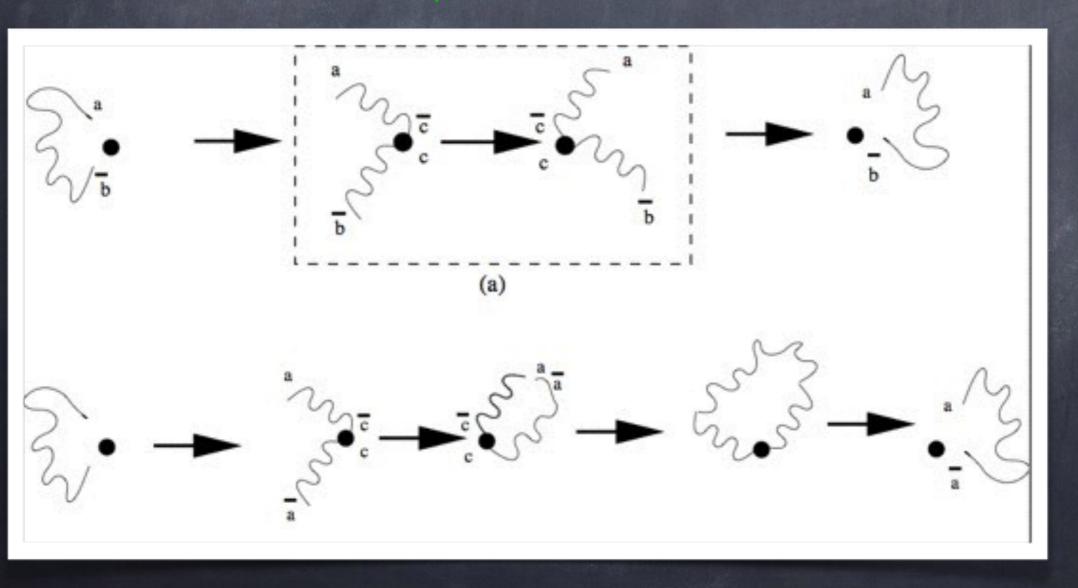
for linearly suppressed terms. GRB observations in TeV γ -rays can therefore probe quantum gravity and may explain that higher energy photons tend to arrive later (Ellis et al.).



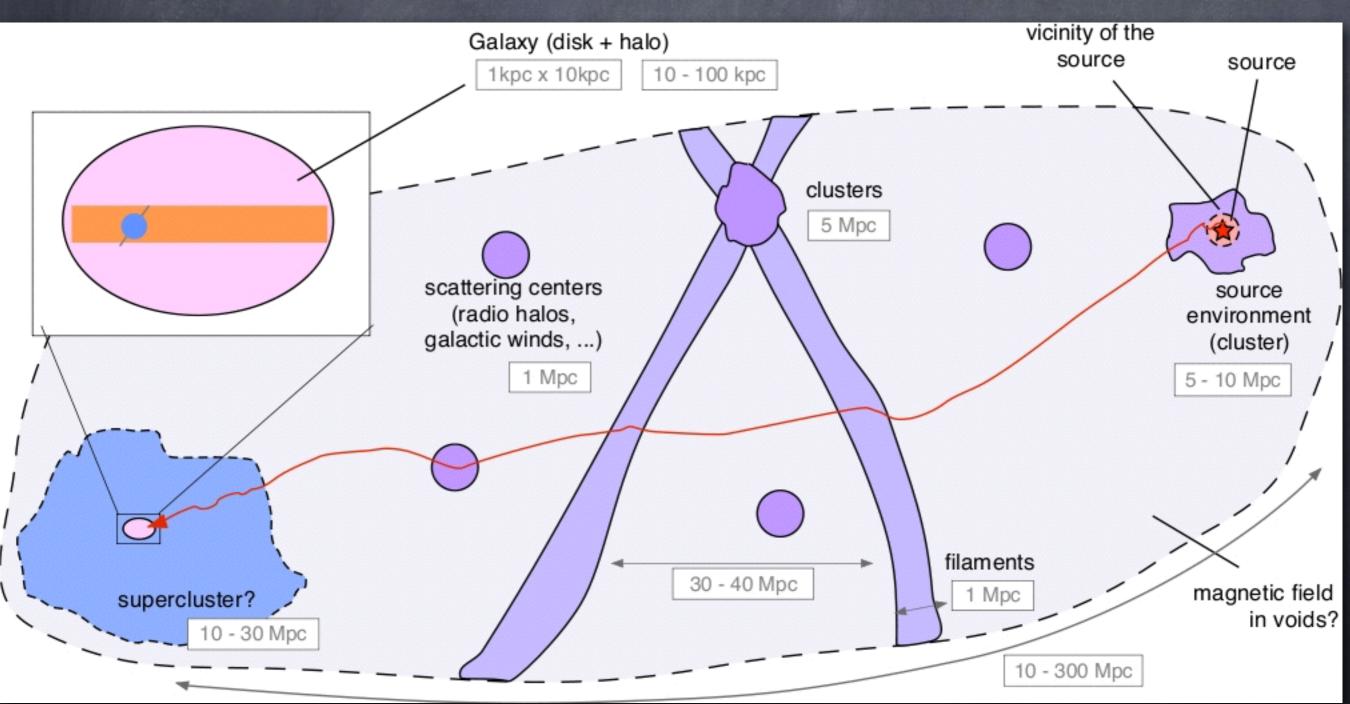
But the UHE photon limits are inconsistent with interpretations of time delays of high energy gamma-rays from GRBs within quantum gravity scenarios based on effective field theory

Maccione, Liberati, Sigl, PRL 105 (2010) 021101

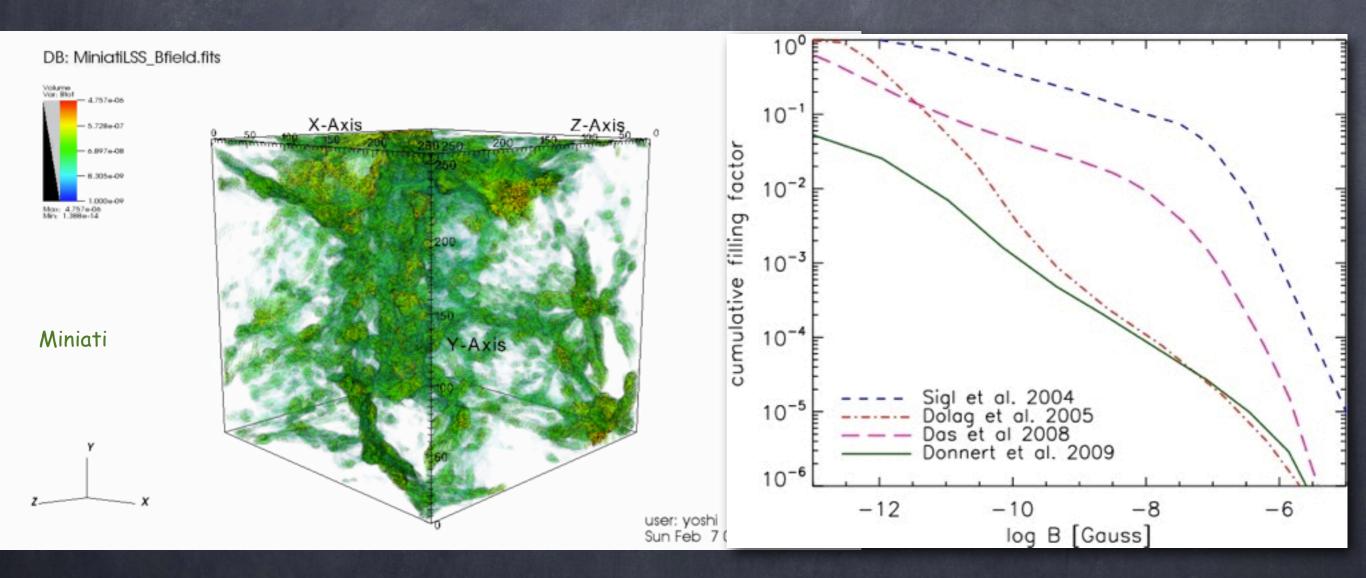
Possible exception in space-time foam models, Ellis, Mavromatos, Nanopoulos, arXiv:1004.4167



3-Dimensional Effects in Propagation



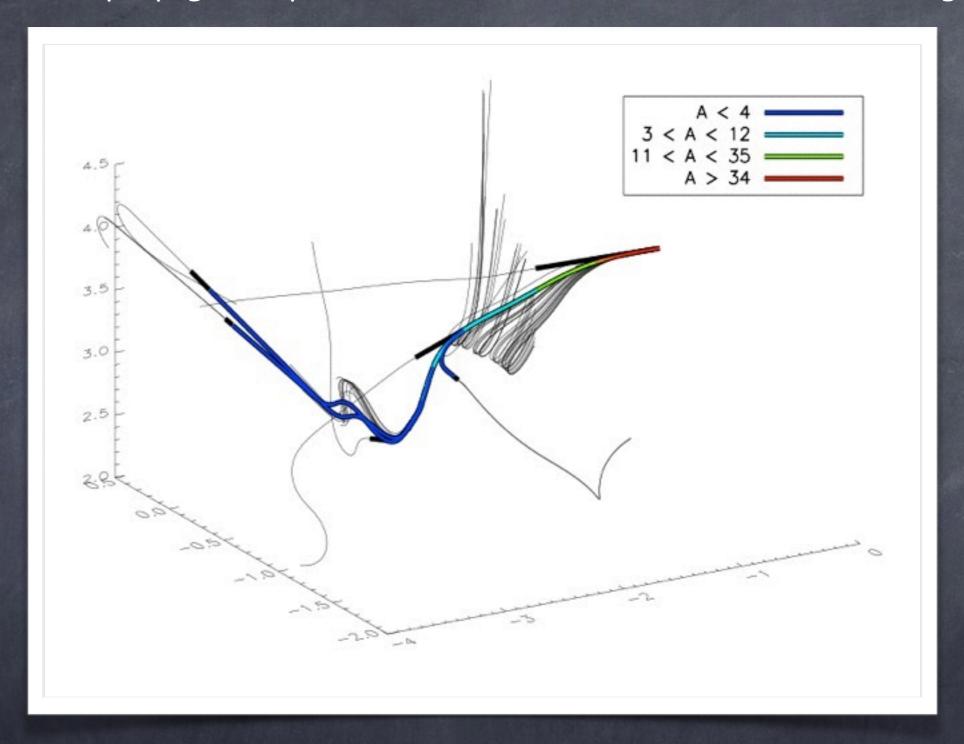
Structured Extragalactic Magnetic Fields



Kotera, Olinto, Ann. Rev. Astron. Astrophys. 49 (2011) 119

Filling factors of extragalactic magnetic fields are not well known and come out different in different large scale structure simulations

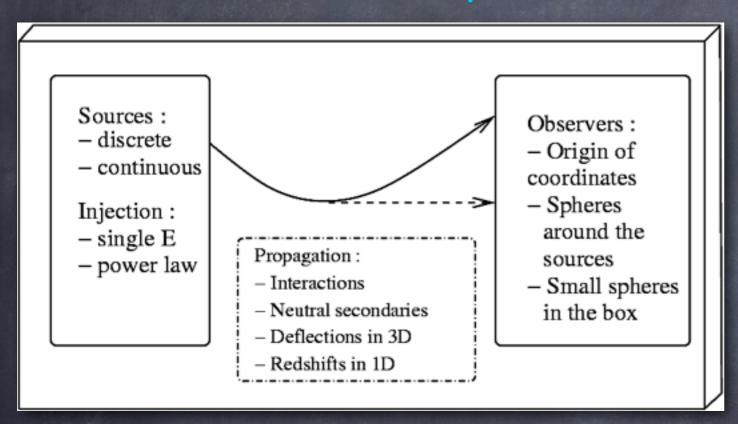
Extragalactic iron propagation produces nuclear cascades in structured magnetic fields:

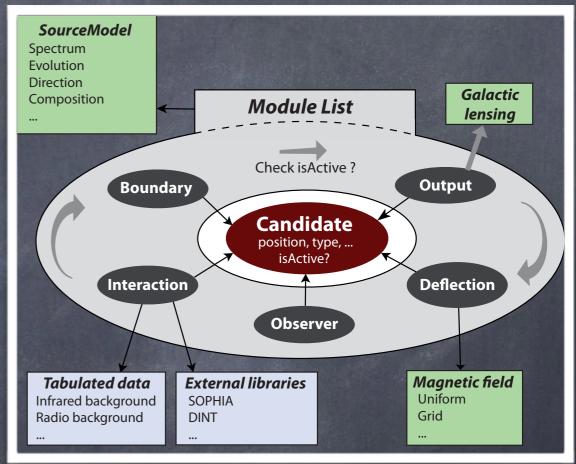


Initial energy 1.2 \times 10²¹ eV, magnetic field range 10⁻¹⁵ to 10⁻⁶ G. Color-coded is the mass number of secondary nuclei

CRPropa 2.0/3.0

CRPropa is a public code for UHE cosmic rays, neutrinos and y-rays being extended to heavy nuclei and hadronic interactions





Version 1.4: Eric Armengaud, Tristan Beau, Günter Sigl, Francesco Miniati, Astropart. Phys. 28 (2007) 463.

Version 2.0 at https://crpropa.desy.de/Main_Page

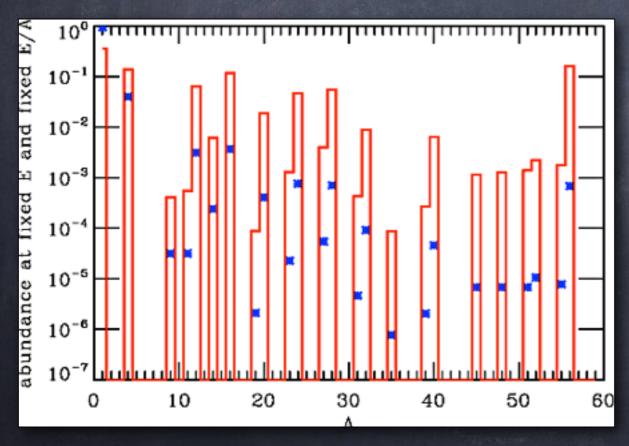
Version 3.0: Luca Maccione, Rafael Alves Batista, David Walz, Gero Müller, Nils Nierstenhoefer, Karl-Heinz Kampert, Peter Schiffer, Arjen van Vliet Astroparticle Physics 42 (2013) 41

Mixed mass compositions

For an injection spectrum $E^{-\alpha}$ elemental abundance at given energy E is modified to

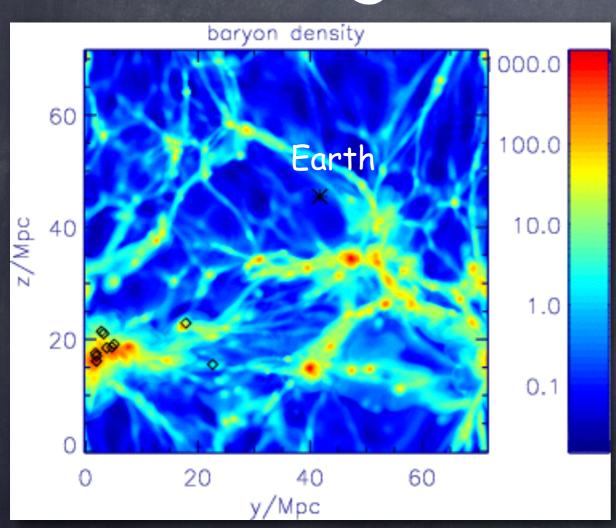
$$\frac{dn_A}{dE}(E) = Nx_A A^{\alpha - 1} E^{-\alpha} g(E)$$

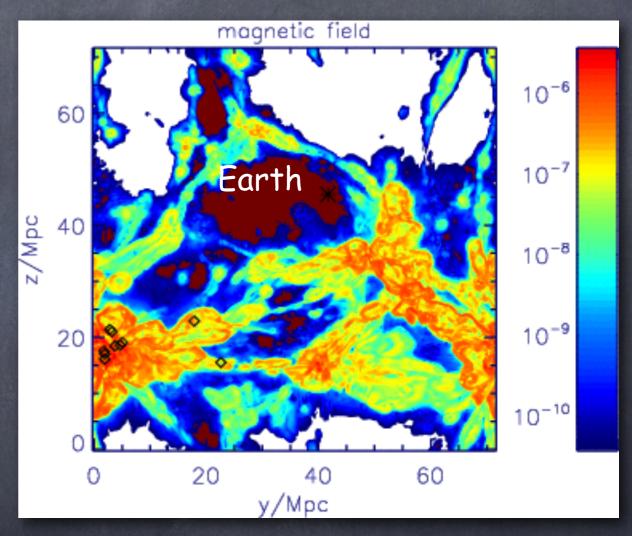
where x_A is the abundance at given energy per nucleon E/A and g(E) is the cut-off shape.



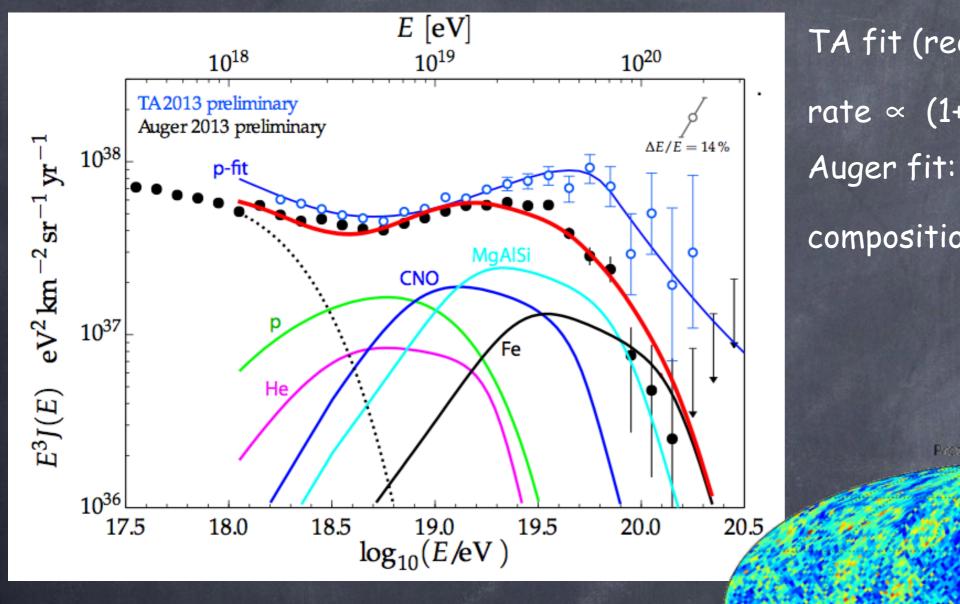
Composition at given E/A (blue) following elemental abundances in the Galaxy Composition at given E for an E^{-2.6} injection spectrum (red).

Discrete Sources in nearby large scale structure

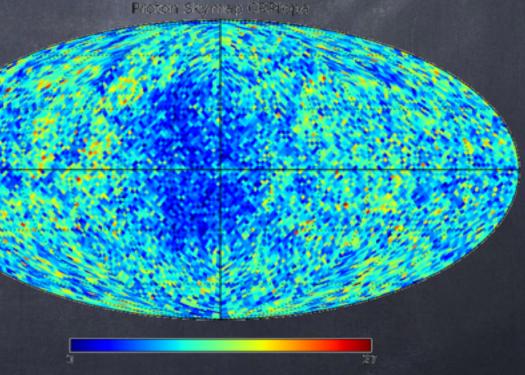




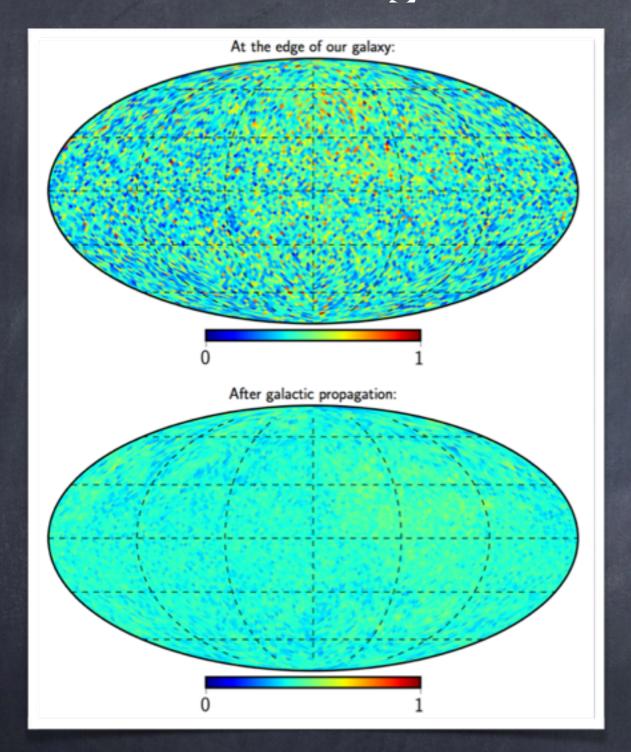
Building Benchmark Scenarios

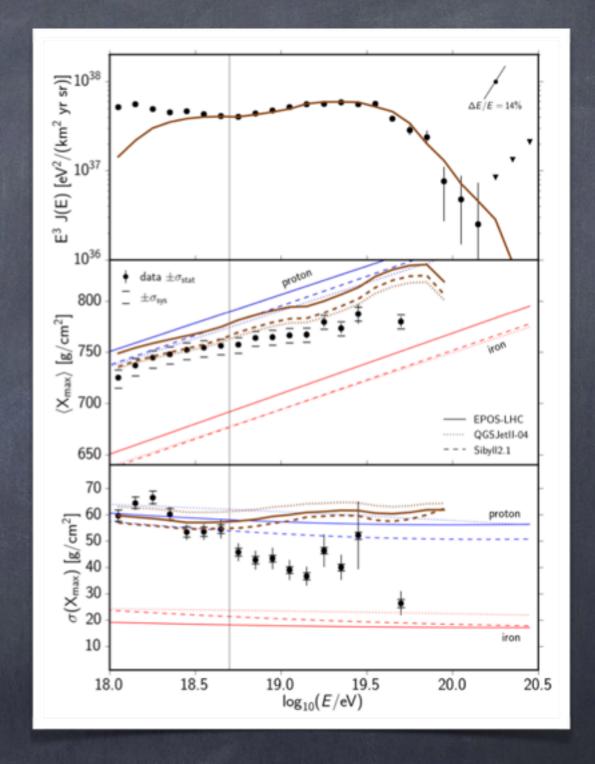


TA fit (red): pure proton injection rate $\propto (1+z)^{4.4} E^{-2.36}$ Auger fit: enhanced galactic composition $\propto E^{-1.8}$ up to $10^{18.7}$ eV*Z



Building Benchmark Scenarios





combining spectral and composition information with anisotropy can considerably strengthen constraints on source characteristics, distributions and magnetization

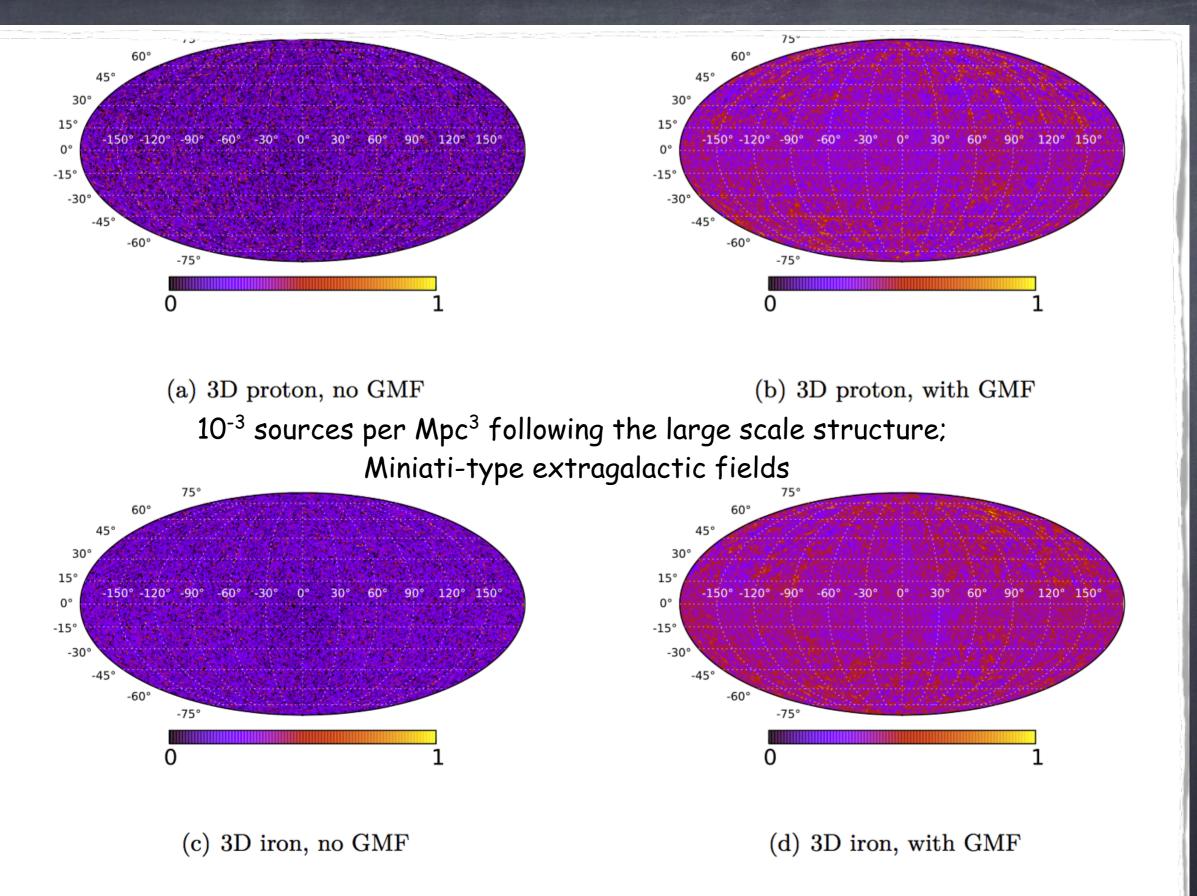
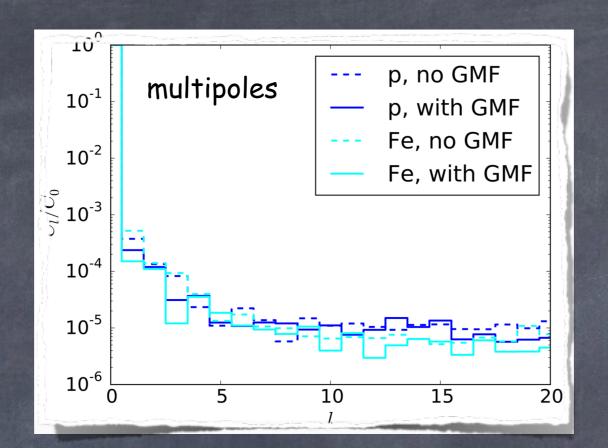
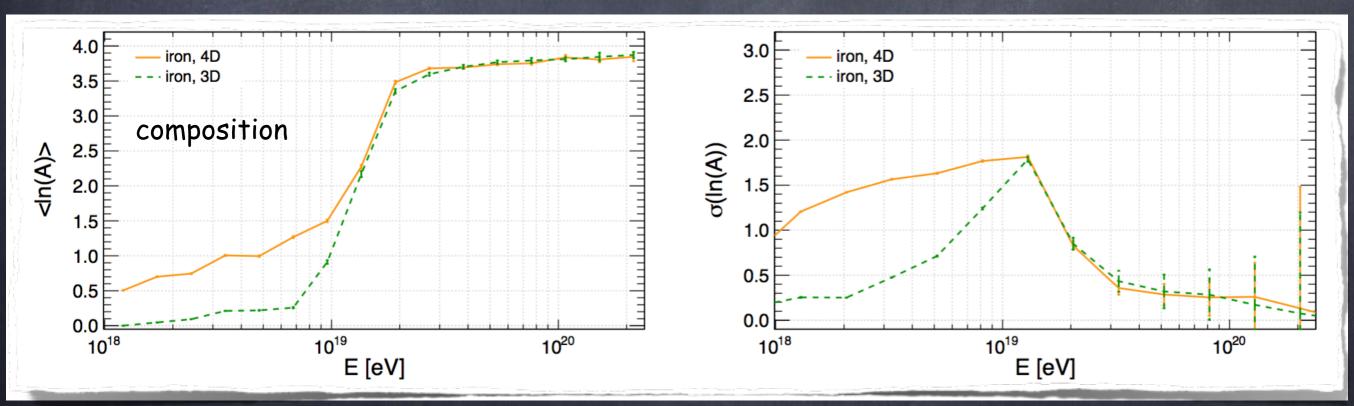
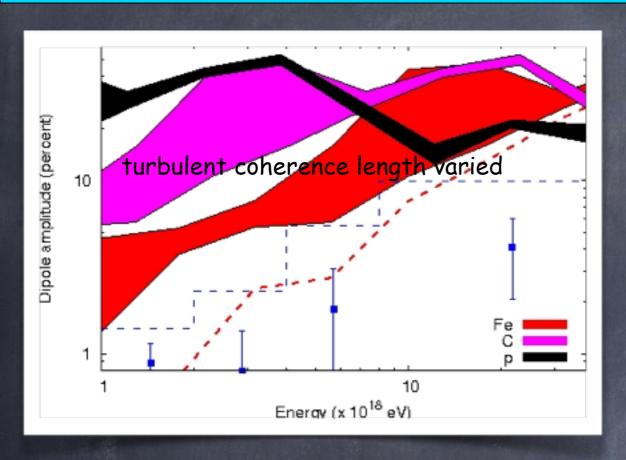


Figure 5. Sky maps for pure proton injection (top) and pure iron injection (bottom) before (left) and after (right) deflections in the galactic magnetic field taken into account using a spectral index of $\gamma = -1.5$ and maximum rigidity $R_{\text{max}} = 500$ EV down to a minimum energy of $E_0 = 1$ EeV. See

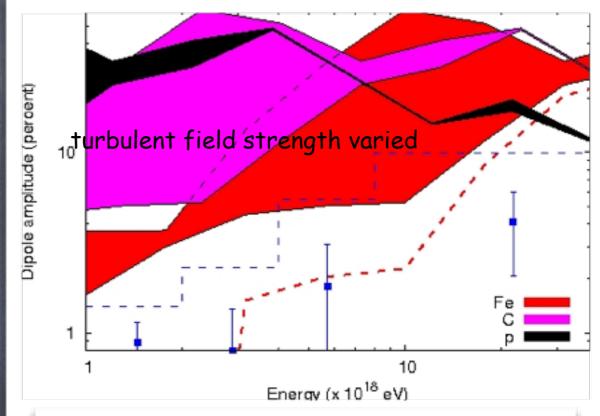


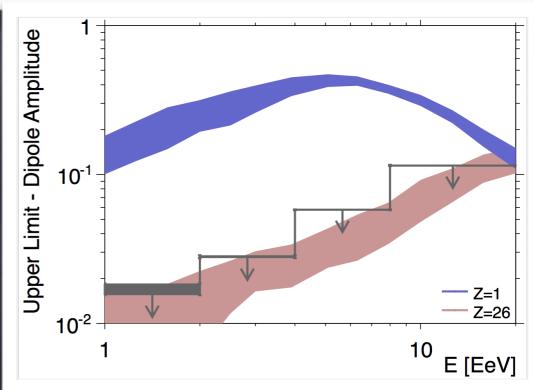


Composition and the Transition Galactic/Extragalactic Cosmic Rays



Giacinti, Kachelriess, Semikoz, Sigl, JCAP 07 (2012) 031 and Pierre Auger Collaboration, Astrophys.J. 762 (2012) L13





Light Galactic Nuclei produce too much anisotropy above ≈ 10¹⁸ eV. This implies:

- 1.) if composition around 10^{18} eV is light => probably extragalactic (and ankle may be due to pair production by protons)
- 2.) if composition around 10^{18} eV is heavy => transition could be at the ankle if Galactic nuclei are produced by sufficiently frequent transients, e.g. magnetars

It is surprisingly difficult to construct simple scenarios with structured sources and magnetic fields that reproduce all observations: spectra, energy dependent composition and anisotropy; to explain them separately is quite easy

Relatively hard injection spectra and low maximal rigidities of few times 10^{18} eV seem to be favored

Conclusions

- 1.) The sources of ultra-high energy cosmic rays are still not identified due to rather small anisotropies; composition seems to become heavier at the highest energies which appears economic in terms of shock acceleration power
- 2.) The observed X_{max} distribution of air showers provides potential constraints on hadronic interaction models: Some models are in tension even when "optimizing" unknown mass composition; however, systematic uncertainties are still high.

Conclusions

- 3.) IceCube neutrinos already constrain their sources which should be sufficiently numerous: Gamma-ray bursts are unlikely as main sources
- 4.) Multi-messenger modeling sources including gamma-rays and neutrinos start to constrain the source and acceleration mechanisms
- 5.) Highest Energy Cosmic Rays, Gamma-rays, and Neutrinos give the strongest constraints on violations of Lorentz symmetry => terms suppressed to first and second order in the Planck mass would have to be unnaturally small