

# Cosmology with Mimetic Matter

Alexander Vikman



26.11.14

# Plan



#### • What is mimetic matter?



- What is mimetic matter?
- Cosmology of "dust" with arbitrary equation of state



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- Conclusions



## This talk is mostly based on

# e-Print: arXiv: **1403.3961,** JCAP 1406 (2014) 017 with *A. H. Chamseddine and V. Mukhanov*

#### and

#### 1412.XXXX with L. Mirzagholi

**Mimetic Matter** 



One can encode the conformal / scalar part of the physical metric  $g_{\mu\nu}$  in a scalar field  $\phi$ :

$$g_{\mu\nu}\left(\tilde{g},\phi\right) = \tilde{g}_{\mu\nu}\,\tilde{g}^{\alpha\beta}\,\partial_{\alpha}\phi\,\partial_{\beta}\phi$$



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 The theory becomes invariant with respect to Weyl transformations:

$$\tilde{g}_{\mu\nu} \to \Omega^2 \left( x \right) \tilde{g}_{\mu\nu}$$



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### But it is still a system with one degree of freedom + standard two polarizations for the graviton!

#### Hamilton-Jacobi equation

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Indeed,

$$u^{\lambda} \nabla_{\lambda} u_{\alpha} = \nabla^{\lambda} \phi \nabla_{\lambda} \nabla_{\alpha} \phi = \frac{1}{2} \nabla_{\alpha} \left( \nabla^{\lambda} \phi \nabla_{\lambda} \phi \right) = 0$$

# **Dissformal Transformation**

Nathalie Deruelle and Josephine Rua (2014)

One obtains the same dynamics (*the same Einstein equations*), if instead of varying the Einstein-Hilbert action with respect to the metric  $g_{\mu\nu}$ one plugs in a *dissformal transformation*   $g_{\mu\nu} = F(\Psi, w) \ell_{\mu\nu} + H(\Psi, w) \partial_{\mu} \Psi \partial_{\nu} \Psi$ with  $w = \ell^{\mu\nu} \partial_{\mu} \Psi \partial_{\nu} \Psi$  and  $w^2 F \frac{\partial}{\partial w} \left(H + \frac{F}{w}\right) \neq 0$ 

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Chamseddine, Mukhanov; Golovnev; Barvinsky (2013) Lim, Sawicki, Vikman; (2010)

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# "Cold Dark Matter"?



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Lim, Sawicki, Vikman; (2010) Chamseddine, Mukhanov, Vikman (2014)

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#### **Enough freedom to obtain any cosmological evolution!**

• In particular  $V(\phi) = \frac{1}{3} \frac{m^4 \phi^2}{e^{\phi} + 1}$  gives the same cosmological inflation as  $\frac{1}{2}m^2\phi^2$  potential in the standard case

The simplest way to generate needed cosmological evolution

$$\ddot{\xi} - \frac{3}{4}V(t)\xi = 0$$
  
where  $\xi = a^{3/2}$ 

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#### even a bouncing universe is possible!

e.g. for 
$$V(\phi) = \frac{4}{3} \frac{\alpha}{(\phi^2 + \phi_0^2)^2}$$
## **Perturbations** I

Lim, Sawicki, Vikman; (2010) Chamseddine, Mukhanov, Vikman (2014)

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# Even with potential, the energy still moves along the timelike geodesics

 $c_{\rm S}=0$ 

Newtonian potential:

$$\Phi = C_1 \left( \mathbf{x} \right) \left( 1 - \frac{H}{a} \int a dt \right) + \frac{H}{a} C_2 \left( \mathbf{x} \right)$$

Here on **all scales** but in the usual cosmology it is an approximation for **superhorizon** scales

Chamseddine, Mukhanov, Vikman (2014)

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Just add higher derivatives !

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#### Back to waves, oscillators and normal quantum fluctuations!

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- Weigher time derivatives can be eliminated just by the differentiation of this Hamilton-Jacobi equation
- There are only minor changes (rescaling) in the background evolution equations e.g.

$$2\dot{H} + 3H^2 = \frac{2}{2 - 3\gamma}V(t)$$

Mirzagholi, Vikman (to appear 2014)

no potential

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expansion  $\theta = \nabla_{\mu} u^{\mu}$ 

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pressure  $p = -\gamma \left(\dot{\theta} + \frac{1}{2}\theta^2\right)$ 

option for DE  $p \simeq -\varepsilon$ 

#### Violation of the Einstein's equivalence principle?

Raychaudhuri at work!

$$\dot{\theta} = -\frac{1}{3}\theta^2 - \sigma^2 - R_{\mu\nu}u^{\mu}u^{\nu}$$
$$R_{\mu\nu} = T_{\mu\nu} - \frac{1}{2}Tg_{\mu\nu}$$

Mirzagholi, Vikman (to appear 2014)

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Noether current:

$$J_{\mu} = n u_{\mu} + q_{\mu}$$

Mirzagholi, Vikman (to appear 2014)

no potential  $\phi \rightarrow \phi + c$  symmetry  $\nabla_{\mu}J^{\mu} = 0$ Noether current:  $J_{\mu} = nu_{\mu} + q_{\mu}$ charge density  $n=2\lambda-\gamma \dot{\theta}$ 

Charge conservation  $n = 2\lambda - \gamma \dot{\theta} \propto a^{-3}$ 

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Inflation  $\lambda_{\star} = \gamma \theta$ 

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Inflation  $\lambda_{\star} = \gamma \dot{\theta}$ 

$$\frac{p_{\star}}{\varepsilon_{\star}} = -1 - \frac{2}{3} \frac{\dot{H}}{H^2} = w_{\Sigma}$$

$$\varepsilon_{\star} = \frac{12\pi\gamma}{1 - 12\pi\gamma} \rho_{\rm ext}$$

## **Perturbations HD**

Chamseddine, Mukhanov, Vikman (2014)

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Newtonian potential:  $\Phi=\delta\phi$ 

# Quantization

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the action

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short wavelength quantum fluctuations

$$\delta\phi_k \sim \sqrt{\frac{c_s}{\gamma}} \ k^{-3/2}$$

match with the long-wave-length limit

#### **Perturbations in Mimetic Inflation**

Chamseddine, Mukhanov, Vikman (2014)

on scale  $\lambda \simeq k^{-1}$
Chamseddine, Mukhanov, Vikman (2014)

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Newtonian potential  $\Phi_{\lambda} \simeq \sqrt{c_{\rm s}}/\gamma \cdot H_{c_{\rm s}} k \simeq H$ 

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  - Sequences spectral indices  $n_{\rm S} 1 = n_{\rm T}$

Chamseddine, Mukhanov, Vikman (2014)

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 $\begin{aligned} & & & \\ \hline {\bf \Theta} \ {\rm spectral \ indices} \ \ n_{\rm S} - 1 = n_{\rm T} \\ & & \\ \hline {\bf \Theta} \ {\rm small \ sound \ speed} \ \ {\bf \Phi}_{\lambda} \gg h_{\lambda} \end{aligned}$ 

But it seems that there is no usual Non-Gaussianity!

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Thanks a lot for attention!